

Ripples of the QCD Critical Point



Shi Yin

15.05.2024

Based on:

- [1] W.-j. Fu, X.-f. Luo, J. M. Pawłowski, F. Rennecke, R. Wen, S. Yin; Phys. Rev. D.104 (2021), 094047.**
- [2] W.-j. Fu, X.-f. Luo, J. M. Pawłowski, F. Rennecke, S. Yin; arXiv: 2308.15508.**

Outline

- Self introduction
- Introduction
- Looking for ripples
- Summary

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Self introduction

Name: Shi Yin (尹诗)

Education: Dalian University of Technology, Theoretical Physics

Research area: QCD phase structure, Relativistic heavy ion collision phenomenology

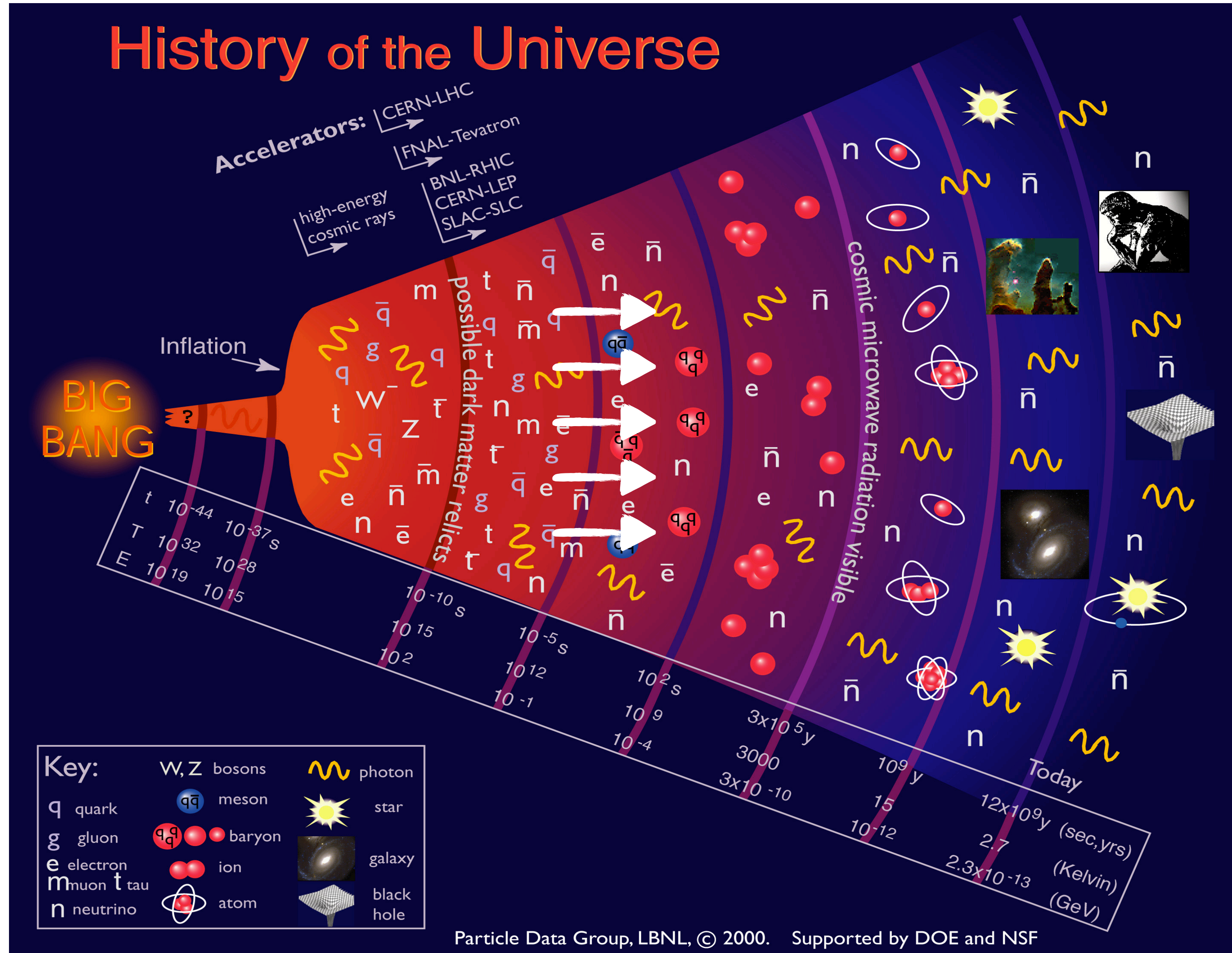
PostDoc: JLU, supported by Alexander von Humboldt foundation



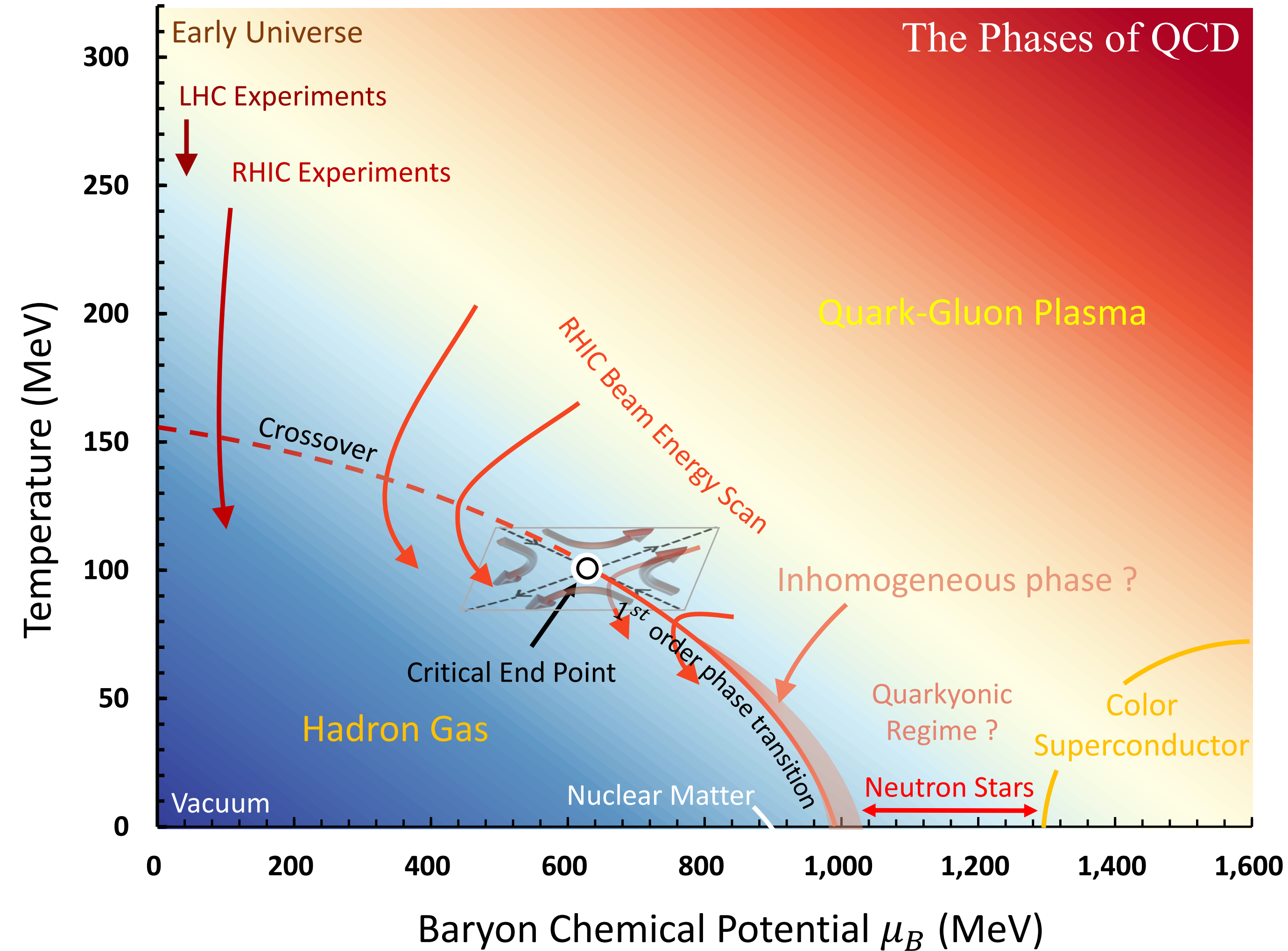
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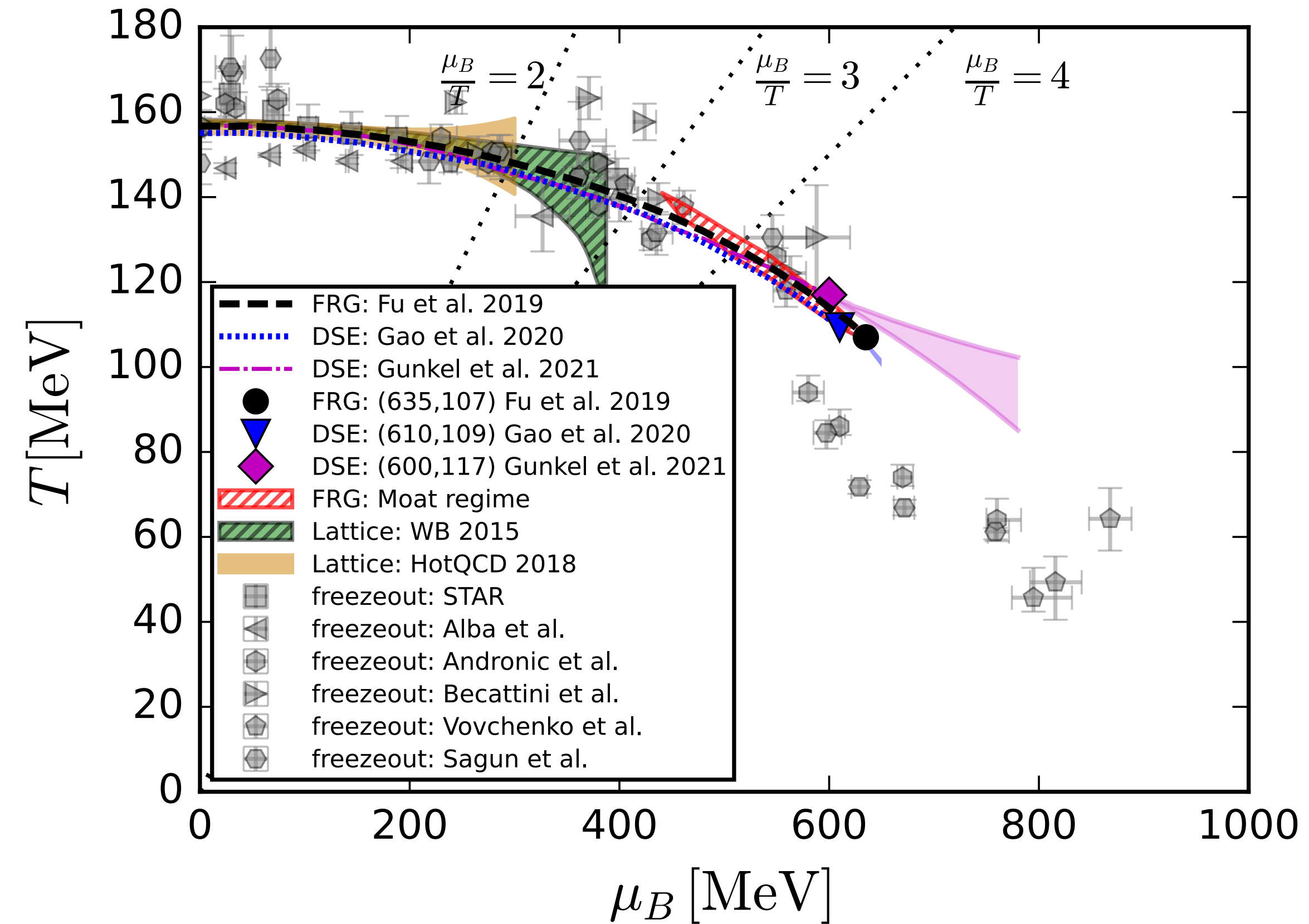
Introduction



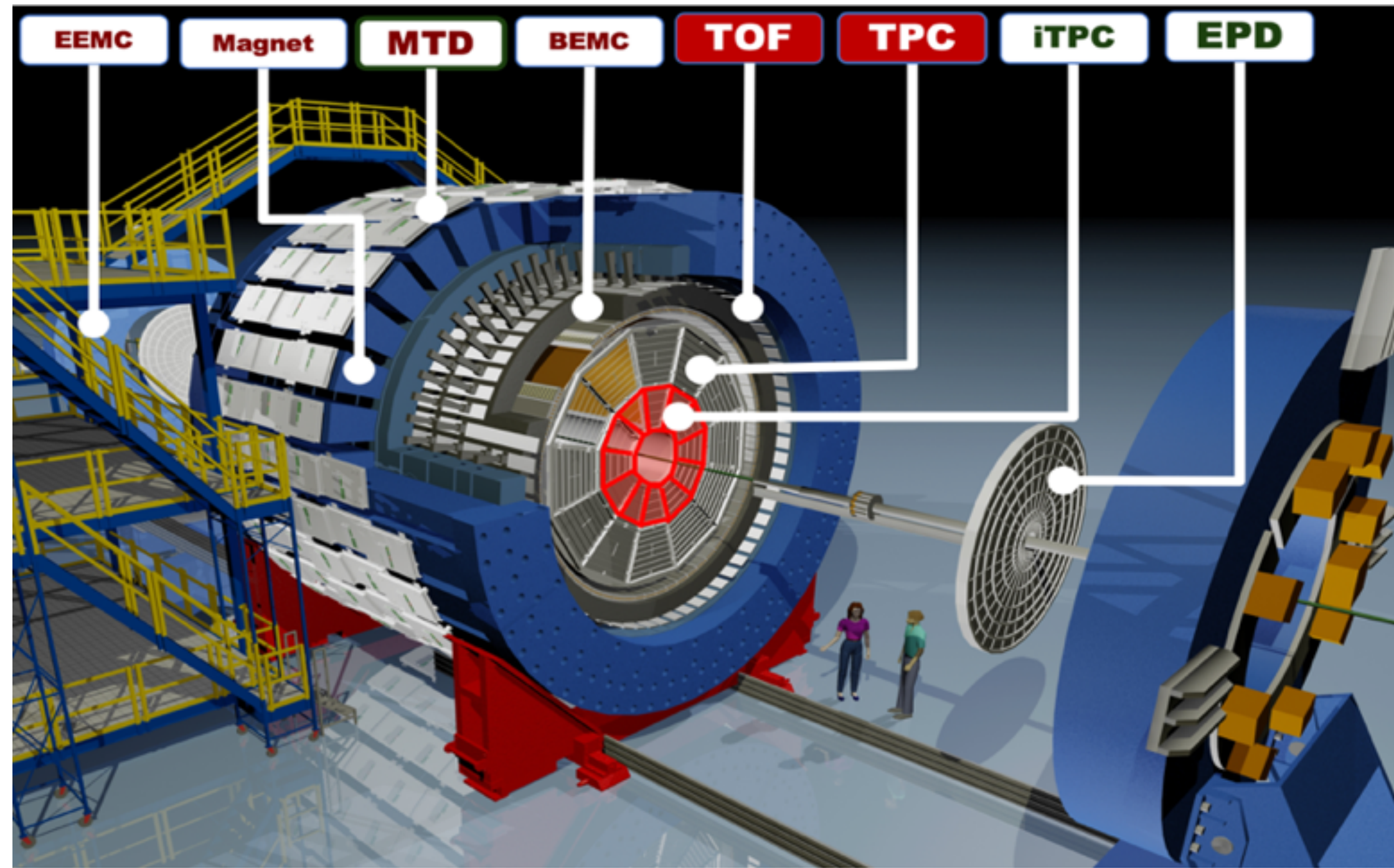
Phase transition



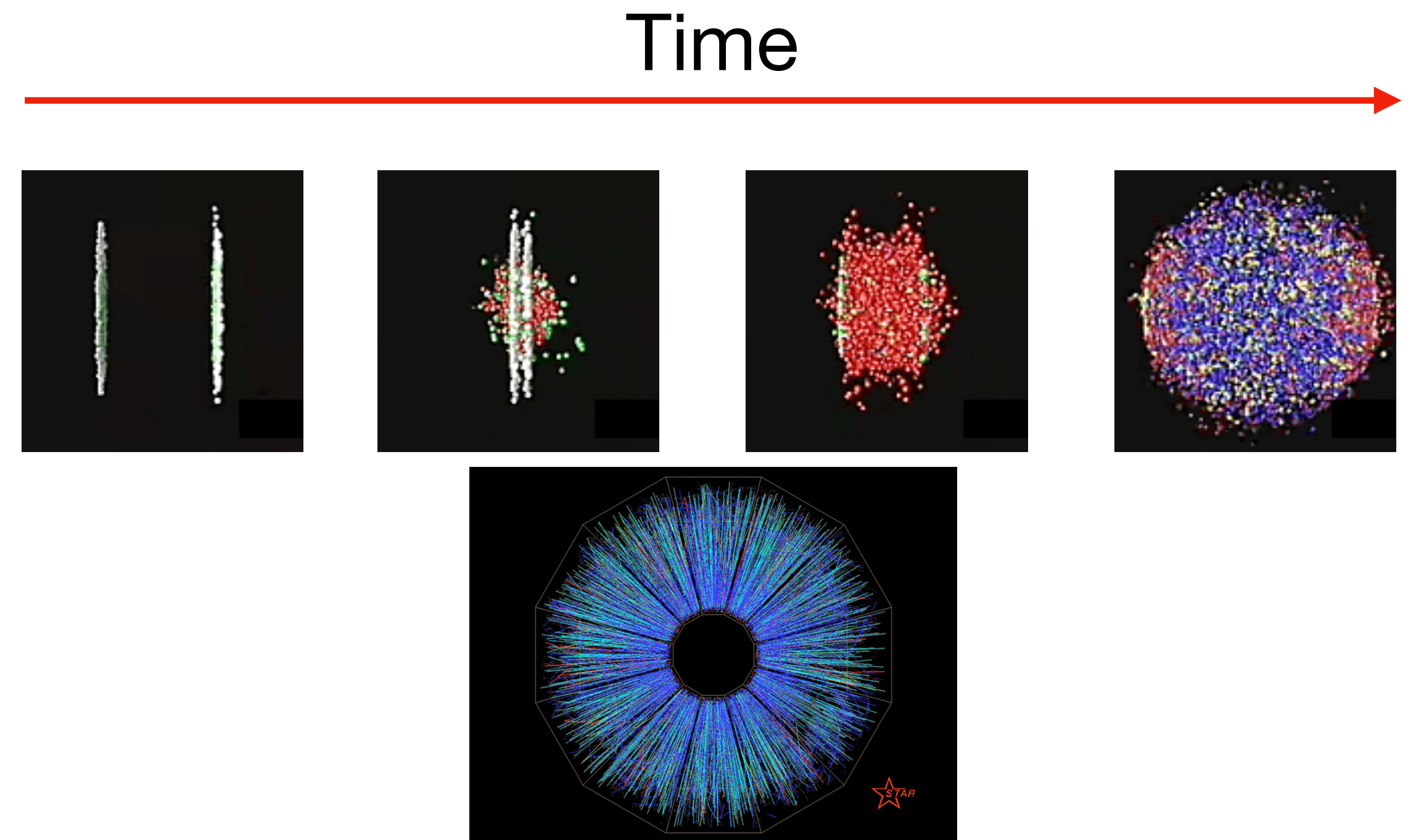
Collection of phase boundaries



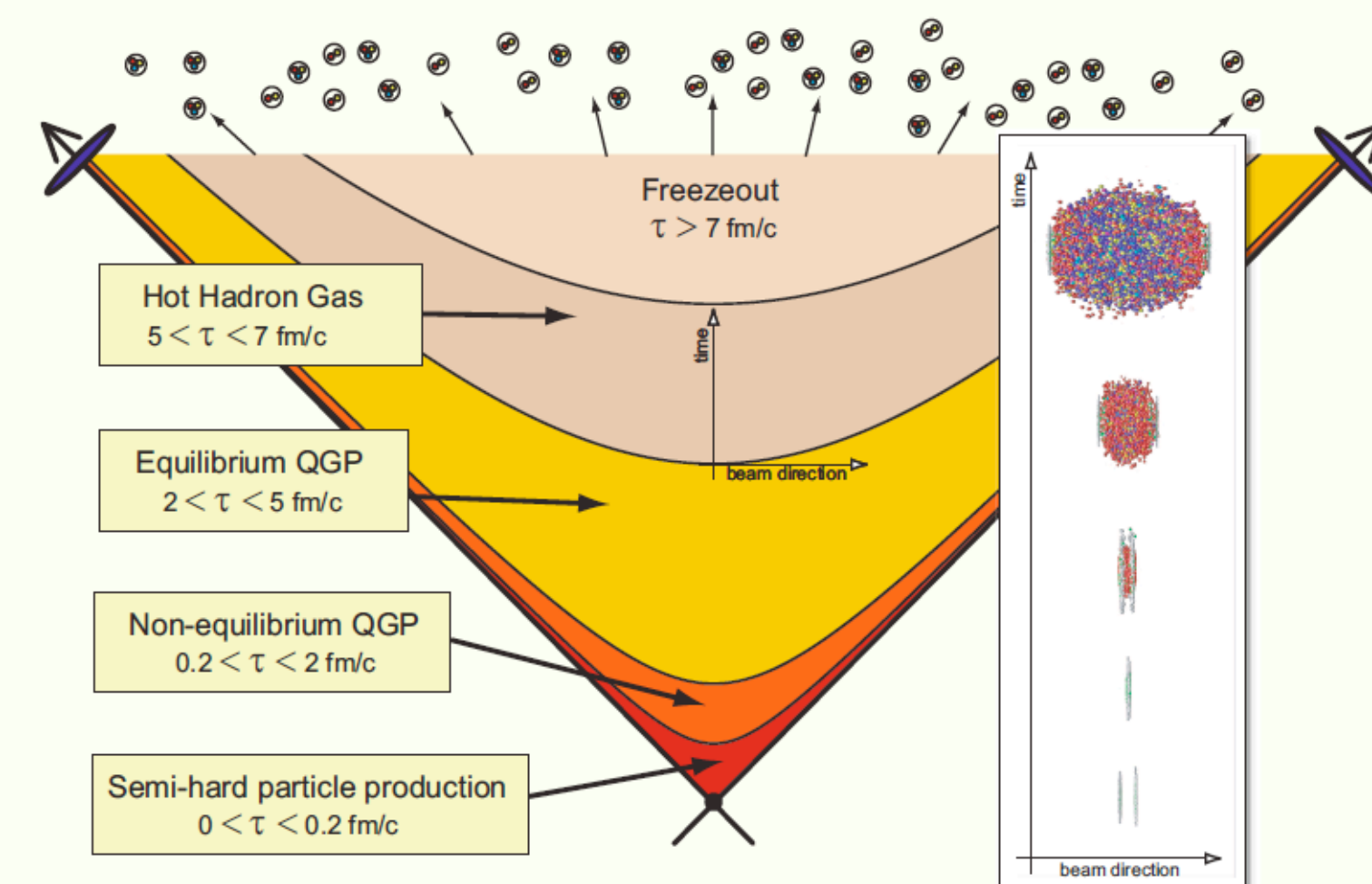
Heavy Ion Collision



STAR Collaboration, K. H. Ackermann et. al., STAR detector overview, Nucl.Instrum.Meth.A 499 (2003) 624-632

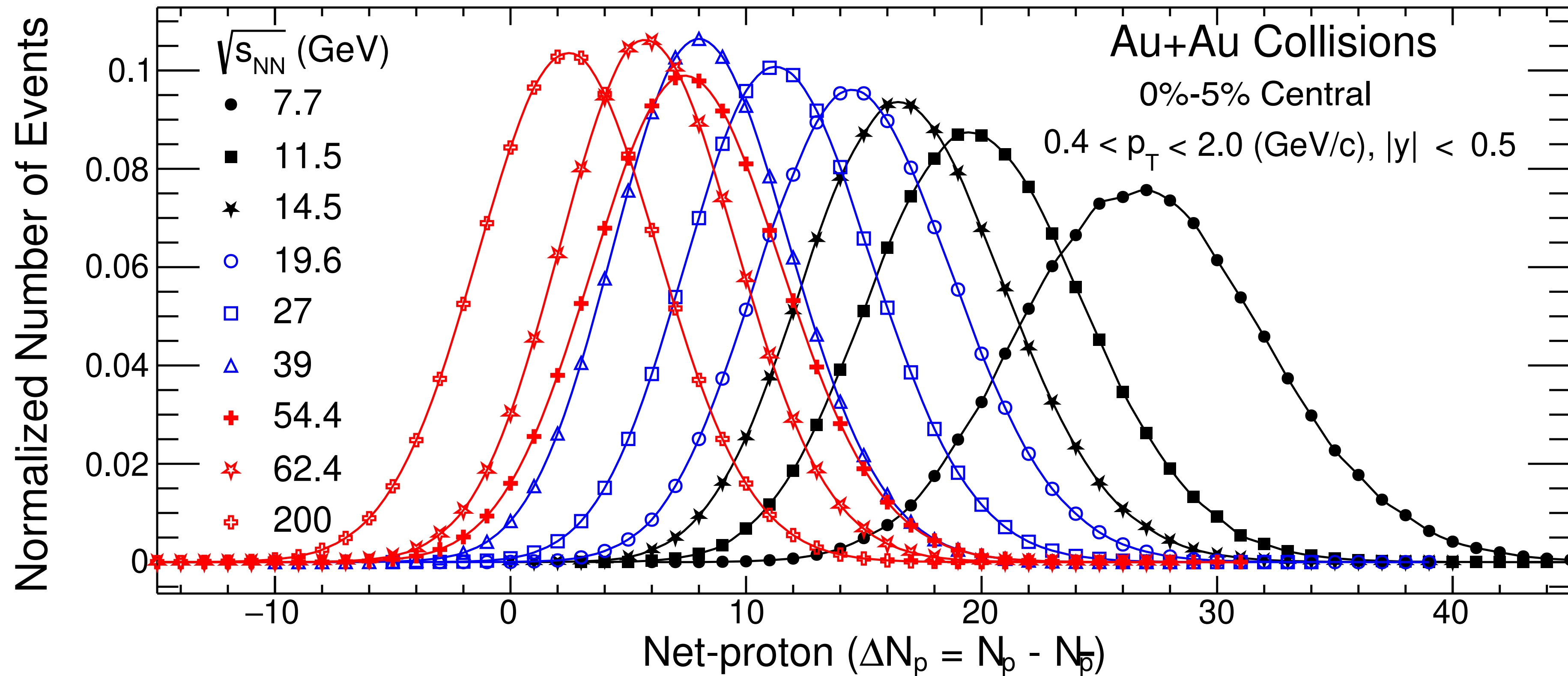


Heavy-ion collision timescales and “epochs” @ RHIC



*1 fm/c $\simeq 3 \times 10^{-24}$ seconds

Baryon number distribution



J. Adam et. Al., *Phys.Rev.Lett.* 126 (2021) 9, 092301

Baryon number fluctuations

$$\chi_1^B = \frac{1}{VT^3} \langle N_B \rangle,$$

$$\chi_2^B = \frac{1}{VT^3} \langle (\delta N_B)^2 \rangle,$$

$$\chi_3^B = \frac{1}{VT^3} \langle (\delta N_B)^3 \rangle,$$

$$\chi_4^B = \frac{1}{VT^3} (\langle (\delta N_B)^4 \rangle - 3\langle (\delta N_B)^2 \rangle^2), \dots$$

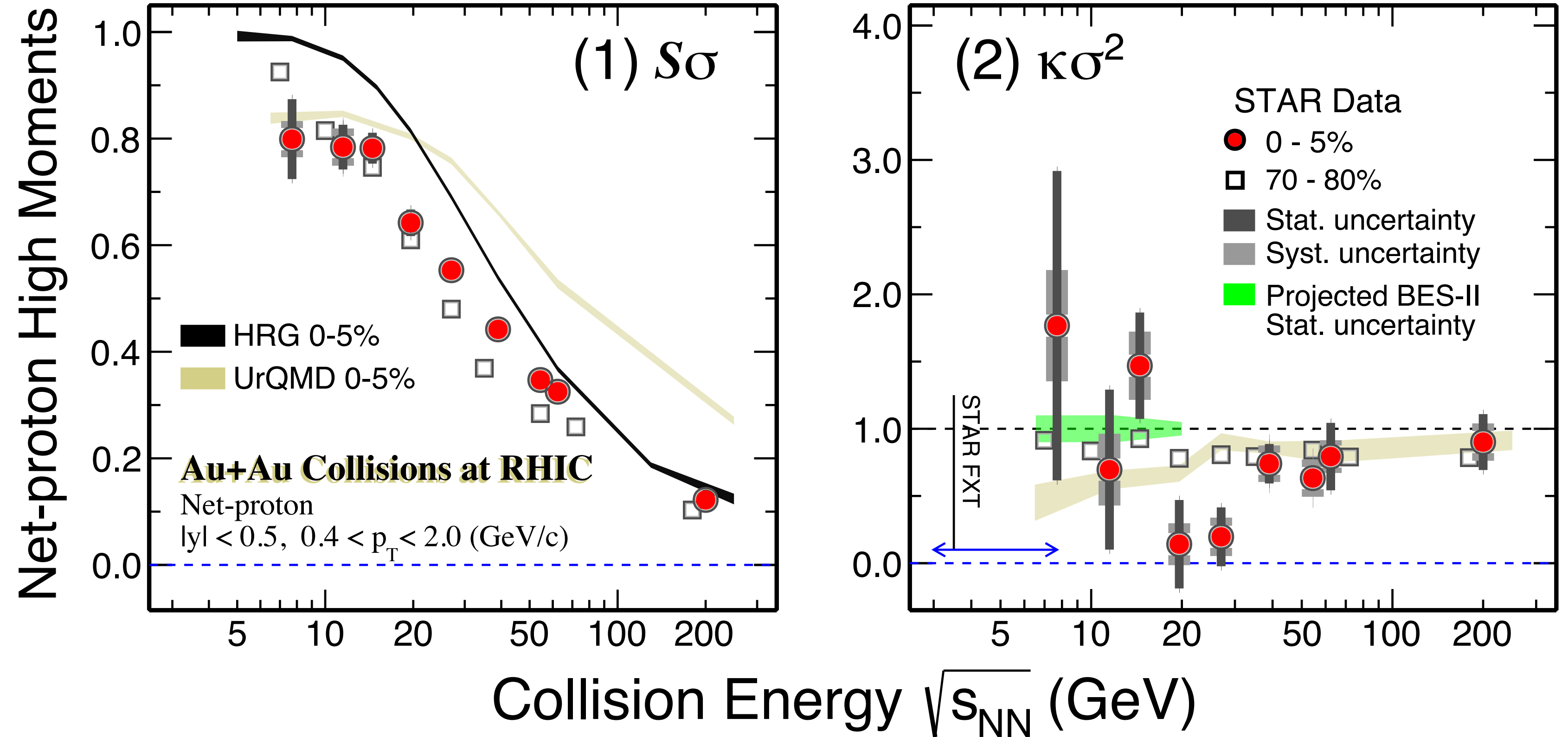
$$\delta N_B = N_B - \langle N_B \rangle$$

Mean value $M = VT^3 \chi_1^B$

Variance $\sigma^2 = VT^3 \chi_2^B$

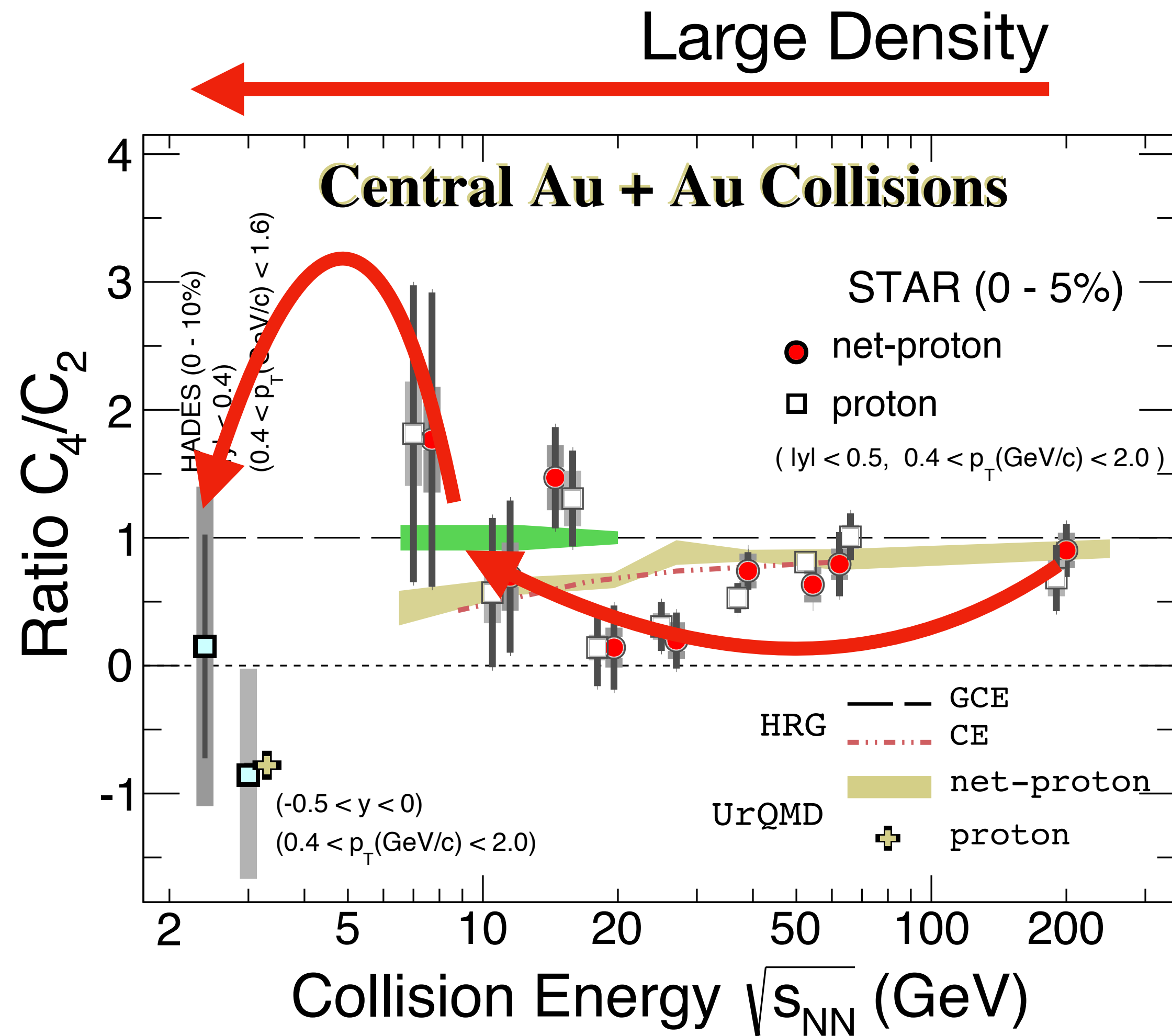
Skewness $S = \frac{\chi_3^B}{\chi_2^B \sigma}$

Kurtosis $\kappa = \frac{\chi_4^B}{\chi_2^B \sigma^2}$

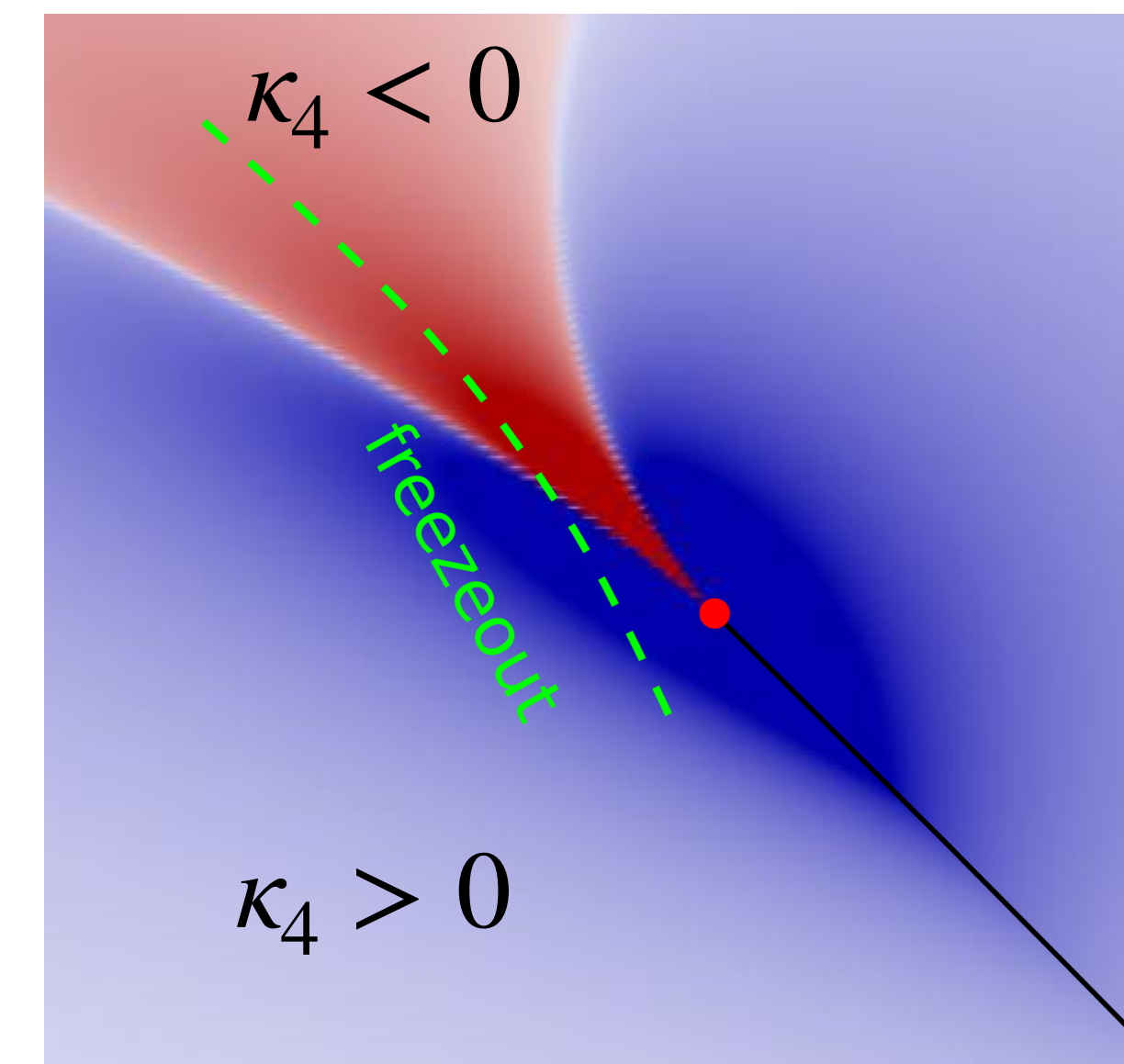
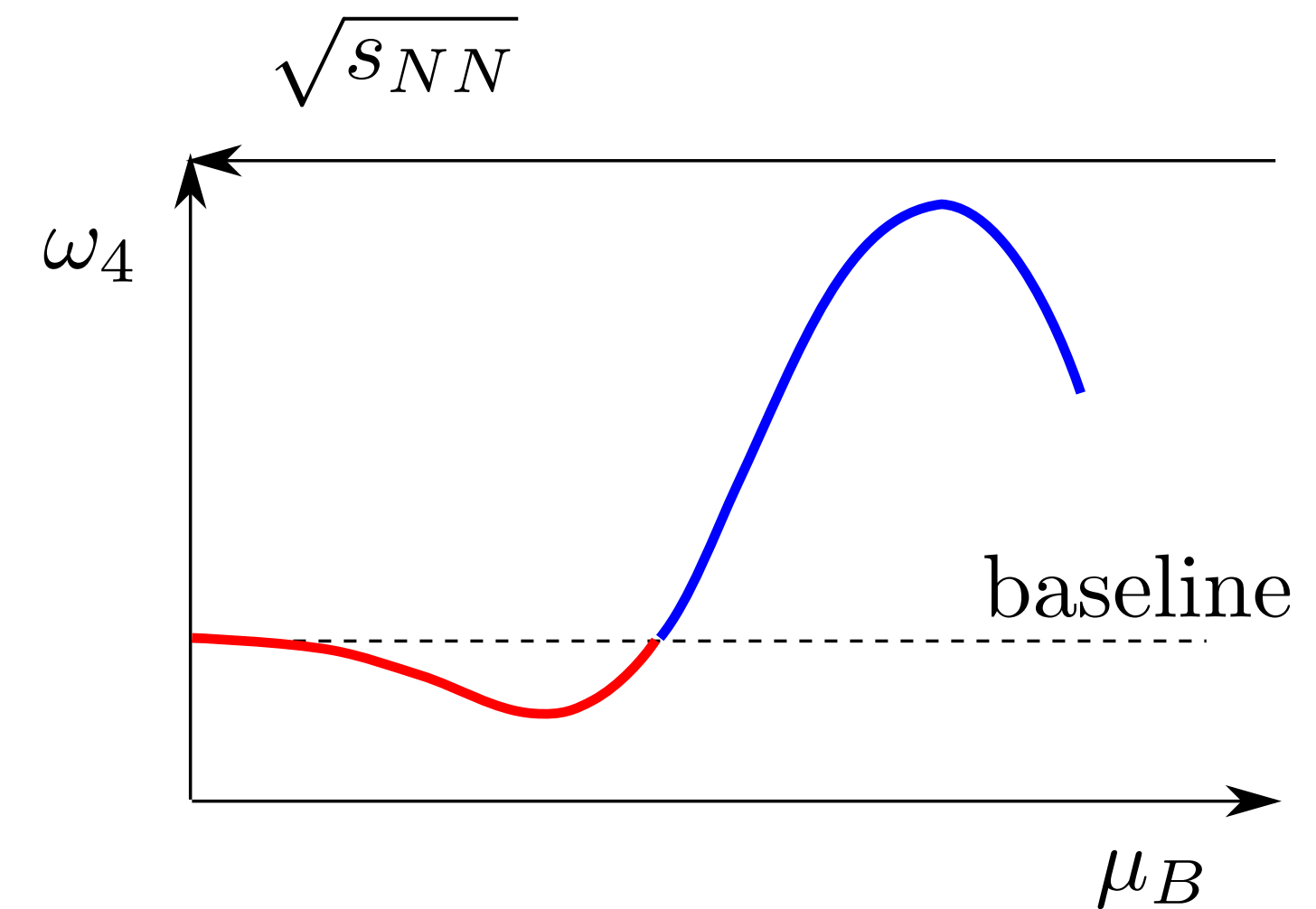


J. Adam *et al.* (STAR), arXiv: 2001.02852

Baryon number fluctuations



M. S. Abdallah et. al., *Phys.Rev.Lett.* 128 (2022) 20, 202303



A.Aprahamian, et al., **Reaching for the horizon: The 2015 long range plan for nuclear science(2015).**

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Looking for ripples

First Step: QCD-assisted LEFT within fRG

$$N_f = 2$$

$$\Gamma_k = \int_x \left\{ Z_{q,k} \bar{q} \left[\gamma_\mu \partial_\mu - \gamma_0 (\mu + igA_0) \right] q + \frac{1}{2} Z_{\phi,k} (\partial_\mu \phi)^2 + h_k \bar{q} (\tau^0 \sigma + \vec{\tau} \cdot \vec{\pi}) q + V_k(\rho, A_0) - c\sigma \right\}$$

$$\Omega[T, \mu_B] = V_{\text{glue}}(L, \bar{L}) + V_{\text{matter}}(\rho, L, \bar{L}) - c\sigma$$

$$p = -\Omega[T, \mu_B]$$

$$\chi_n^B = \frac{\partial^n}{\partial (\mu_B/T)^n} \frac{p}{T^4}$$

$$R_{nm}^B = \frac{\chi_n^B}{\chi_m^B}$$

$$\partial_t \Gamma_k[\Phi] = - \text{[Diagram 1]} + \frac{1}{2} \text{[Diagram 2]}$$

Only Grand Canonical Ensemble !!!

QCD-assisted LEFT

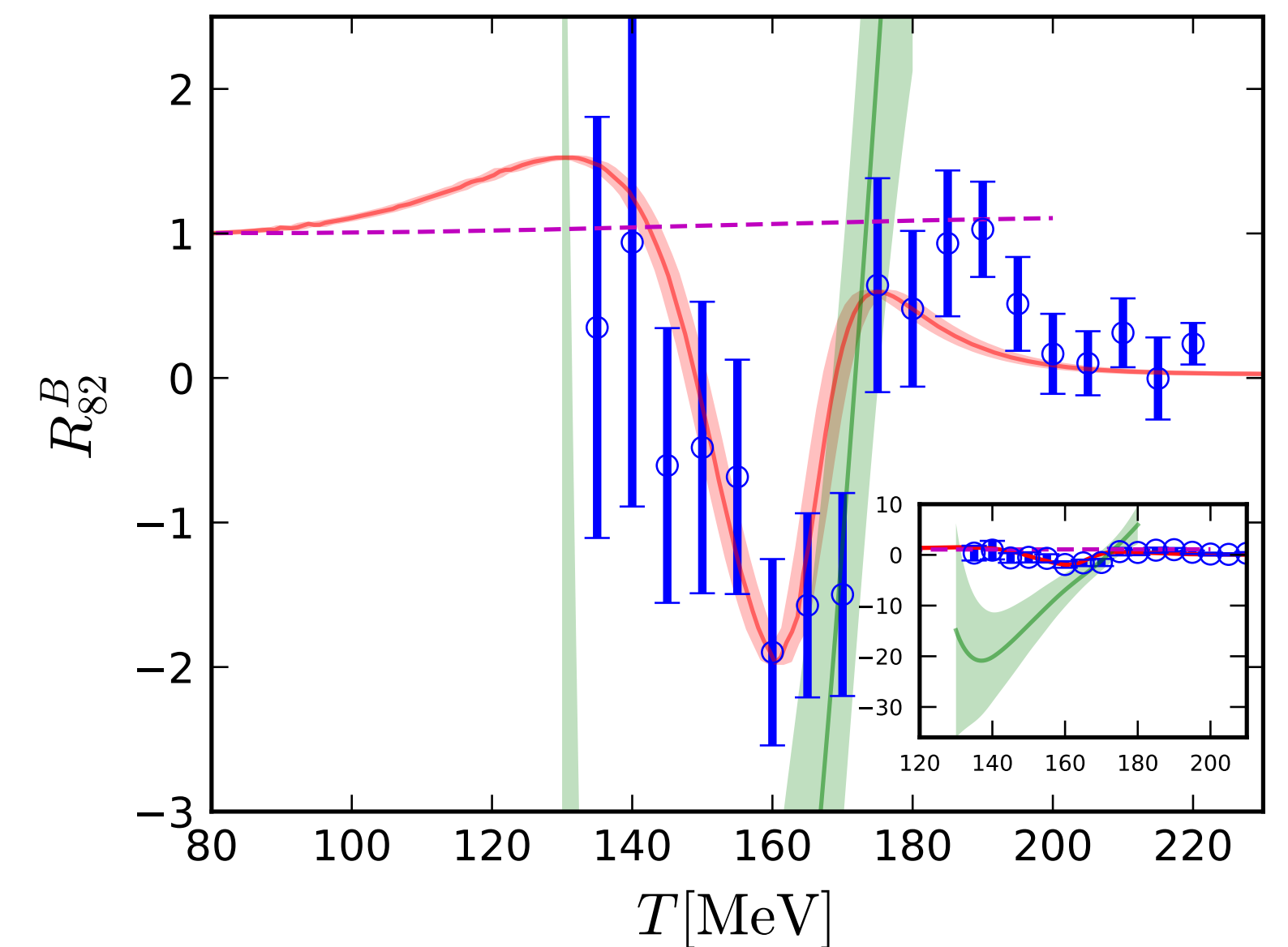
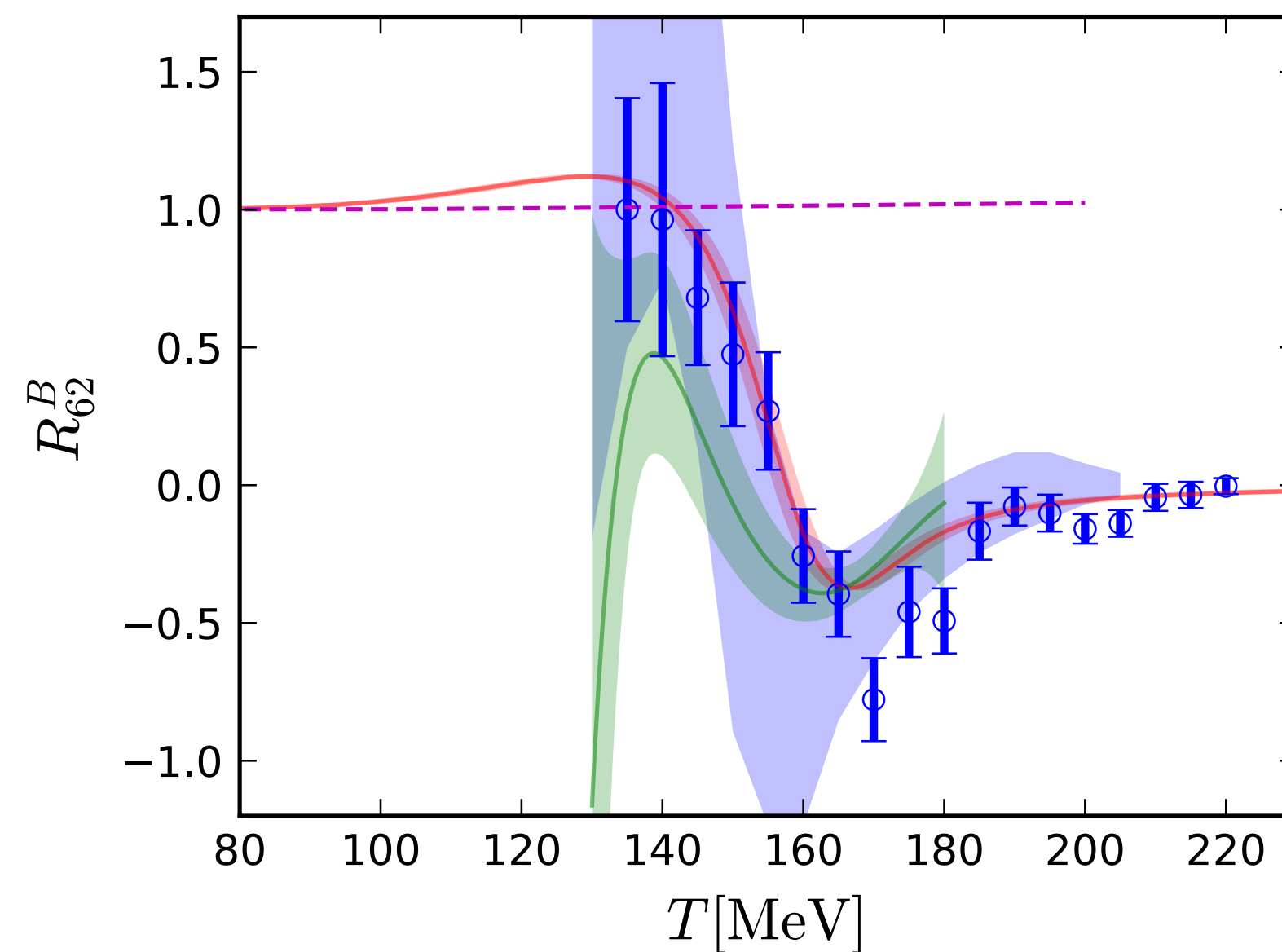
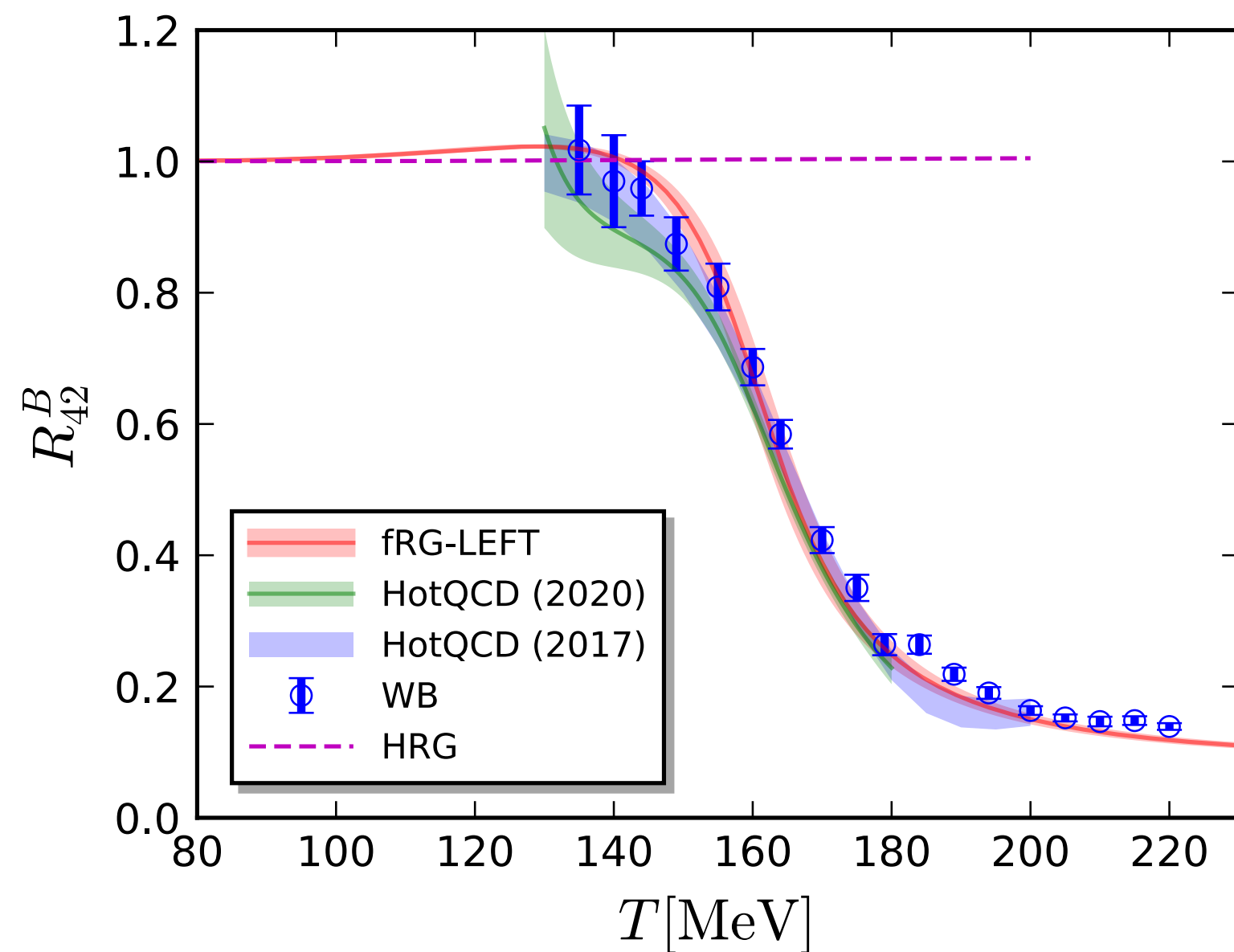
Benchmarked by linear scale-matching

$$T^{\text{LEFT}} = c_T T^{\text{QCD}}$$

$$\mu_B^{\text{LEFT}} = c_\mu \mu_B^{\text{QCD}}$$

* c_T is determined by matching the pseudo-critical temperatures

* c_μ is determined by matching the curvature of the phase boundaries



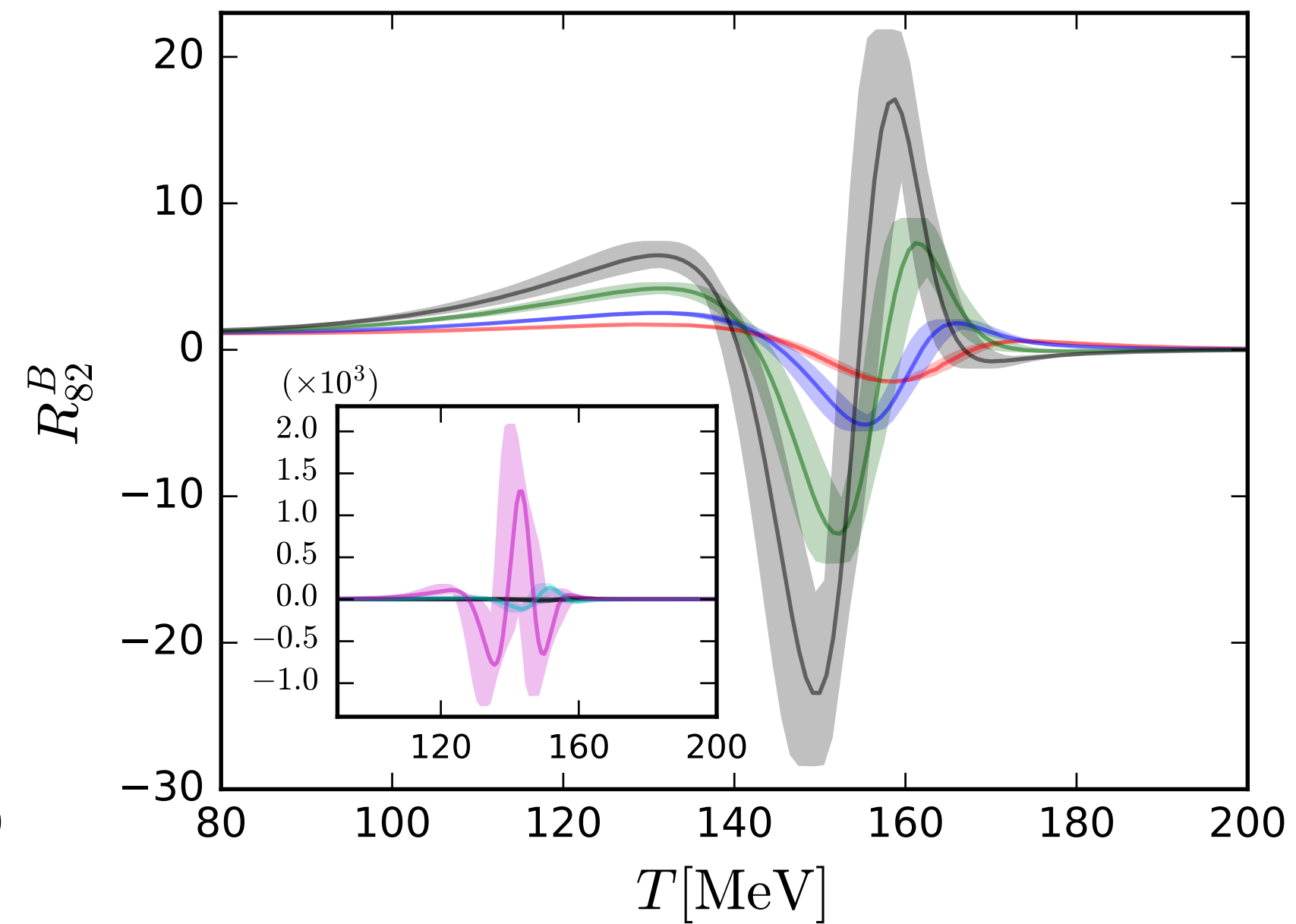
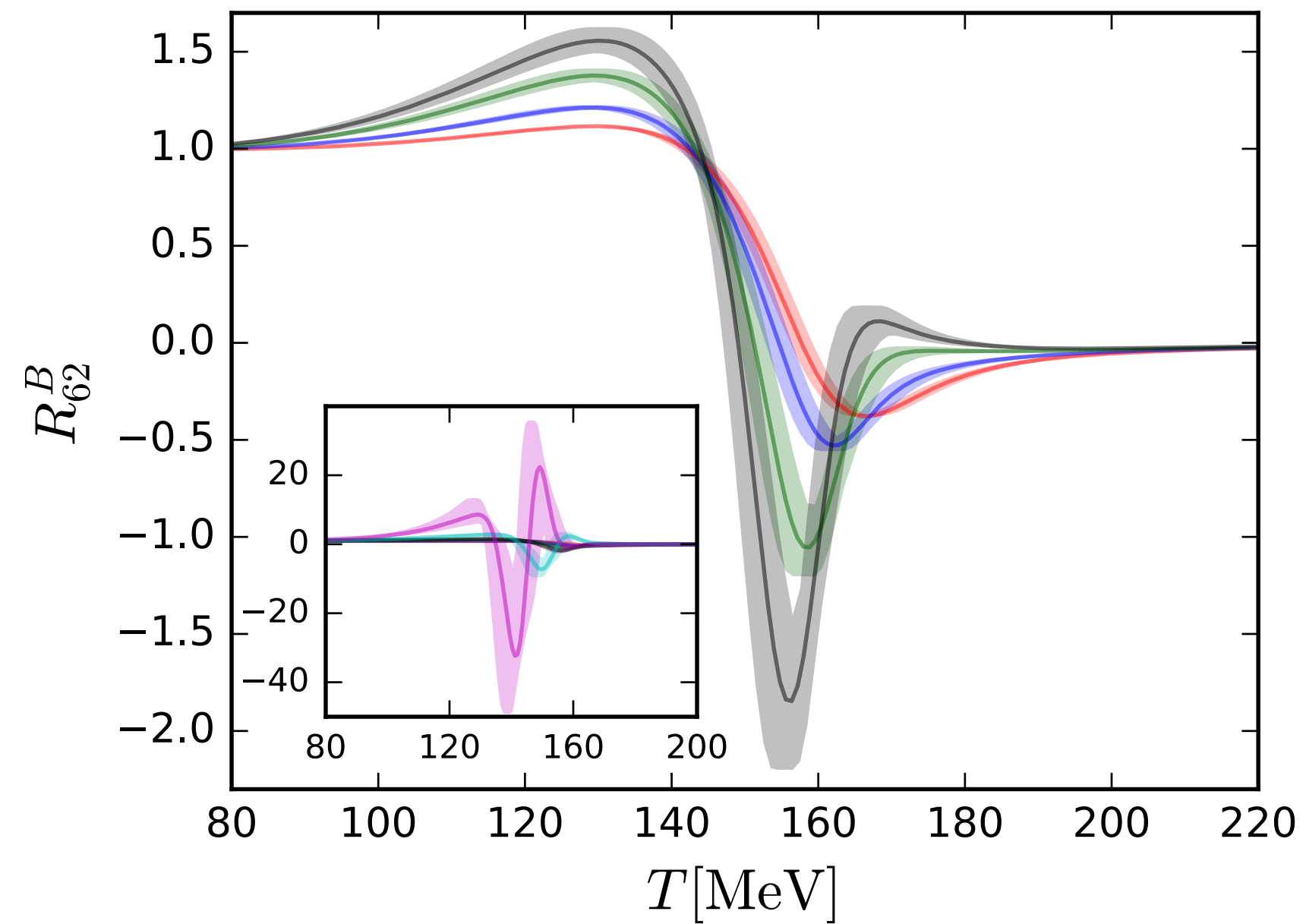
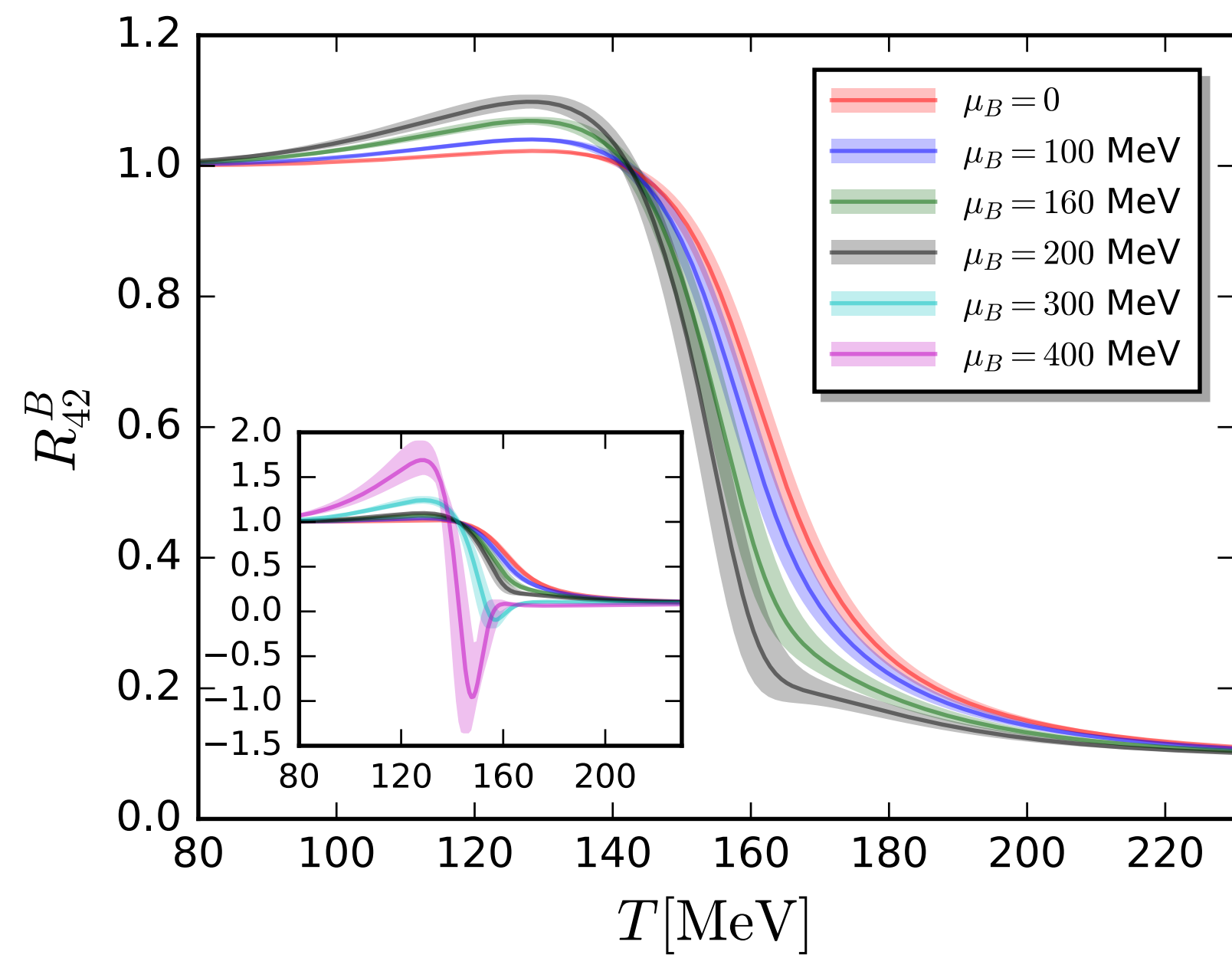
fRG-LEFT: Fu, Luo, Pawłowski, Rennecke, Wen, SY, *PRD* 104 (2021) ,094047

HotQCD: A. Bazavov *et al.*, *PRD* 95 (2017), 054504; *PRD* 101 (2020), 074502

WB: S. Borsanyi *et al.*, *JHEP* 10 (2018) 205

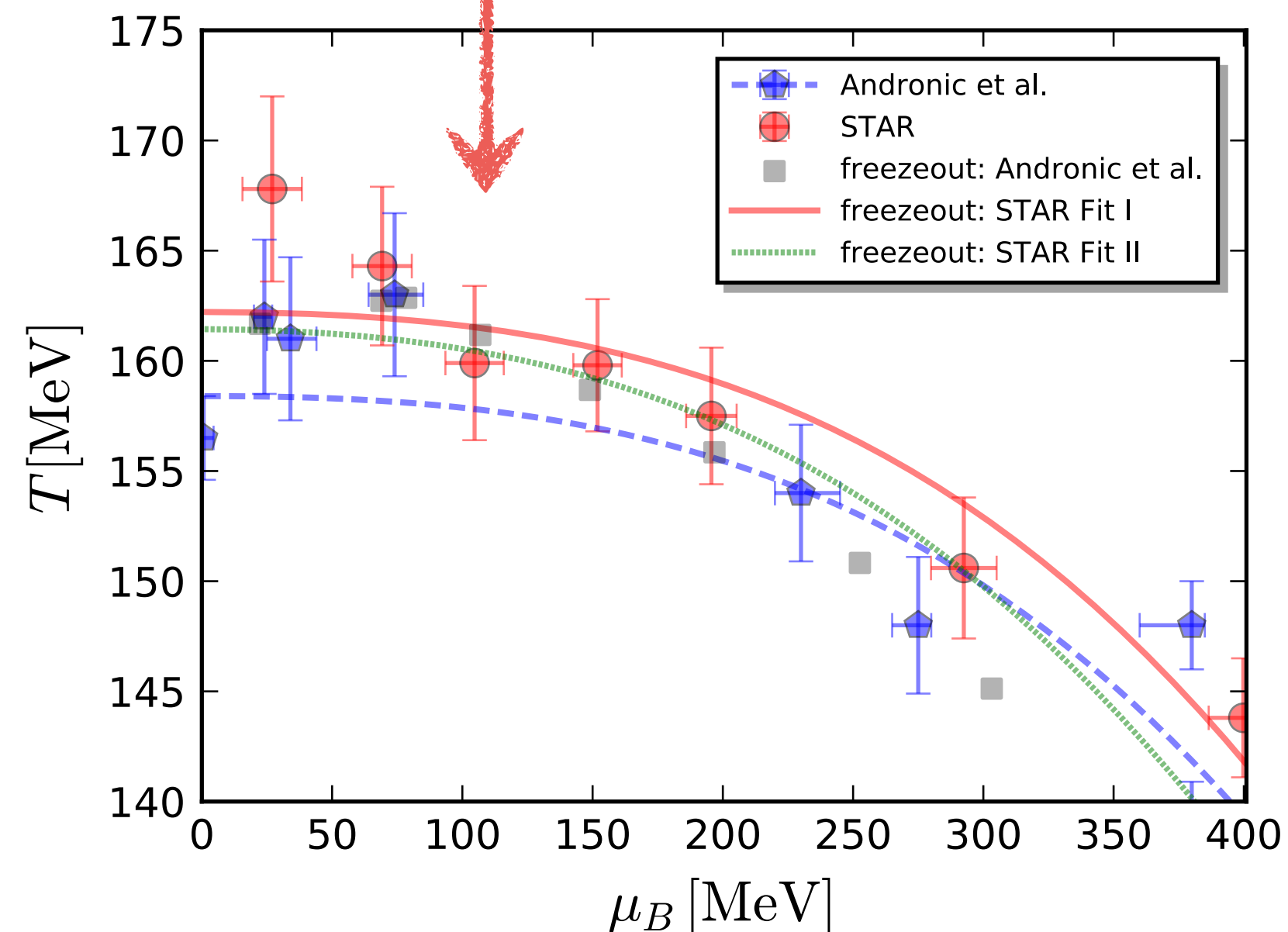
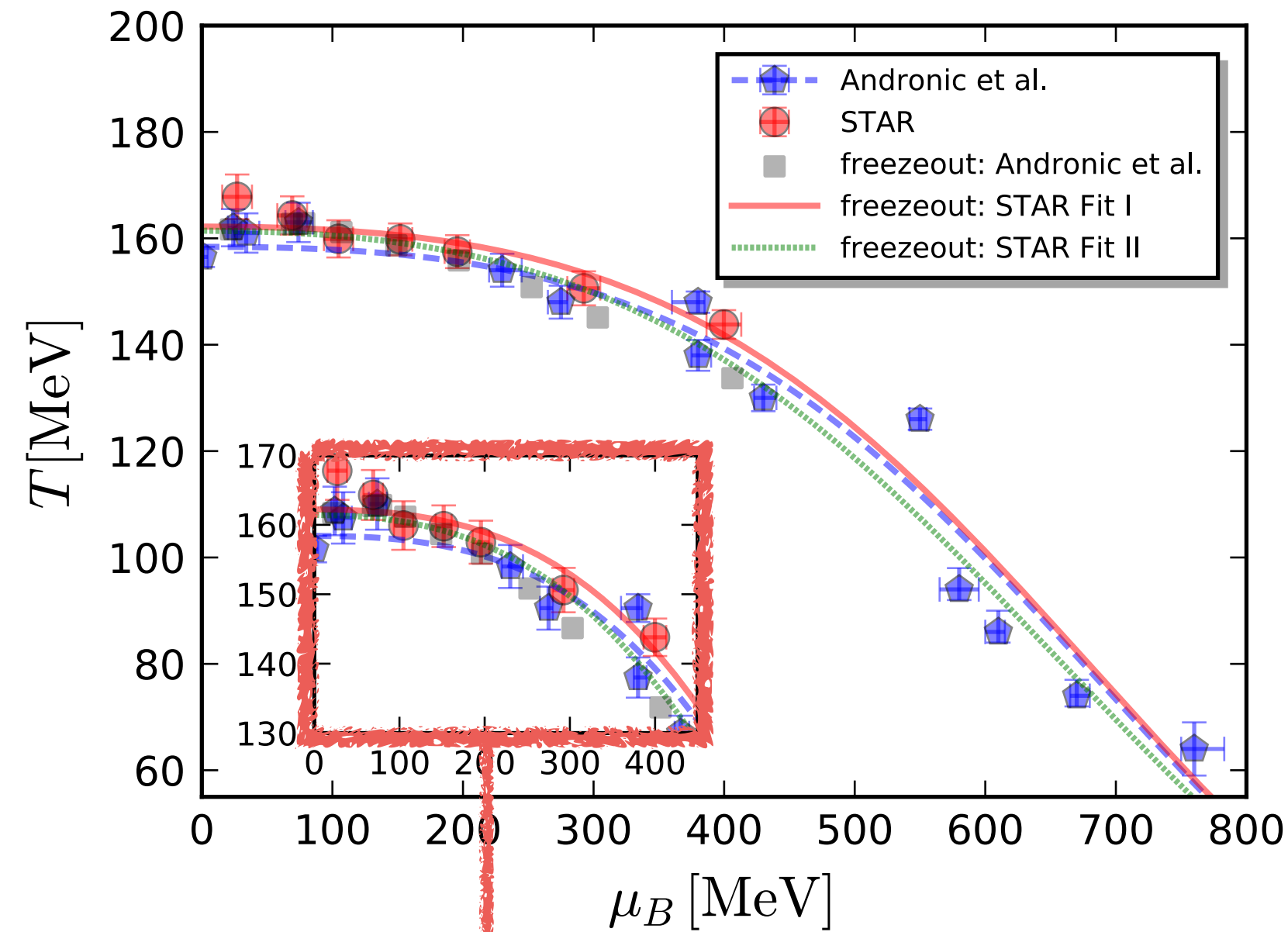
QCD-assisted LEFT

Finite density



Fu, Luo, Pawłowski, Rennecke, Wen, SY, *PRD* 104 (2021) ,094047

Freeze-out curves



three freeze-out curves

1. freeze-out: Andronic *et al.*

Andronic, Braun-Munzinger, Redlich,
Nature 561 (2018) 7723, 321

2. freeze-out: STAR Fit I

L. Adamczyk *et al.* (STAR), *PRC* 96 (2017), 044904

3. freeze-out: STAR Fit II

Fu, Luo, Pawłowski, Rennecke, Wen, SY, *PRD* 104 (2021) ,094047

$$\mu_{B_{CF}} = \frac{a}{1 + 0.288\sqrt{s_{NN}}},$$

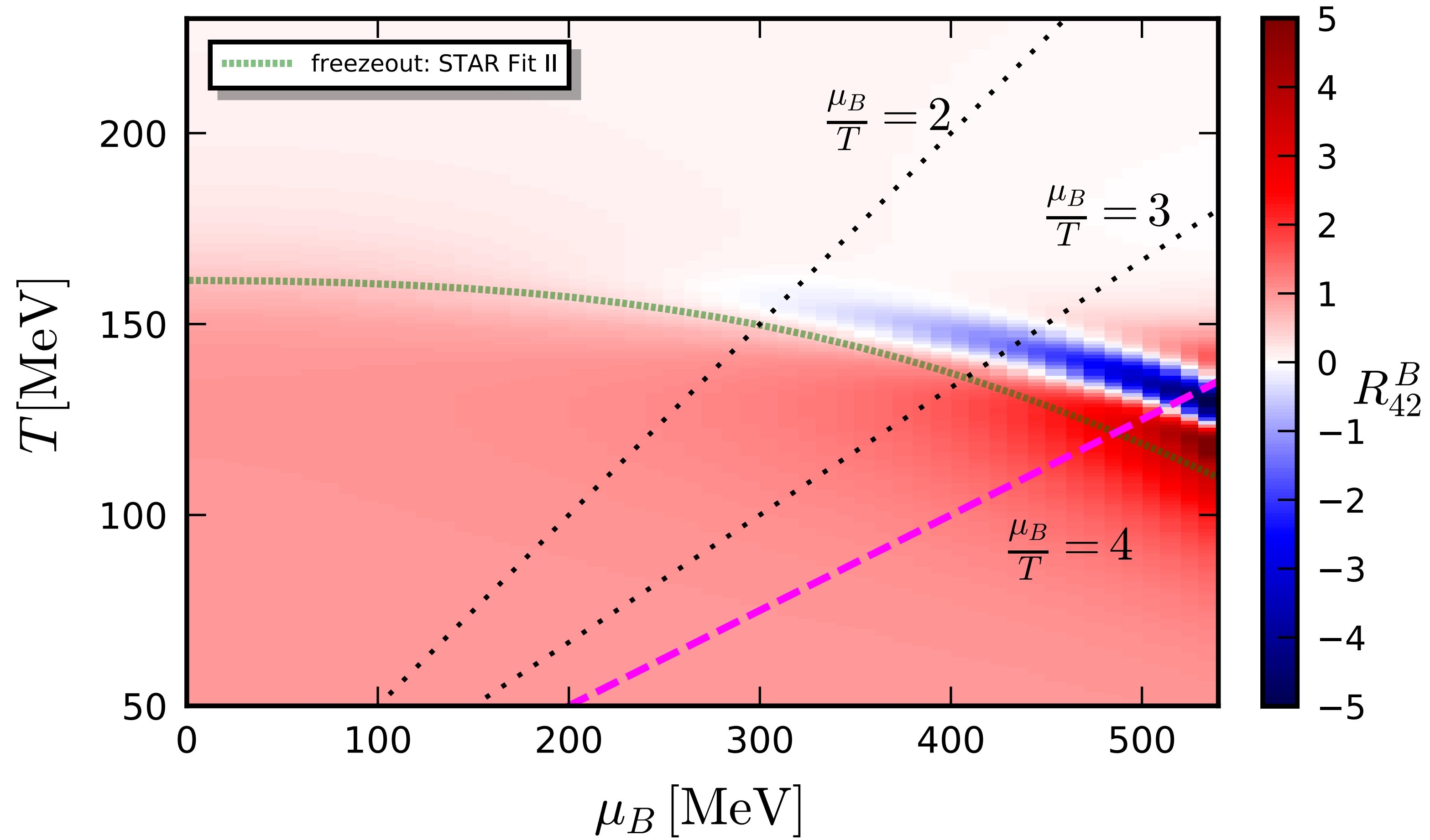
$$T_{CF} = \frac{T_{CF}^{(0)}}{1 + \exp(2.60 - \ln(\sqrt{s_{NN}})/0.45)}$$

all data points

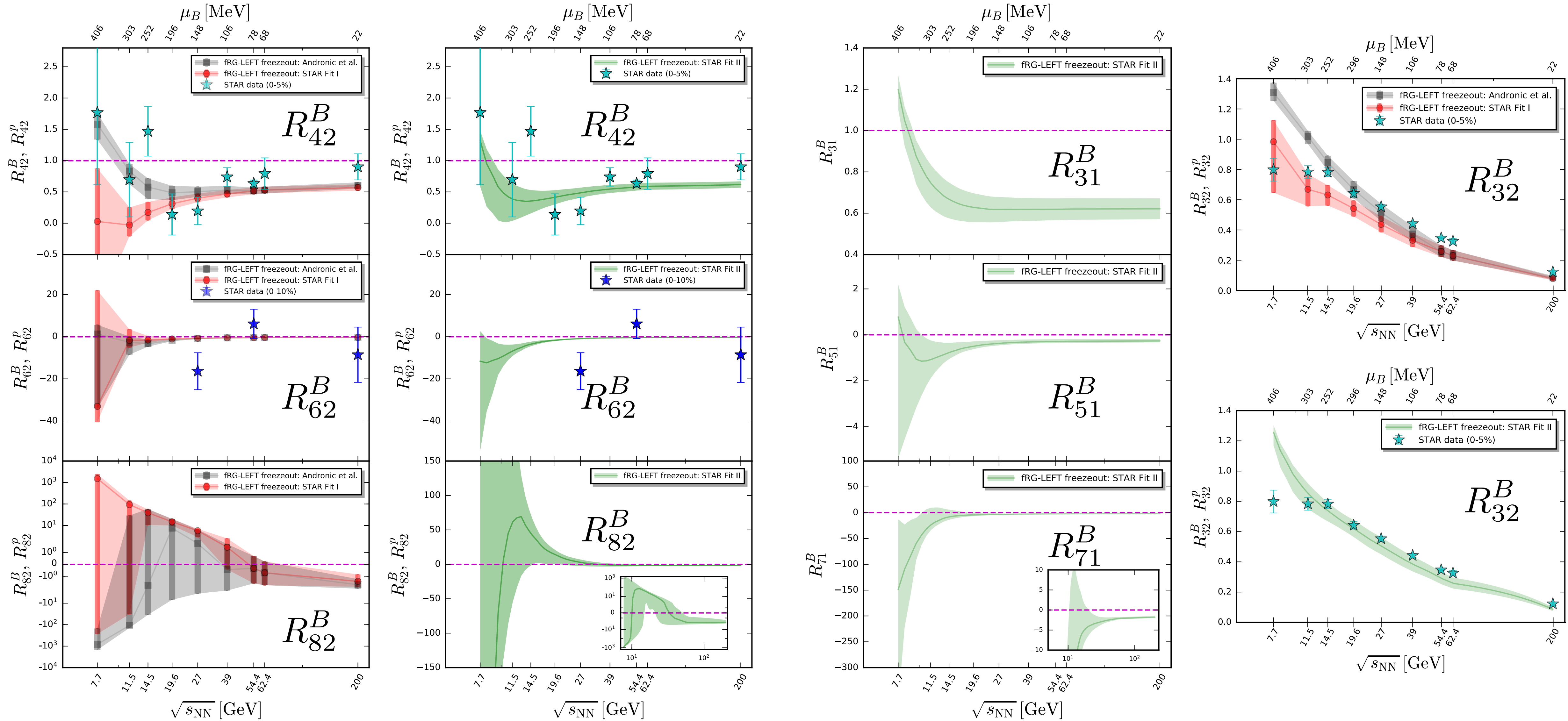
neglecting first two at low μ_B
and the last one

- freeze-out curve should not rise with μ_B
- convexity of the freeze-out curve

Freeze-out curves



Baryon number fluctuations on Freeze-out curves



Summary of our first step

1. We use a QCD-assisted LEFT to compute baryon number fluctuations.
2. Fluctuations have a non-monotonic energy dependence.
3. Non-monotonic behavior can arise with the sharper crossover.

Shortcoming

- A. Reliable only in low-density areas
- B. Only Grand Canonical Ensemble results

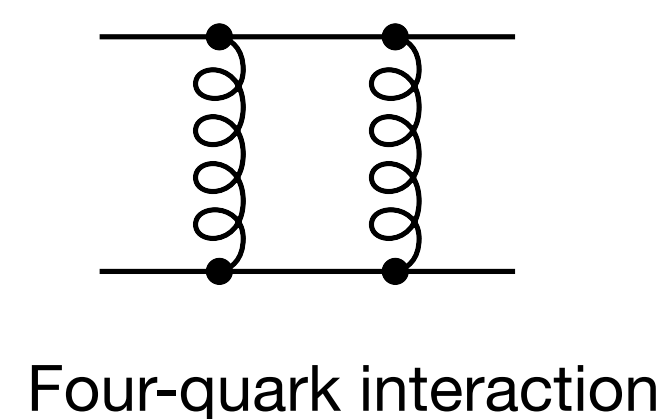
Looking for ripples

Second Step: Updated QCD-assisted LEFT

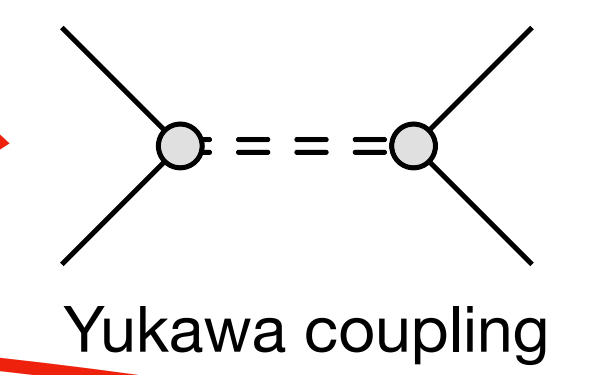
Call support from fRG-QCD

First Principle QCD flow equation:

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{[Orange loop]} - \text{[Dotted loop]} - \text{[Black loop]} + \frac{1}{2} \text{[Blue loop]}$$



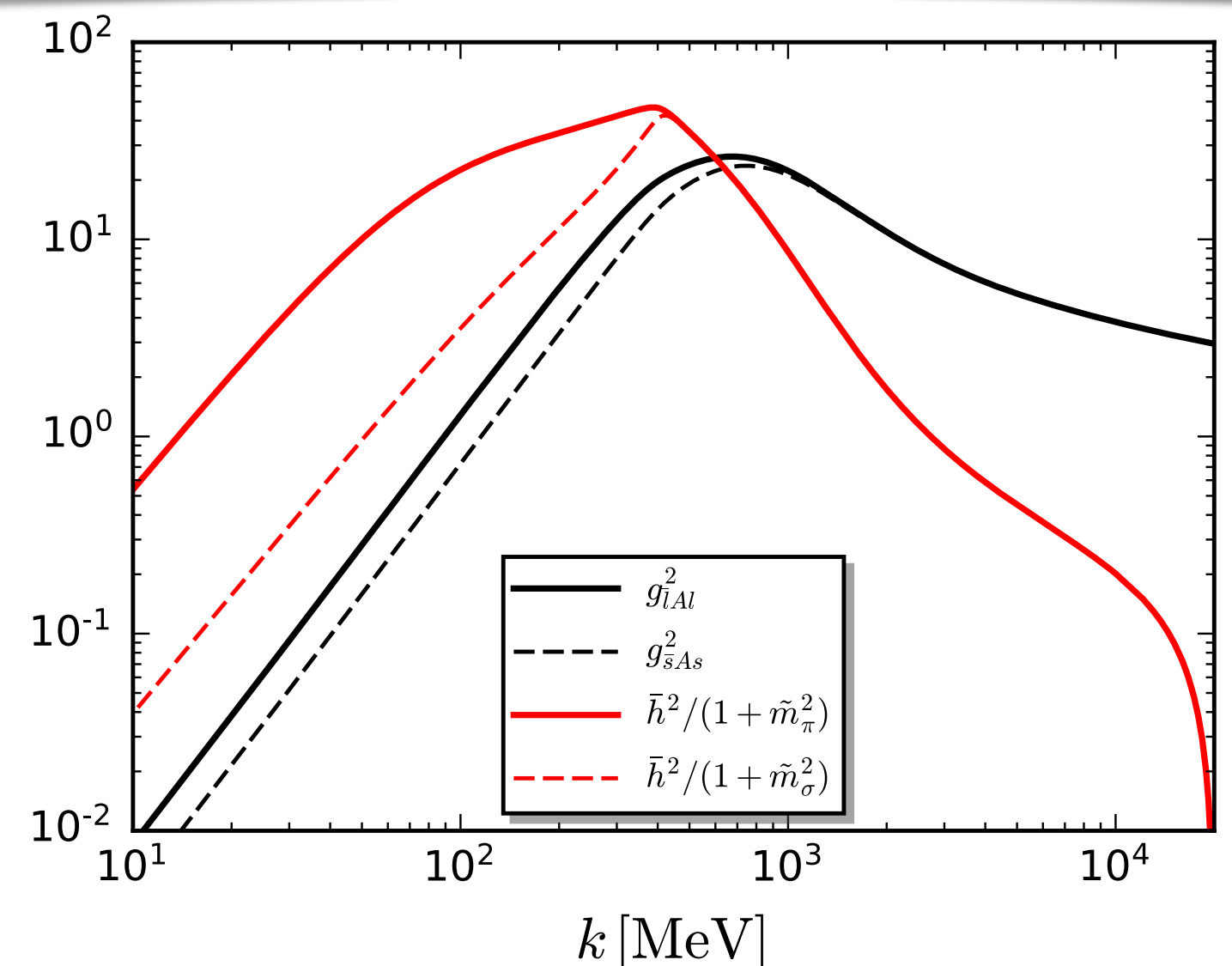
Dynamical hadronization



$$\partial_t \bar{h} = \left(\frac{1}{2} \eta_\phi + \eta_q \right) \bar{h} - \bar{m}_\pi^2 \dot{A} + \overline{\text{Flow}}_{(\bar{q}\tau q)\bar{\pi}}^{(3)}$$

LEFT flow equation:

$$\partial_t \Gamma_k[\Phi] = - \text{[Black loop]} + \frac{1}{2} \text{[Blue loop]}$$

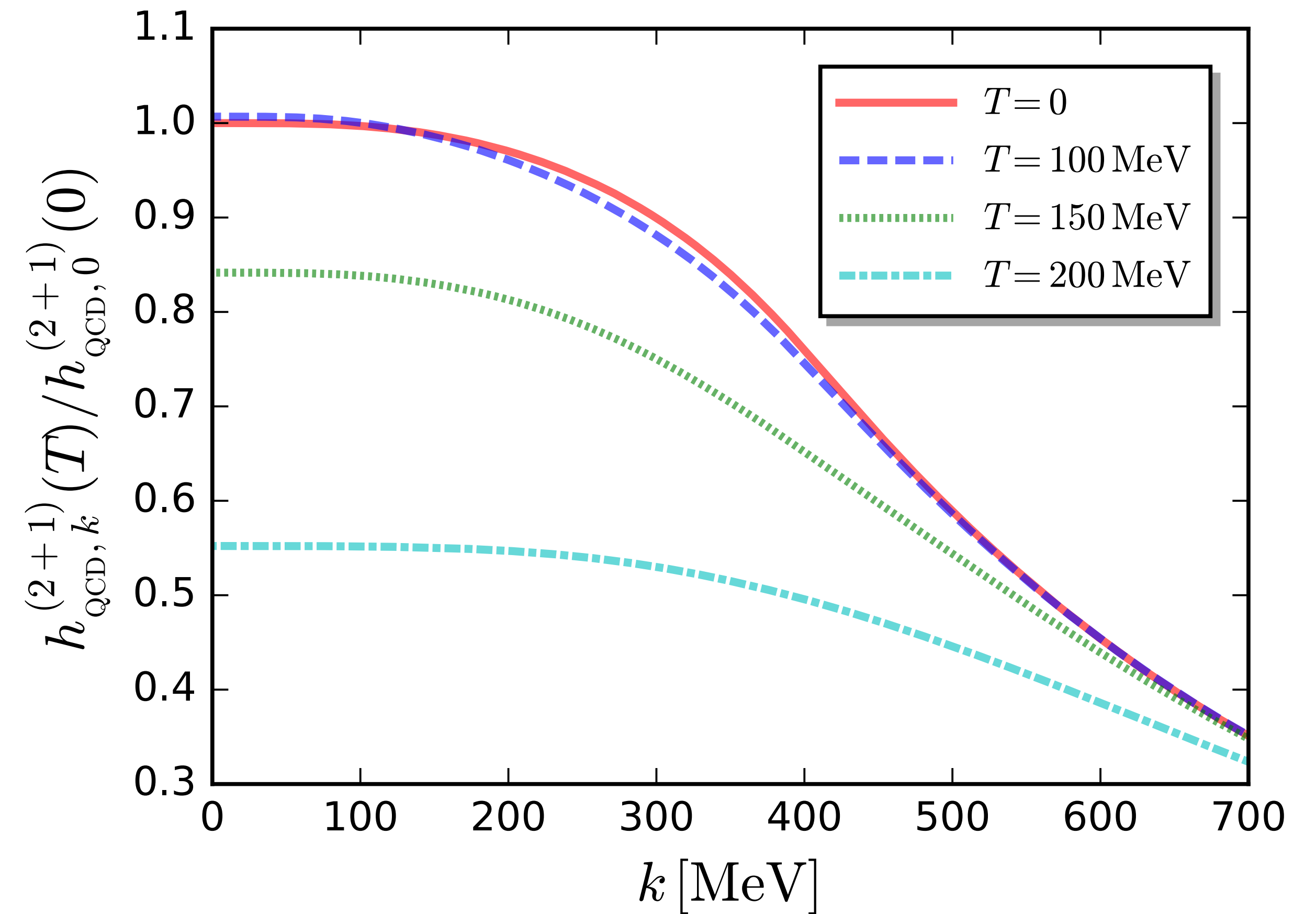
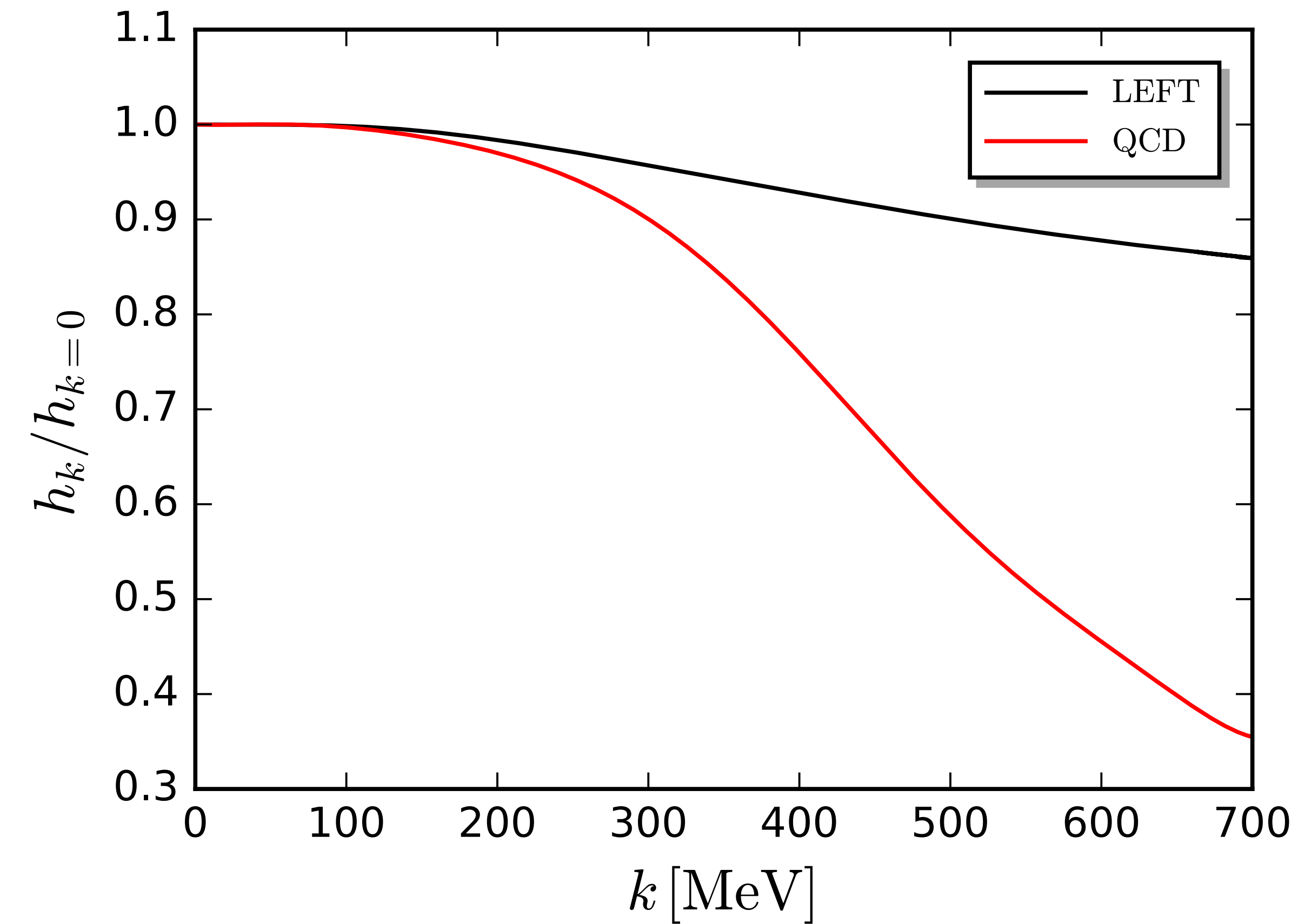


Exchange couplings

- * fRG-LEFT have been improved.
- * Dynamics of first-principle fRG-QCD has been encoded in the LEFT via the Yukawa coupling.

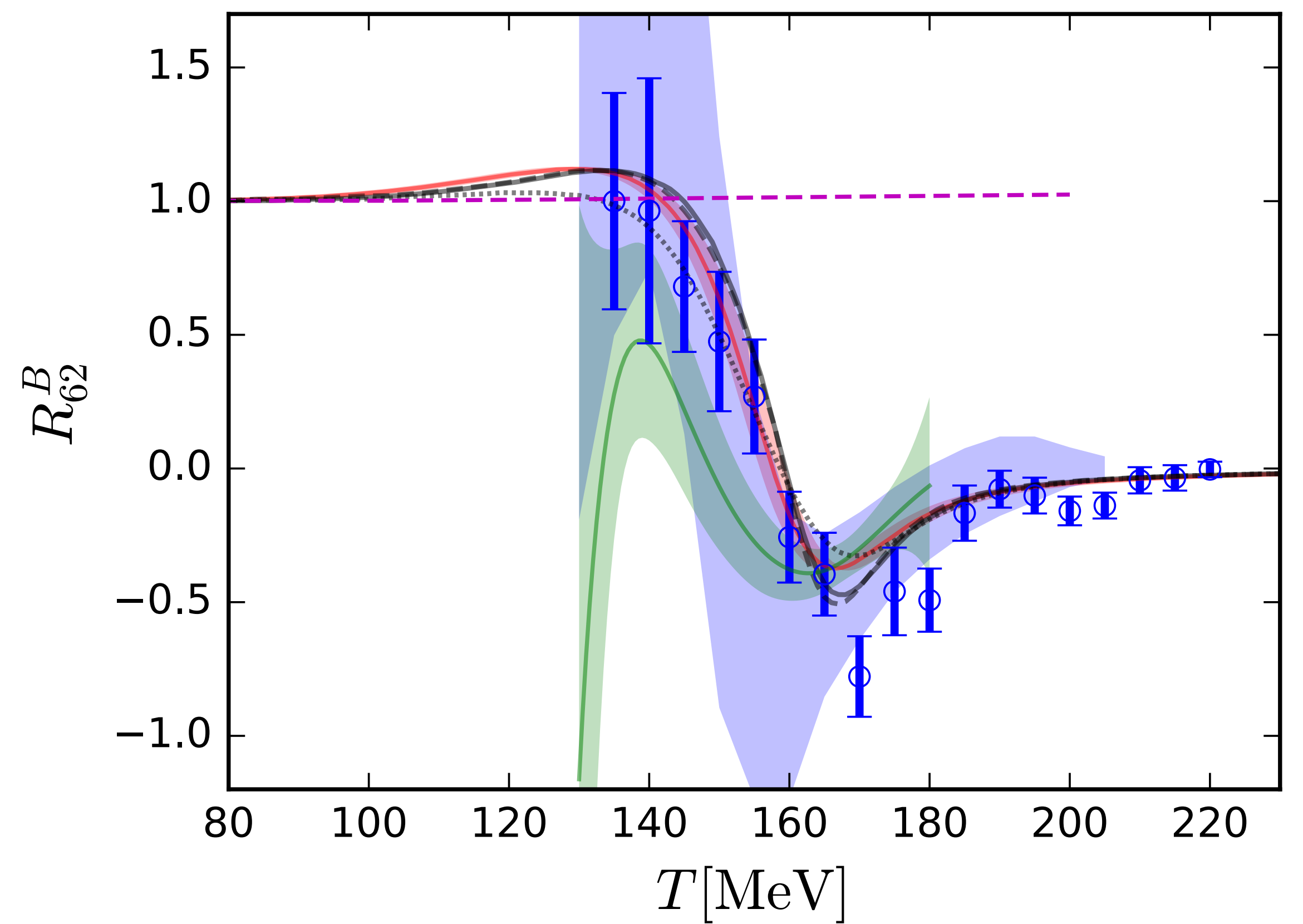
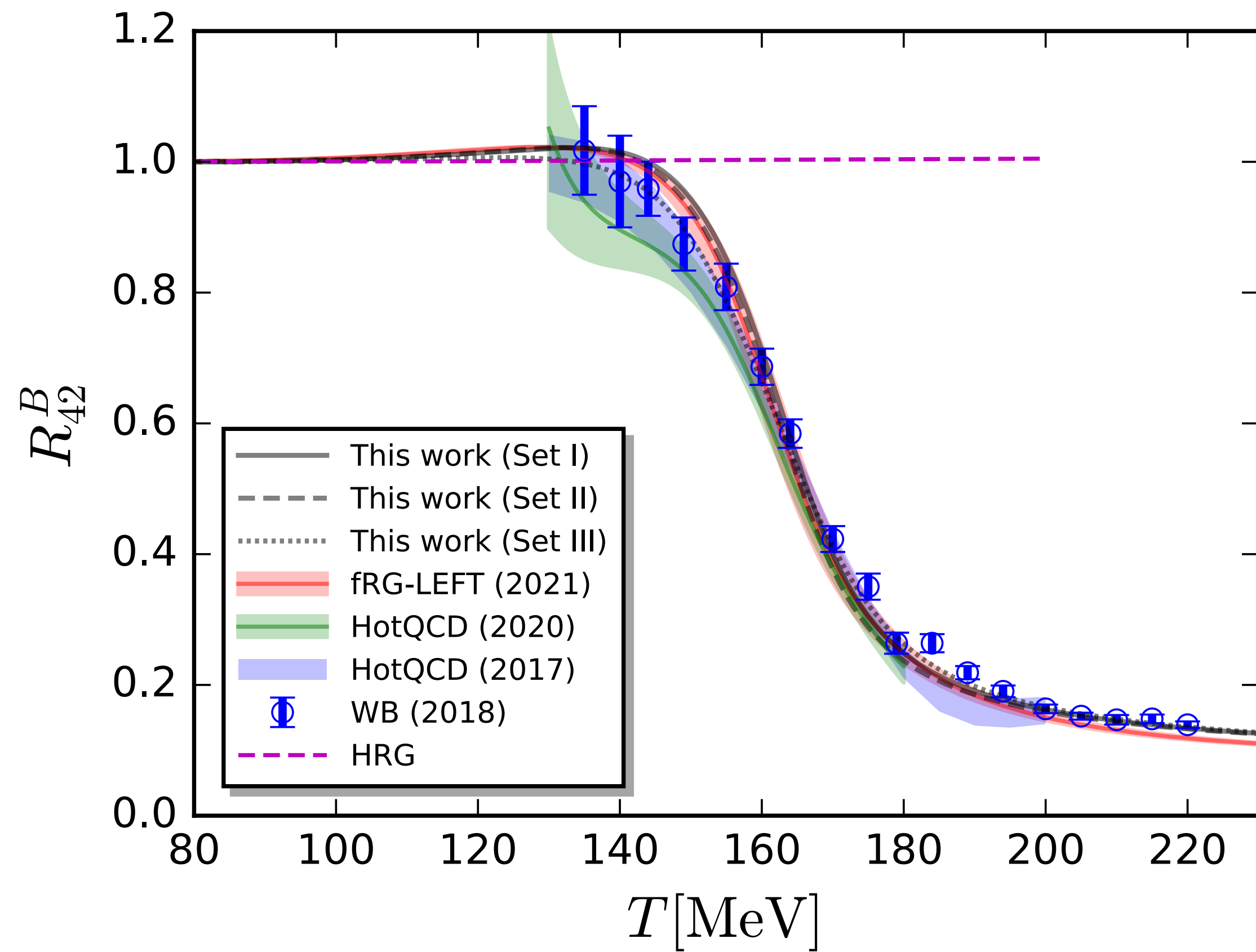
Updated QCD-assisted LEFT

Yukawa coupling

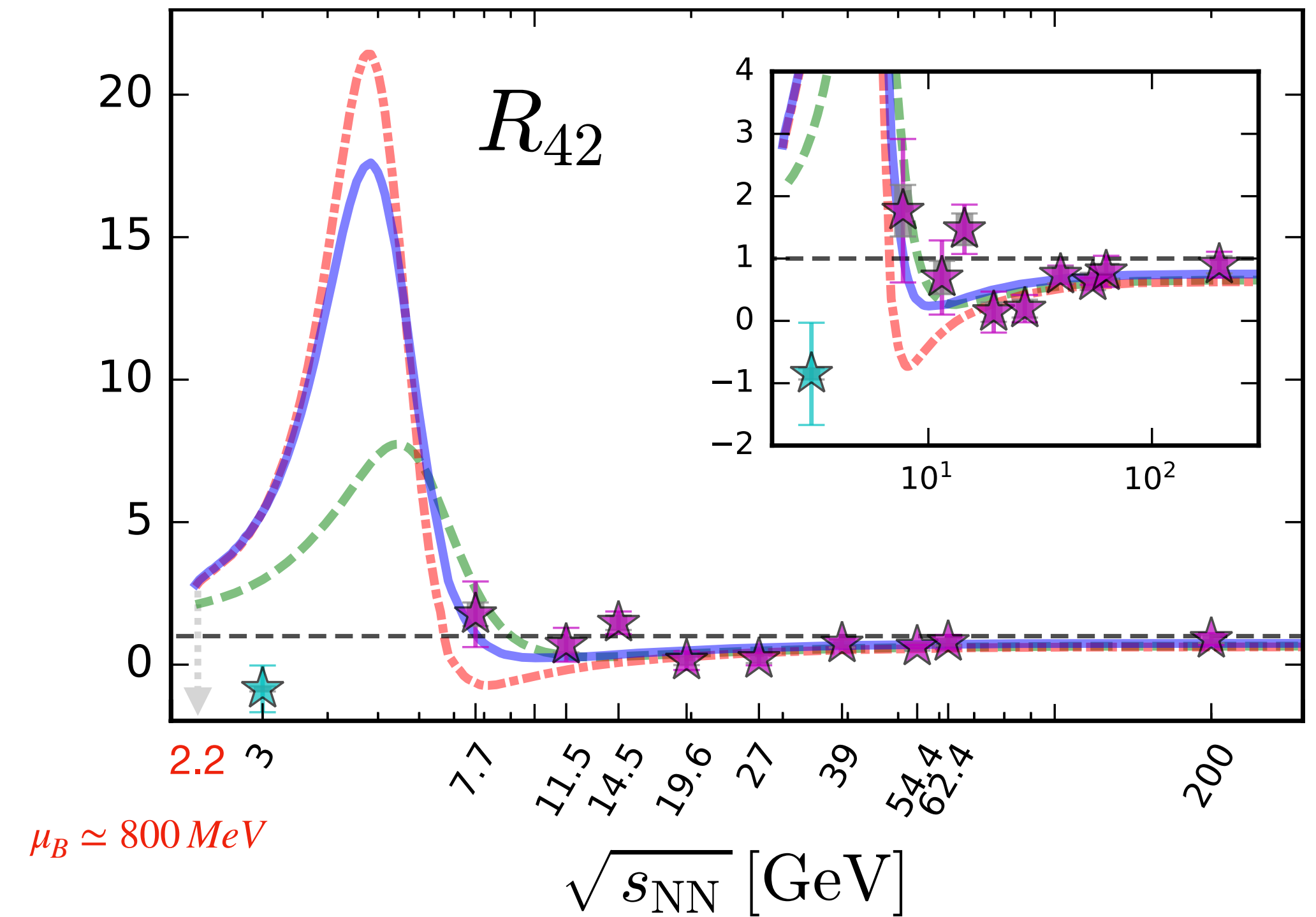
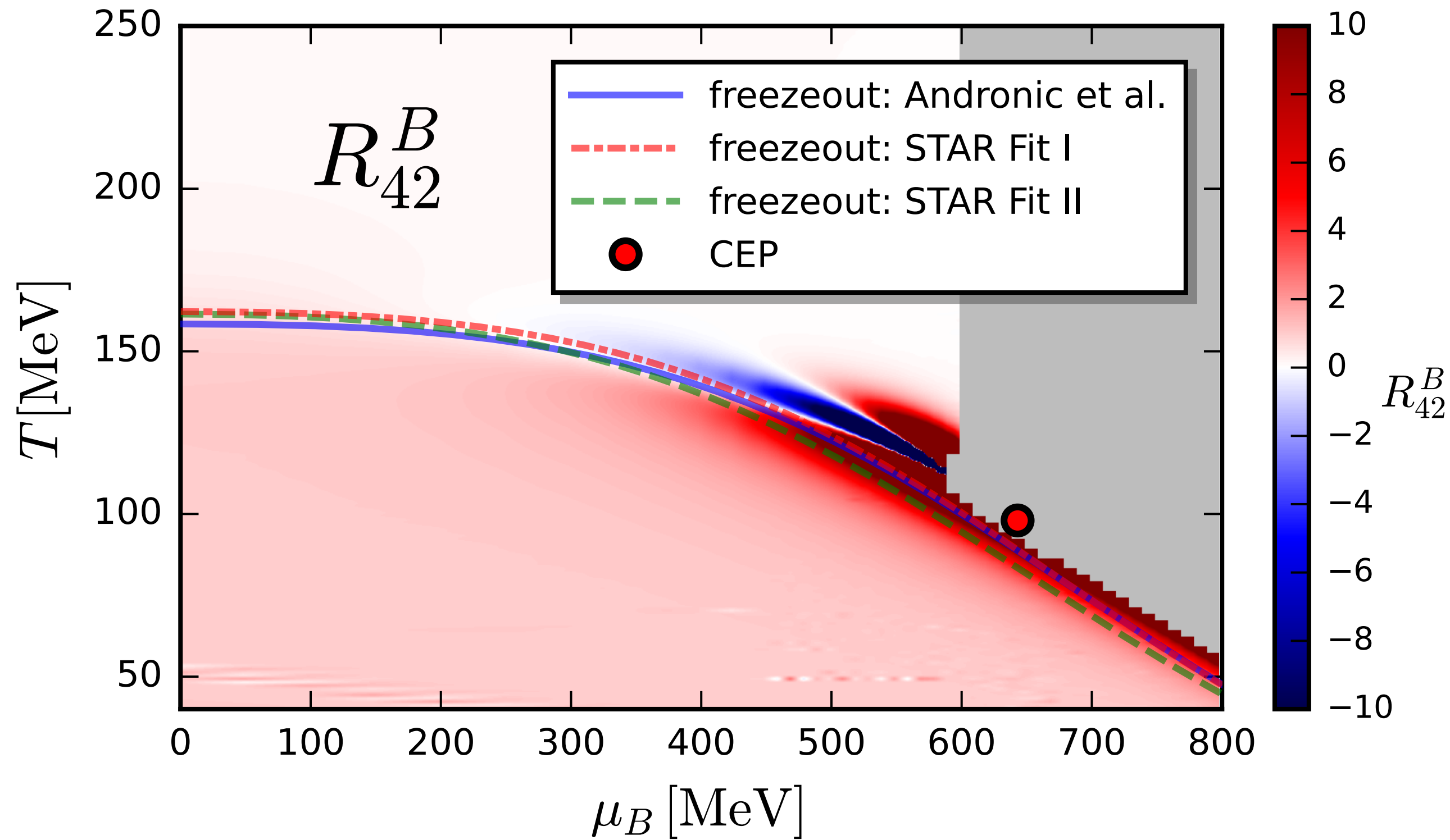


Updated QCD-assisted LEFT

$$\mu_B = 0$$



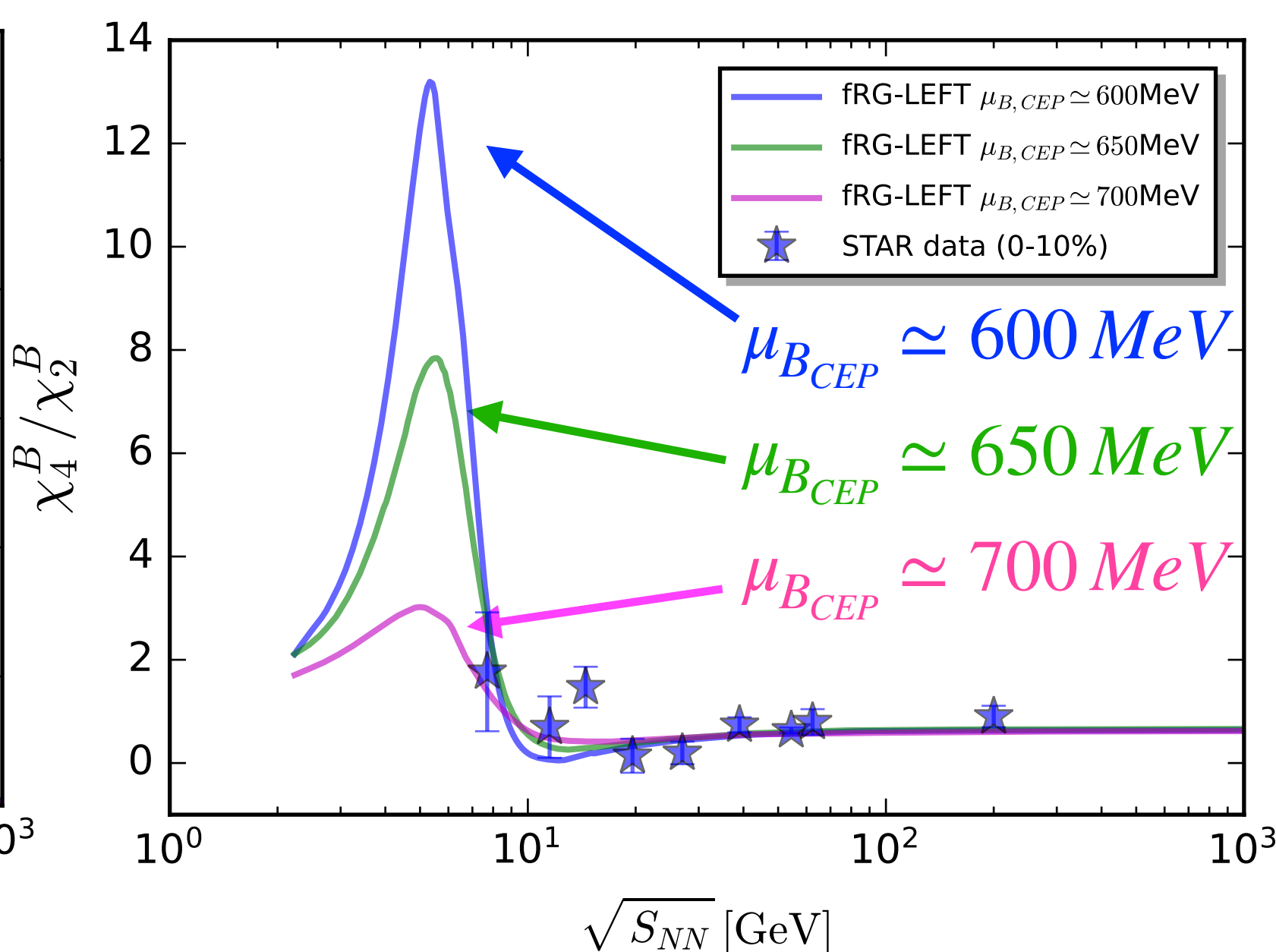
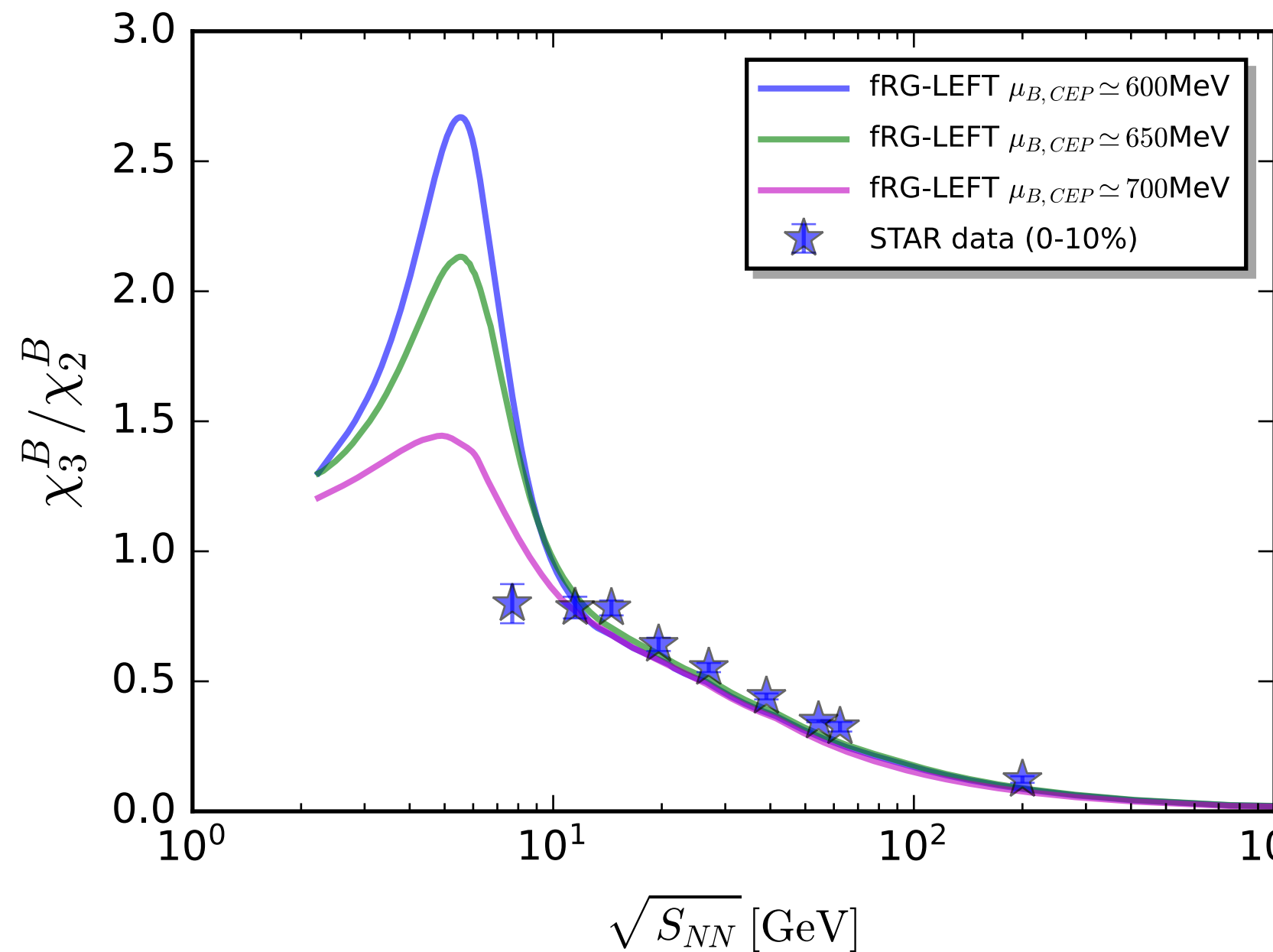
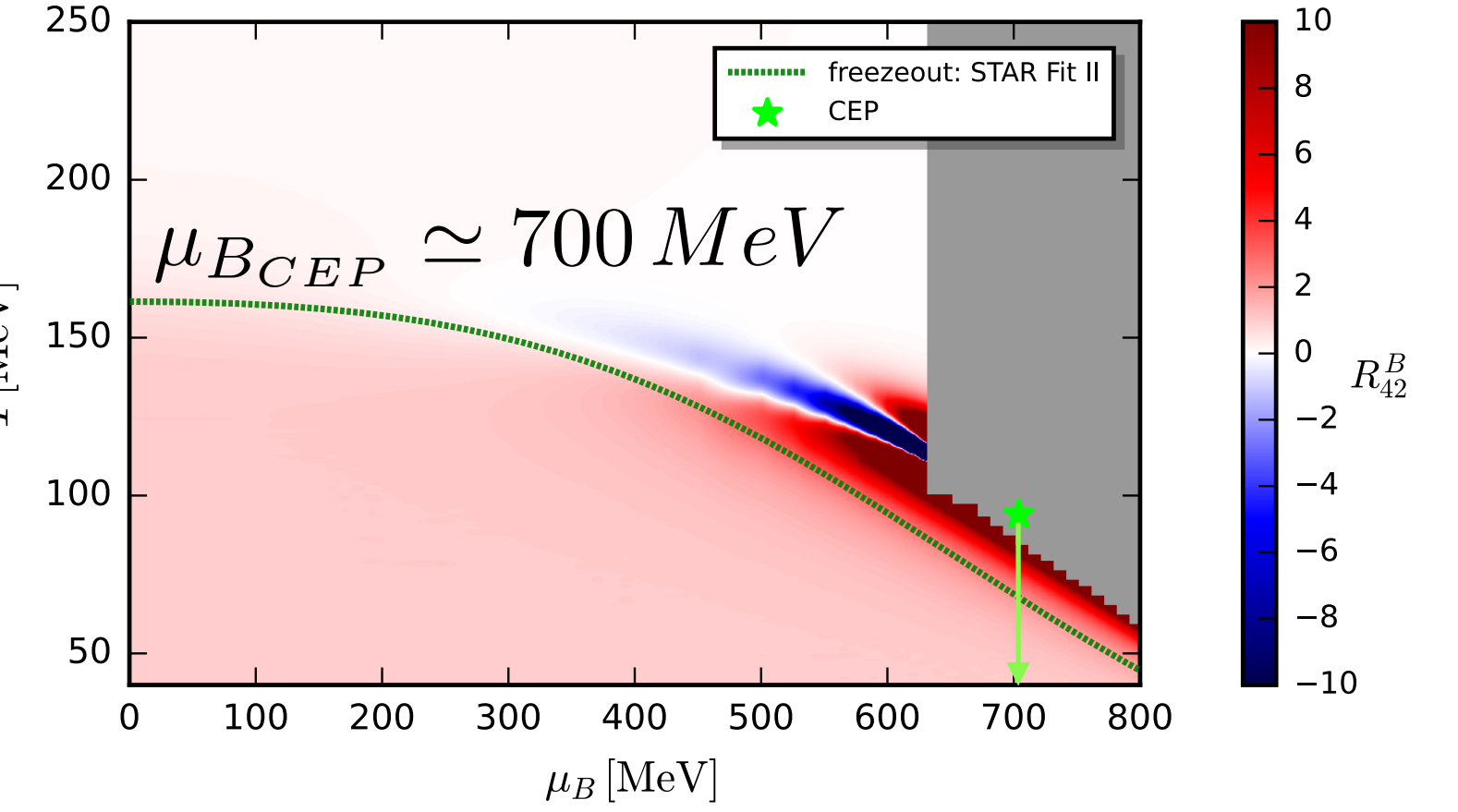
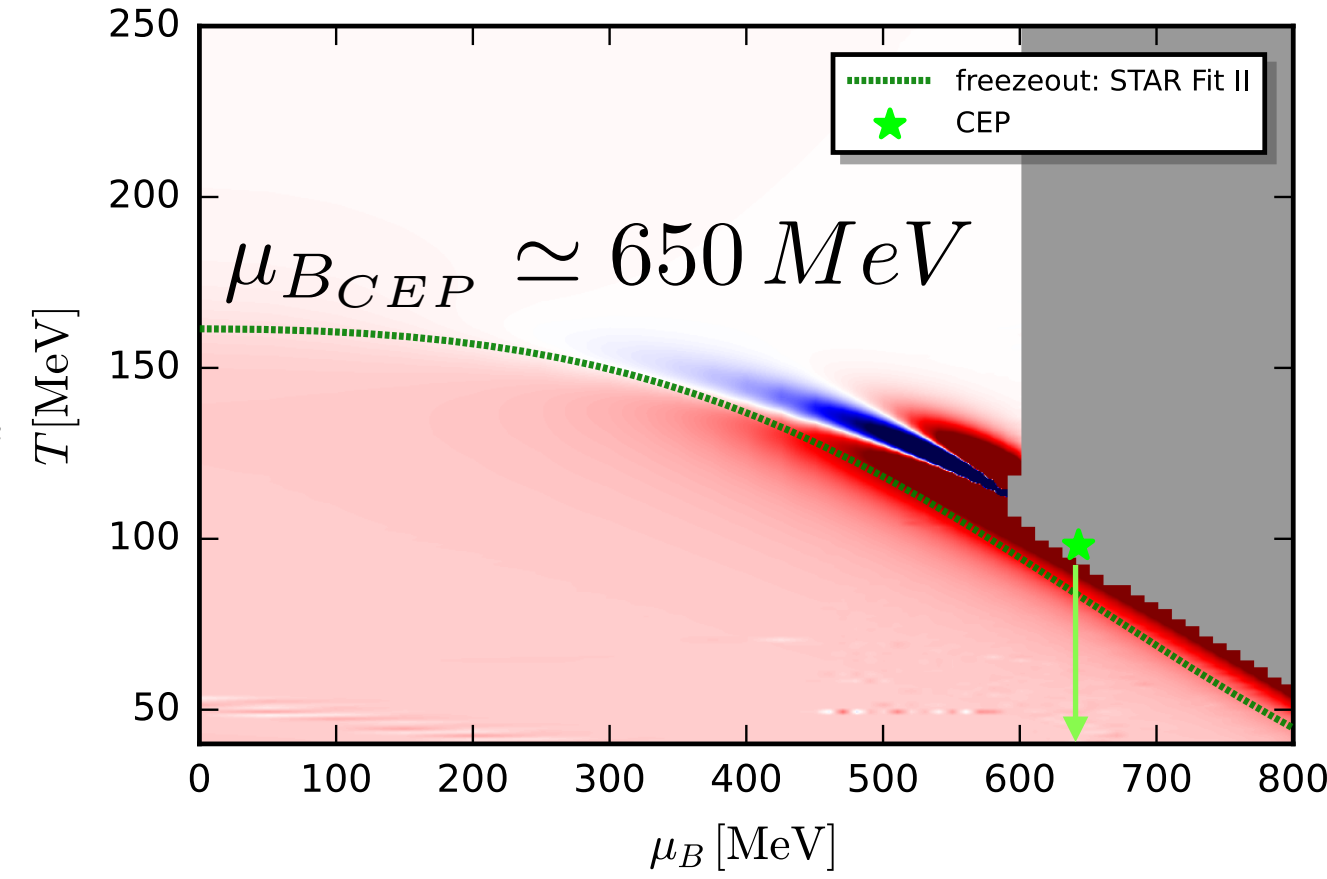
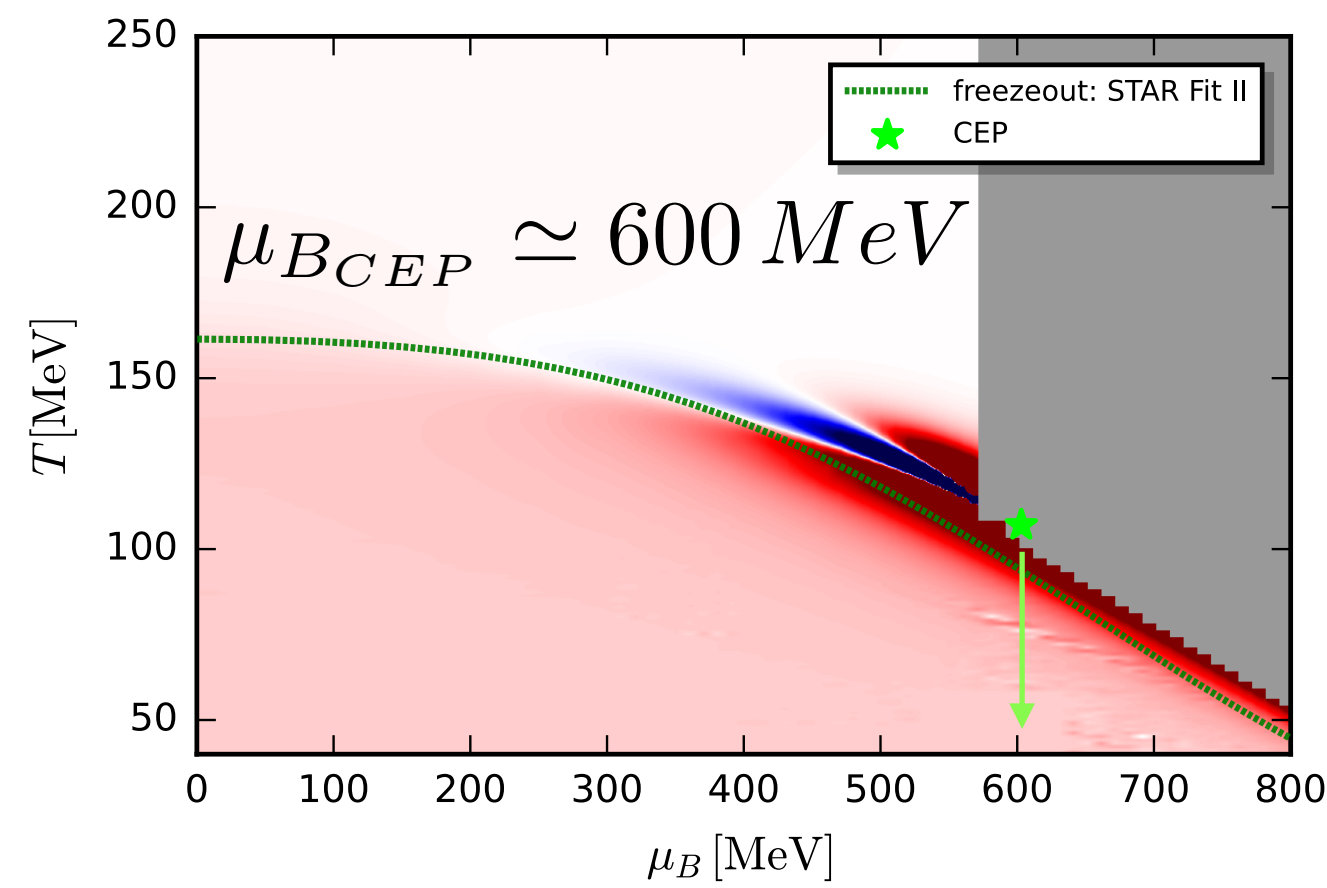
Fluctuations Phase diagram



★ Calculations have been extended from $\mu_B \sim 500$ to 800 MeV in the improved fRG-LEFT.

★ A “peak” structure is found in the regime of low collision energy.

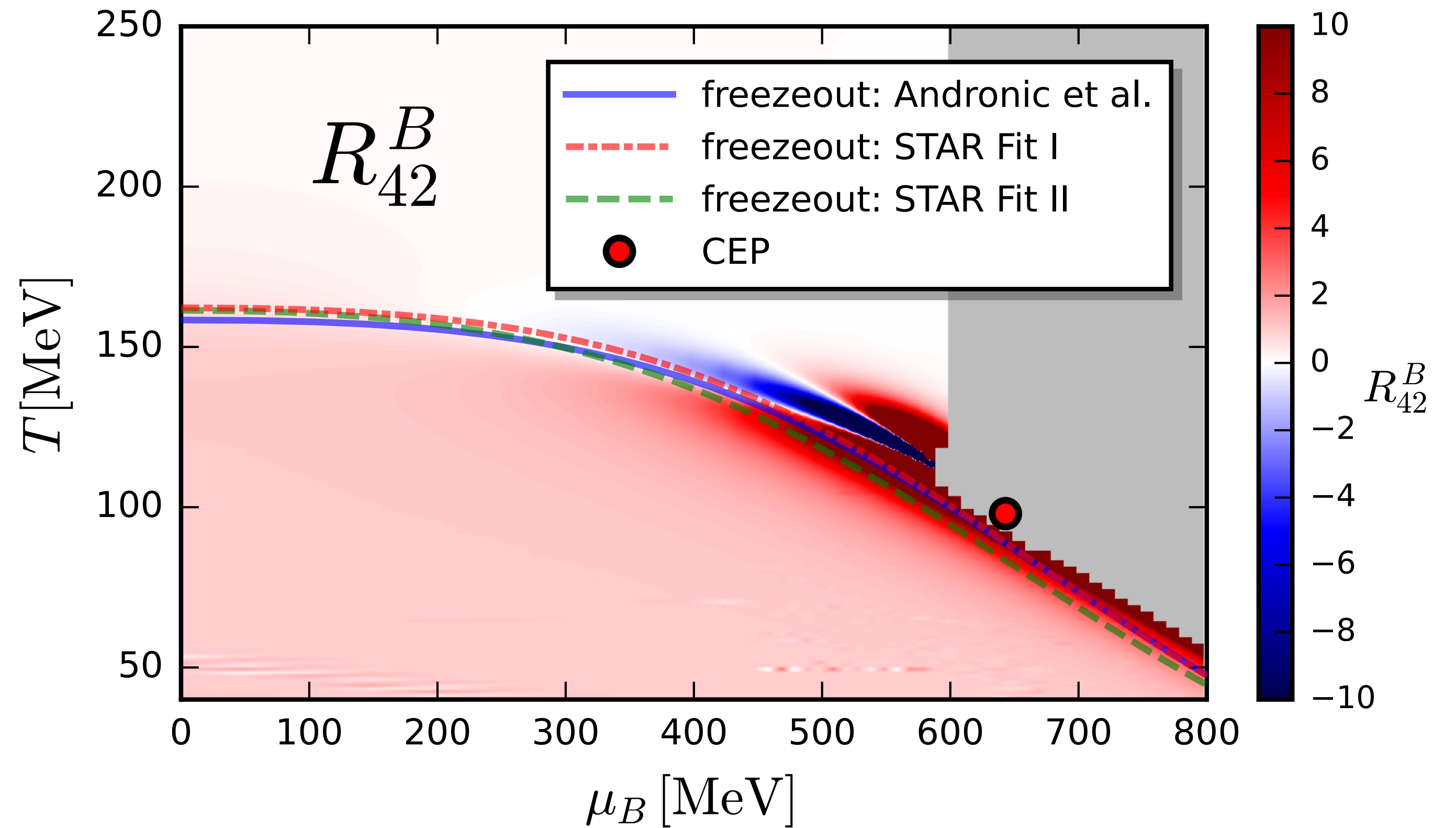
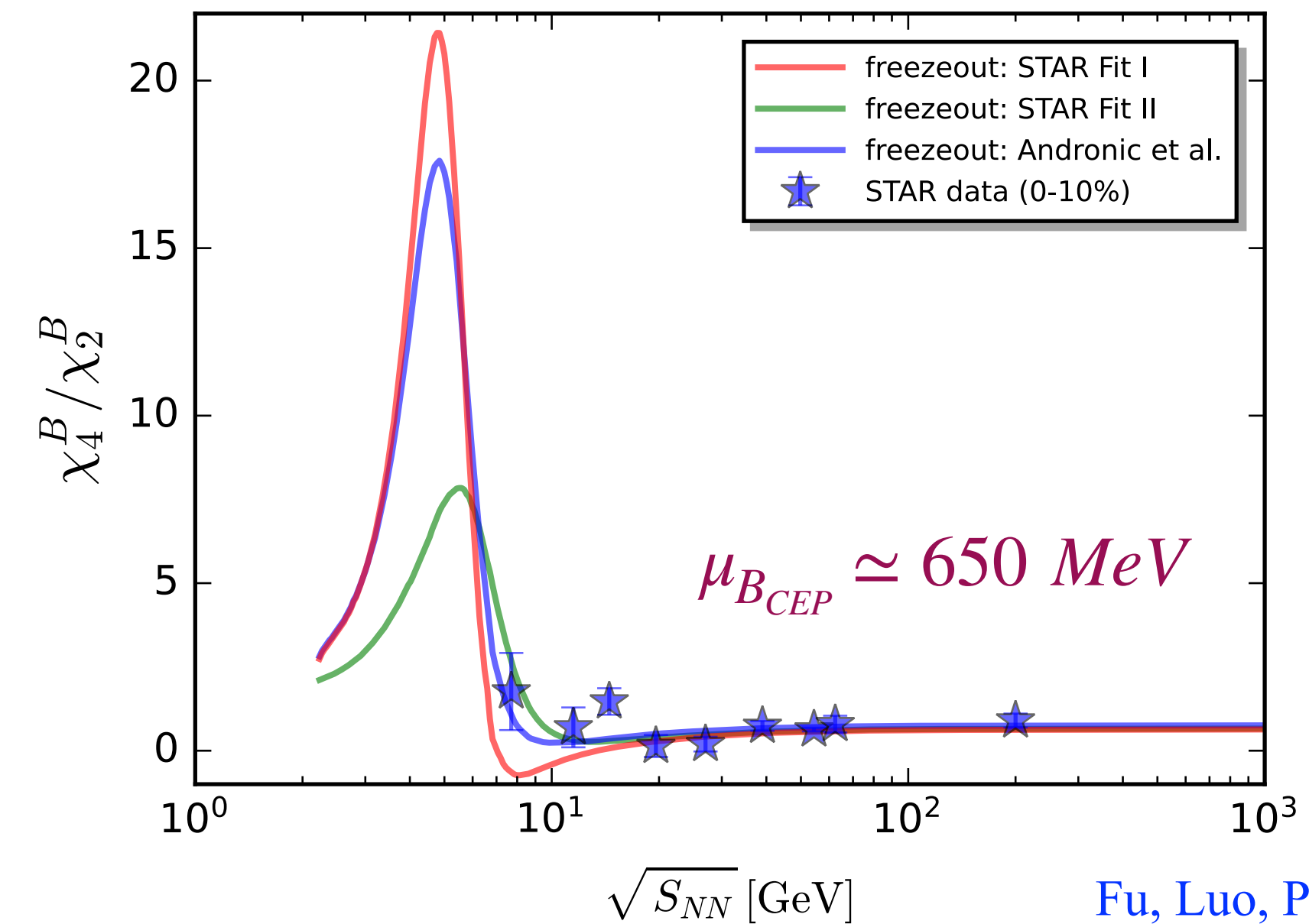
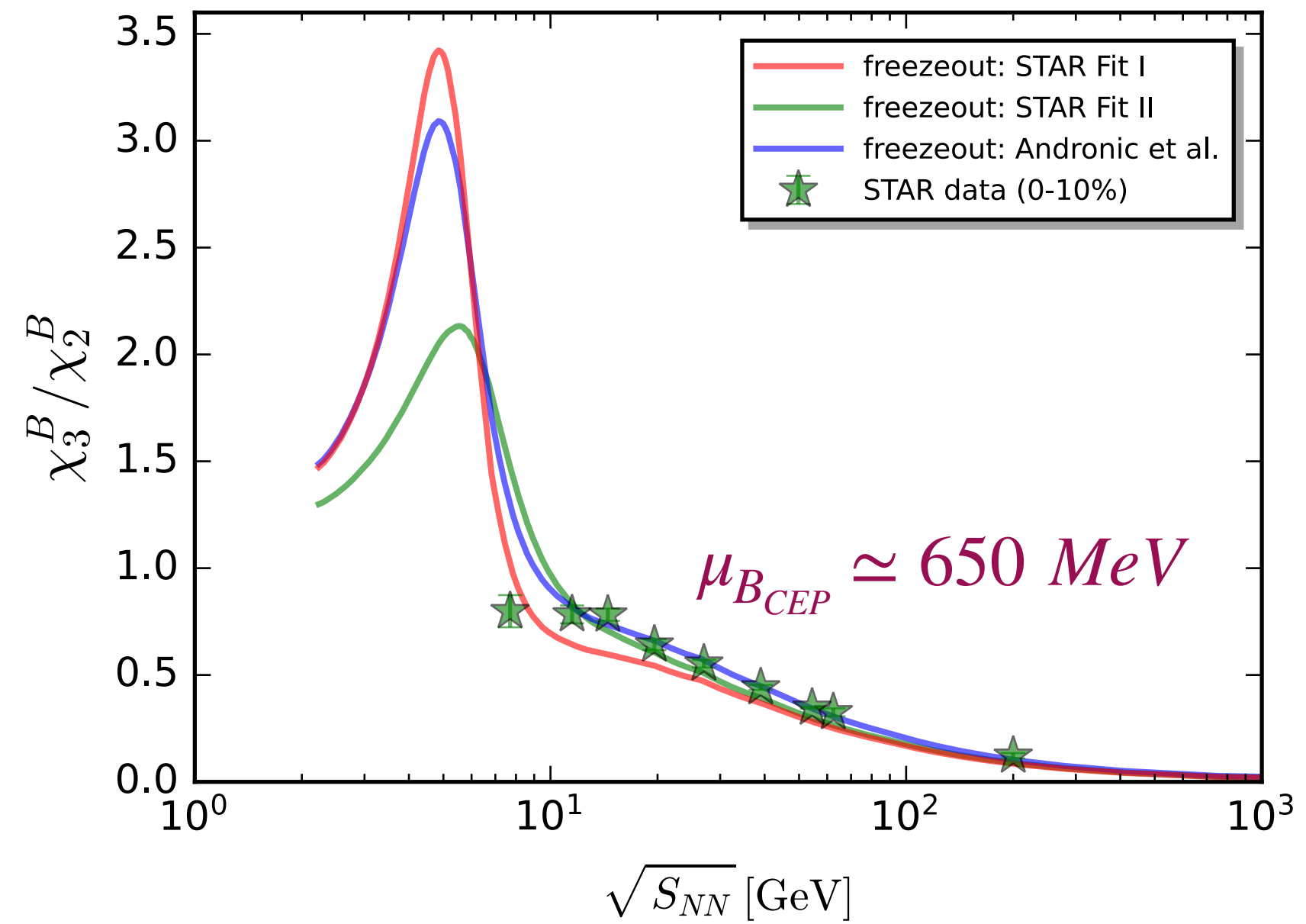
Different location of CEP



★ **Location** of the peak, i.e., the corresponding collision energy shows very **mild** dependence on the location of CEP in the range $\mu_{B,CEP} \simeq (600 \sim 700)$ MeV.

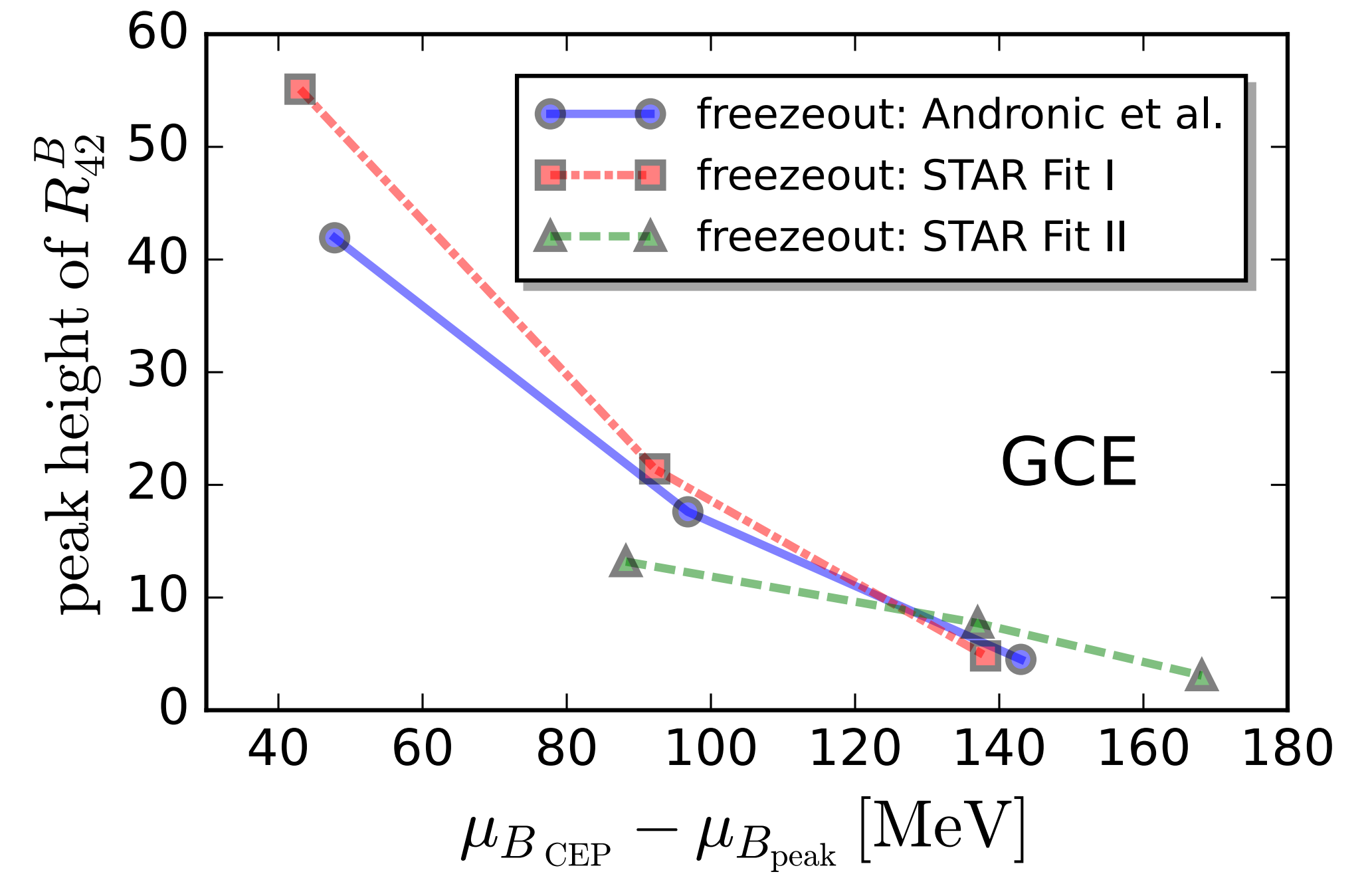
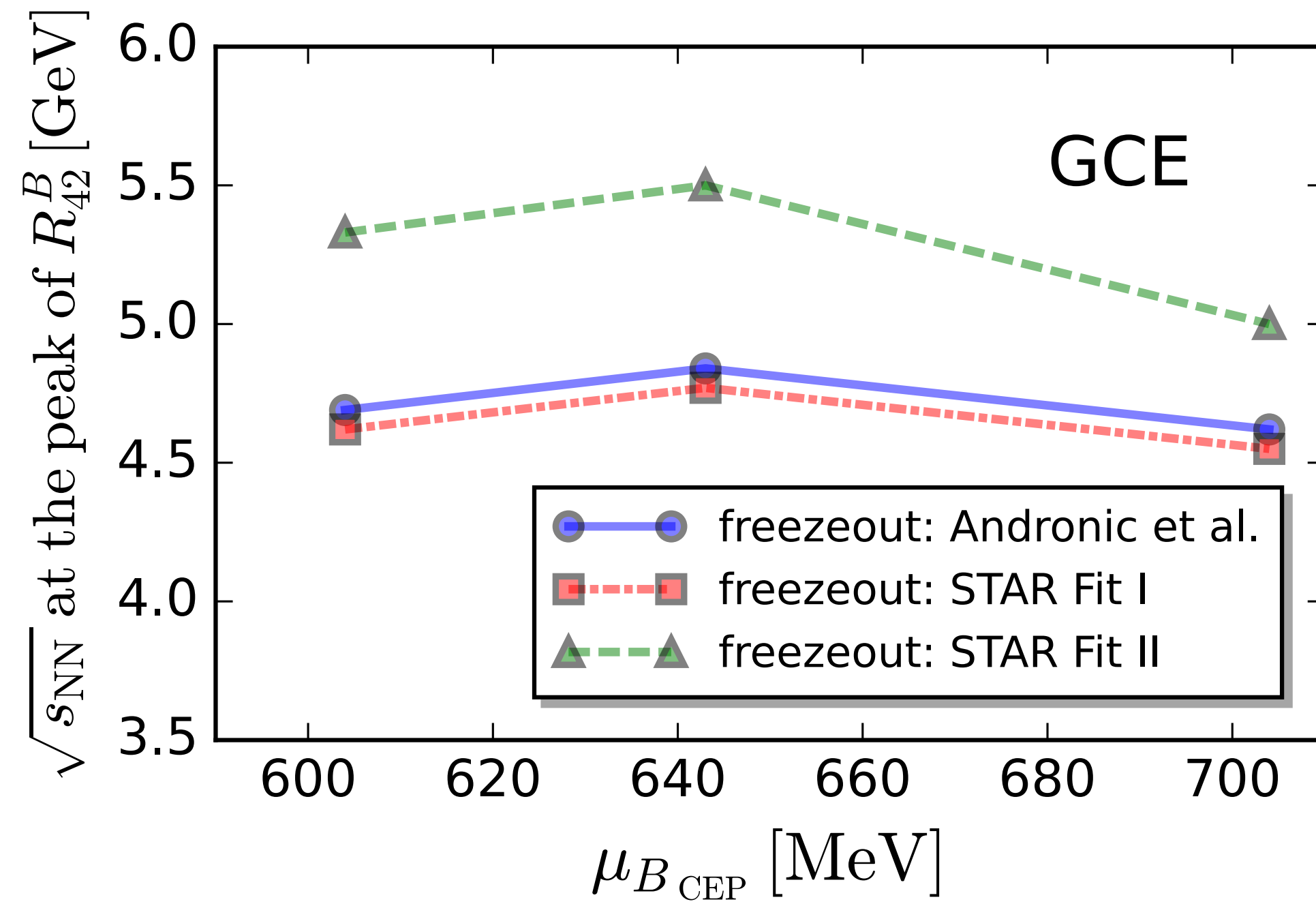
★ While the **height** of the peak shows **significant** dependence.

Different freeze-out curves

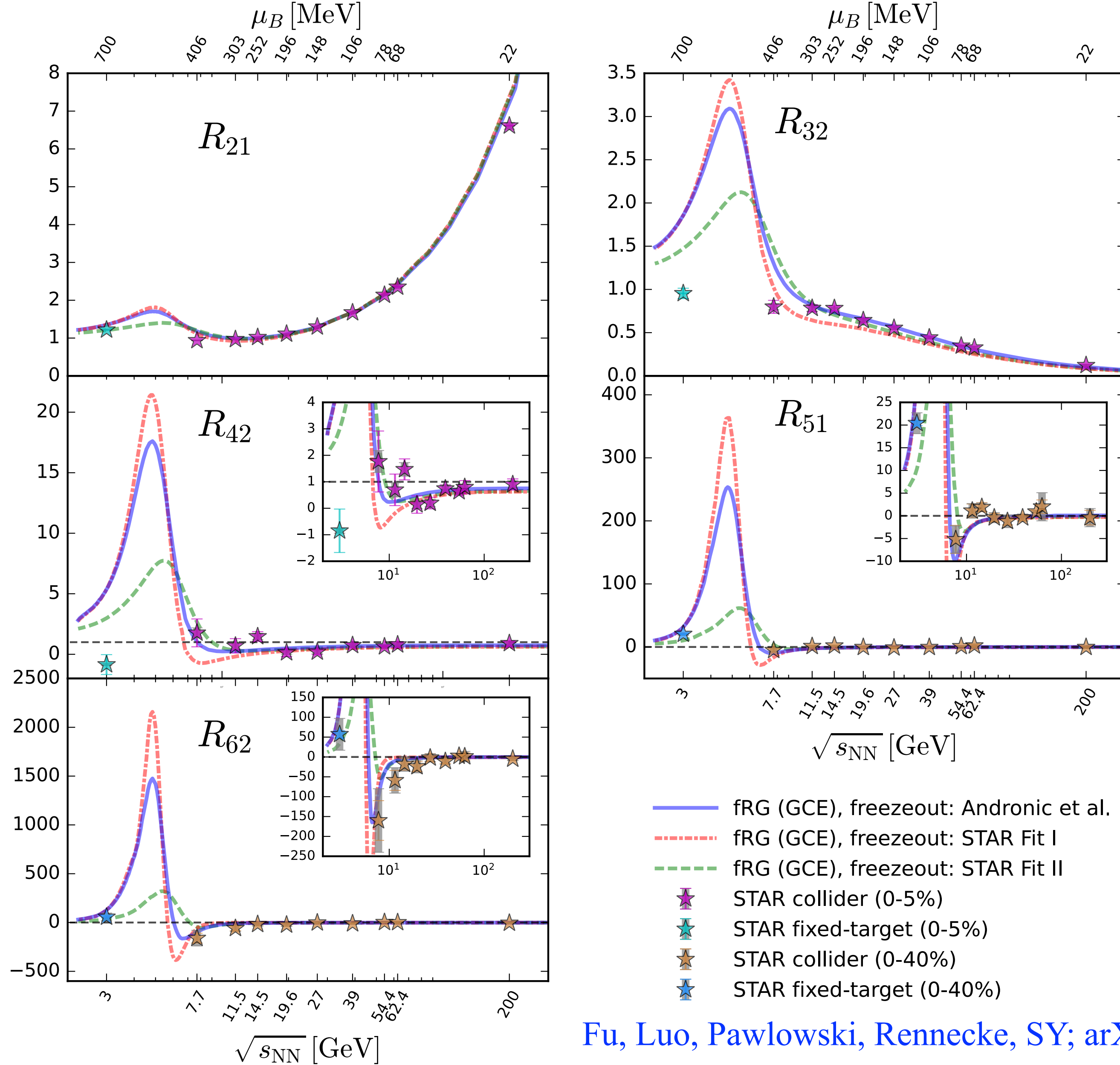


- ★ Height and location of the peak are influenced by different freeze-out curves, especially for the height.
- ★ Closer the freeze-out curve is to the phase boundary, higher the peak.

Different location of CEP & freeze-out curves



Fluctuations within GCE



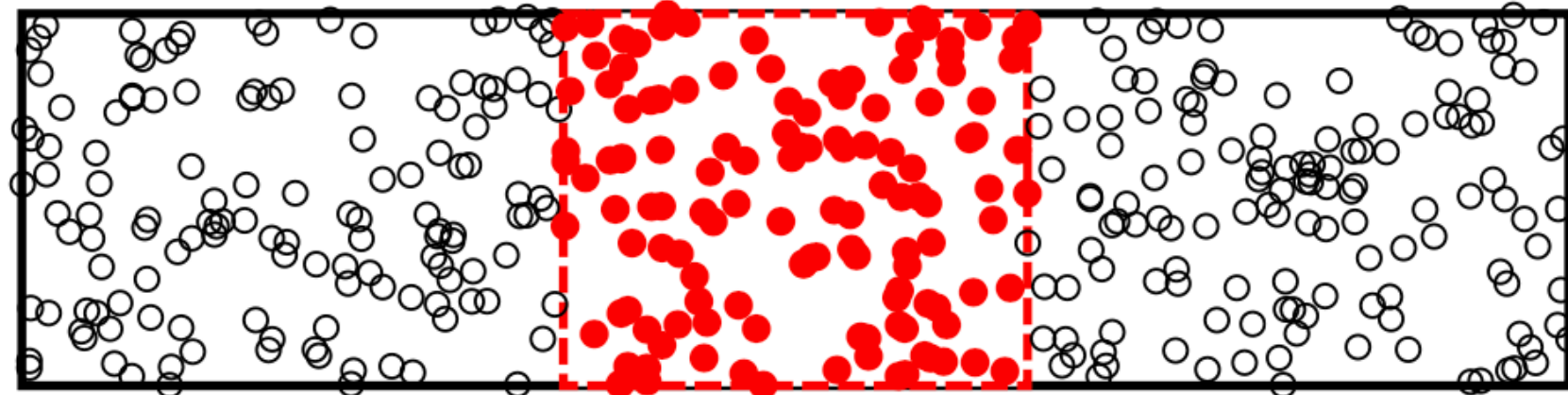
From GCE to CE

- Subensemble acceptance method

$V_1 = \alpha V$ α Proportion of subsystems to total system

$\beta = 1 - \alpha$ β Proportion of remaining systems

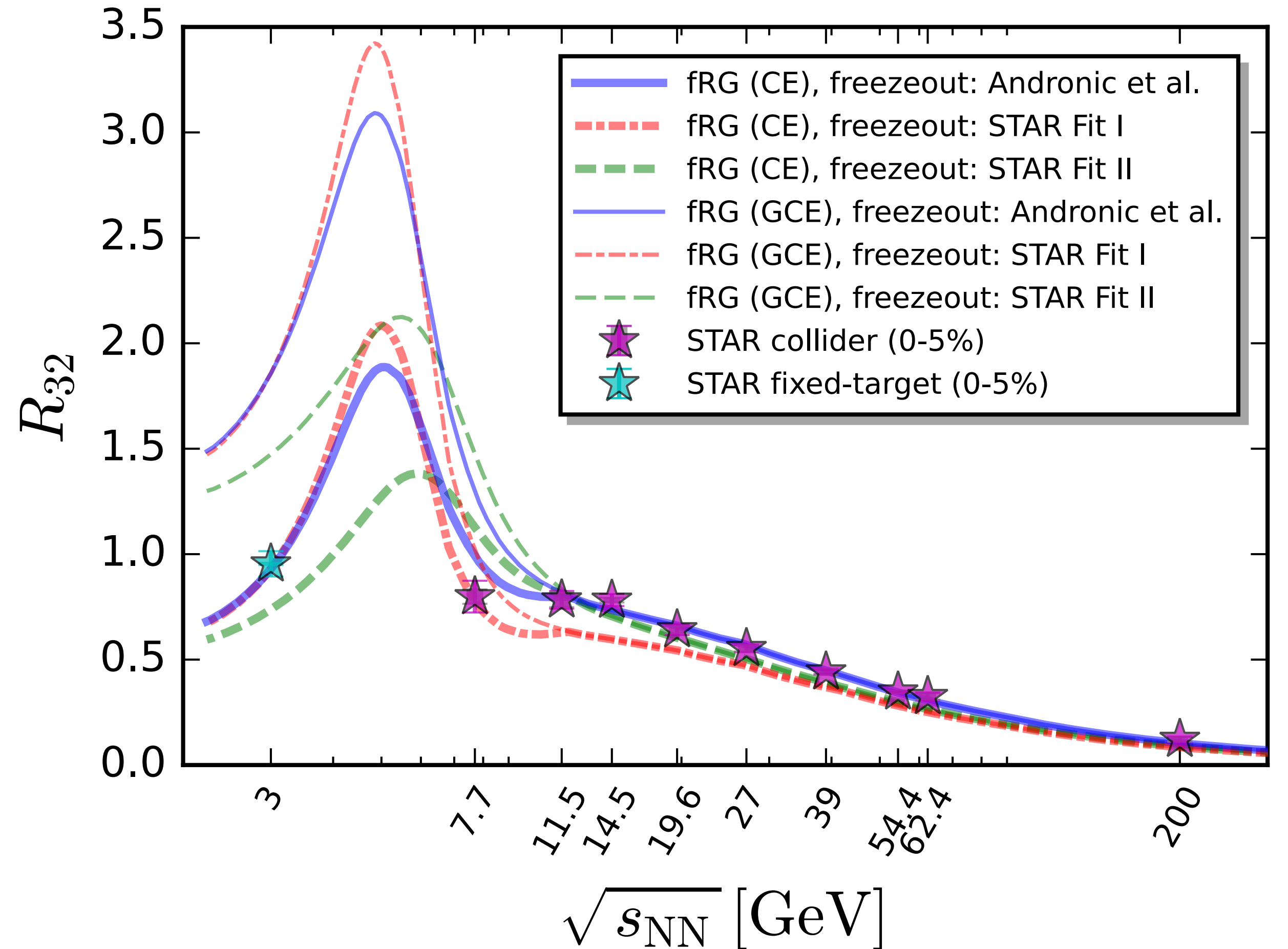
Subensemble acceptance



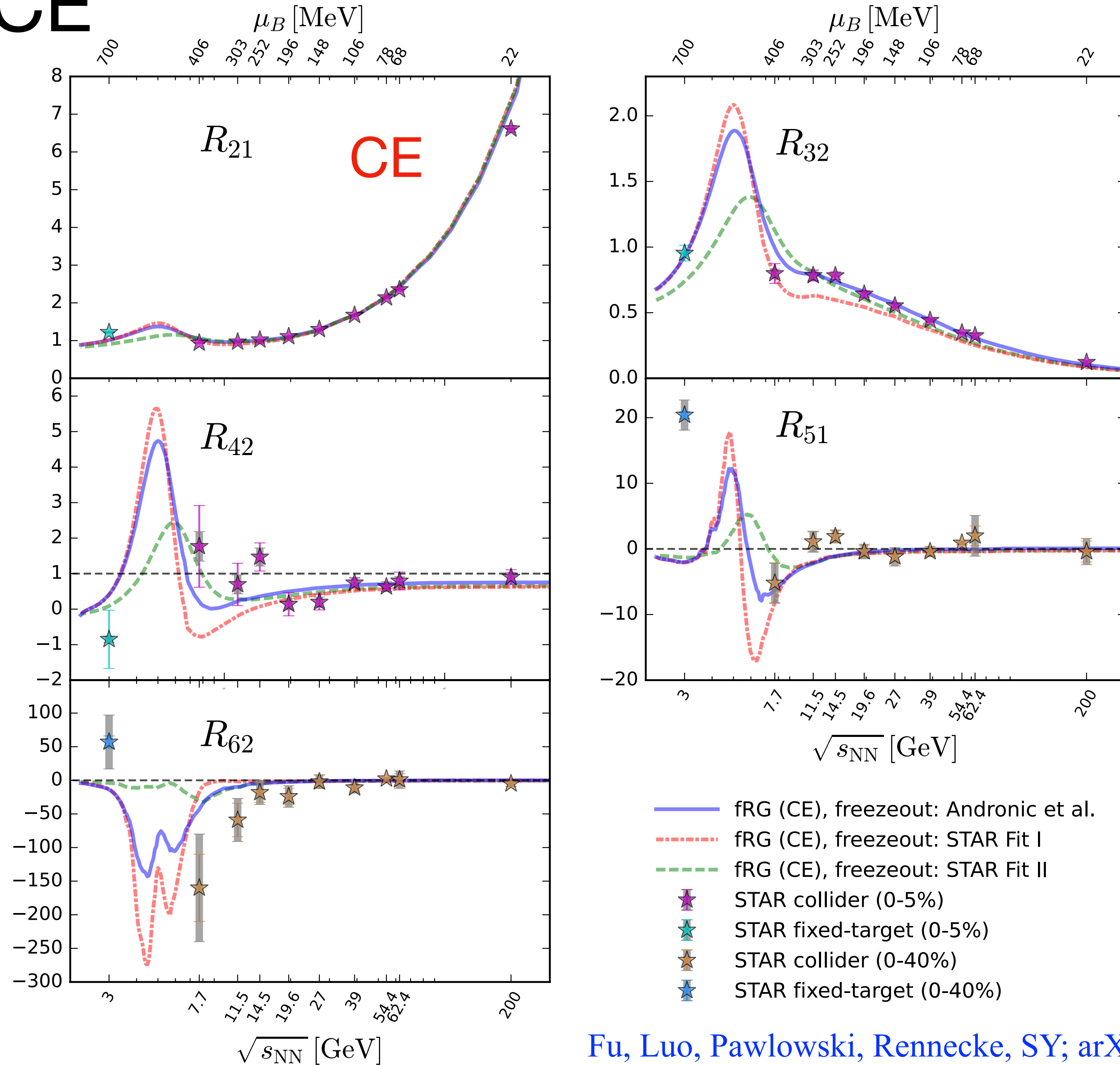
$$\frac{\kappa_2[B_1]}{\kappa_1[B_1]} = (1 - \alpha) \frac{\chi_2^B}{\chi_1^B},$$

$$\frac{\kappa_3[B_1]}{\kappa_2[B_1]} = (1 - 2\alpha) \frac{\chi_3^B}{\chi_2^B},$$

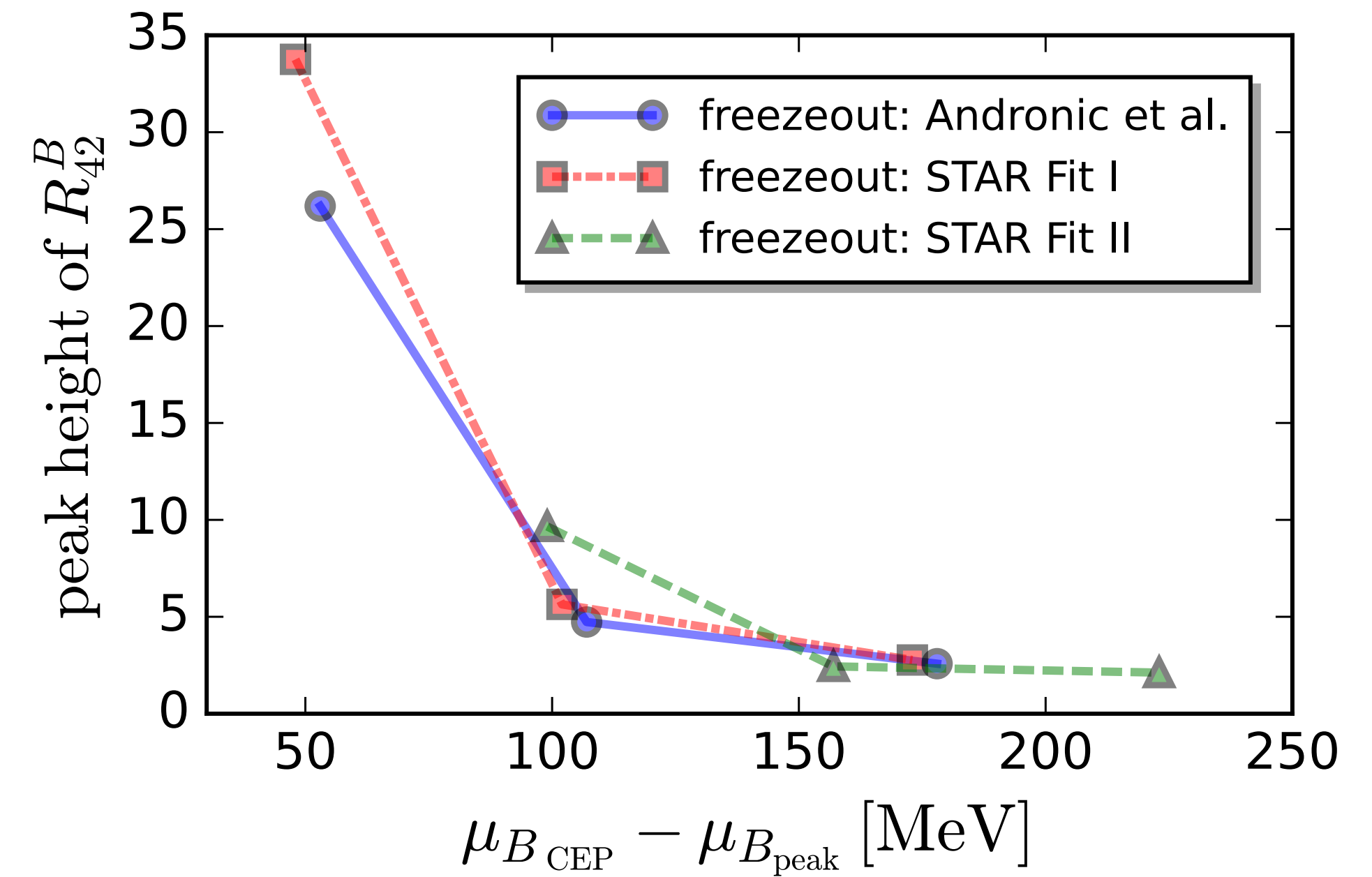
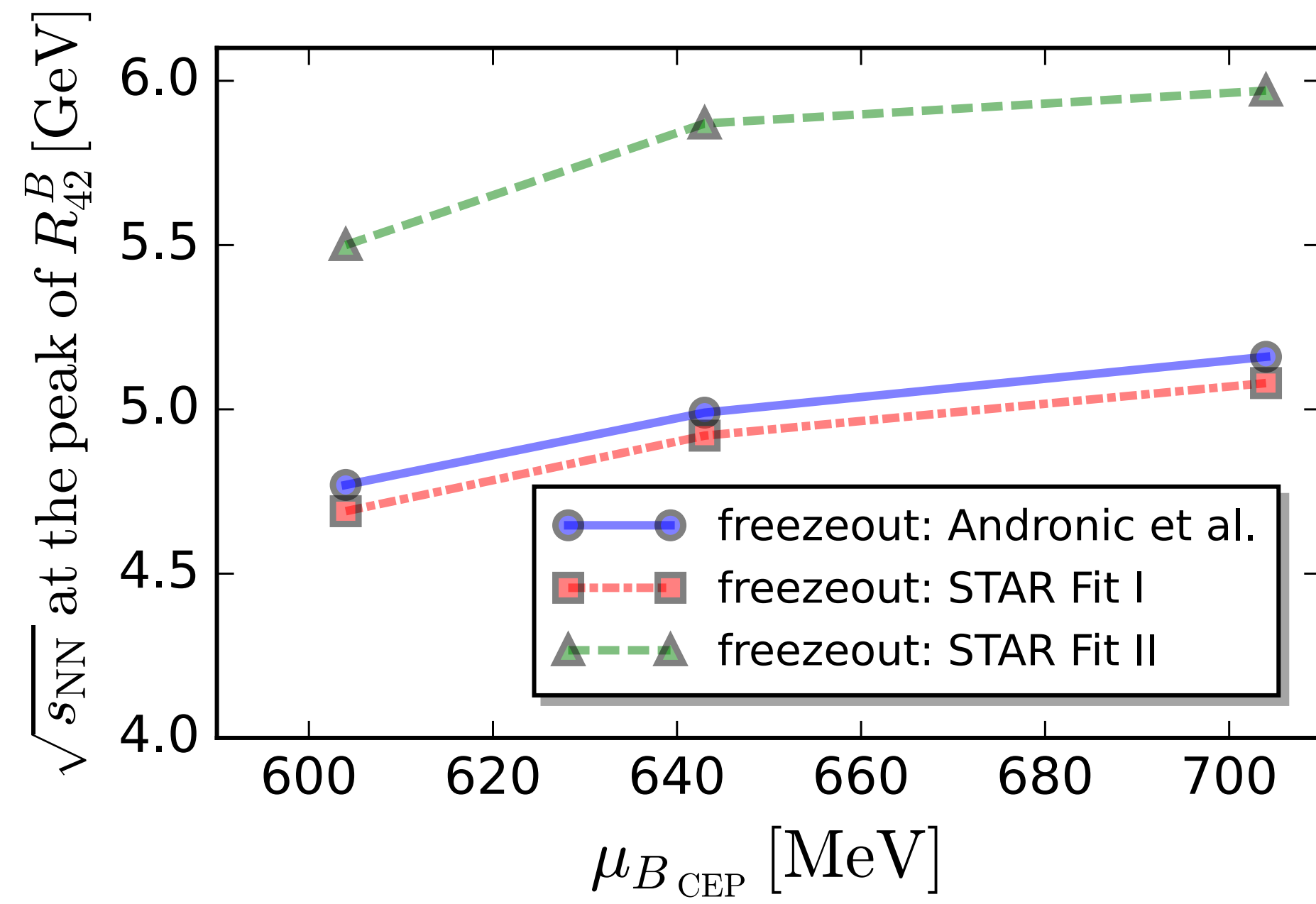
$$\frac{\kappa_4[B_1]}{\kappa_2[B_1]} = (1 - 3\alpha\beta) \frac{\chi_4^B}{\chi_2^B} - 3\alpha\beta \left(\frac{\chi_3^B}{\chi_2^B} \right)^2$$



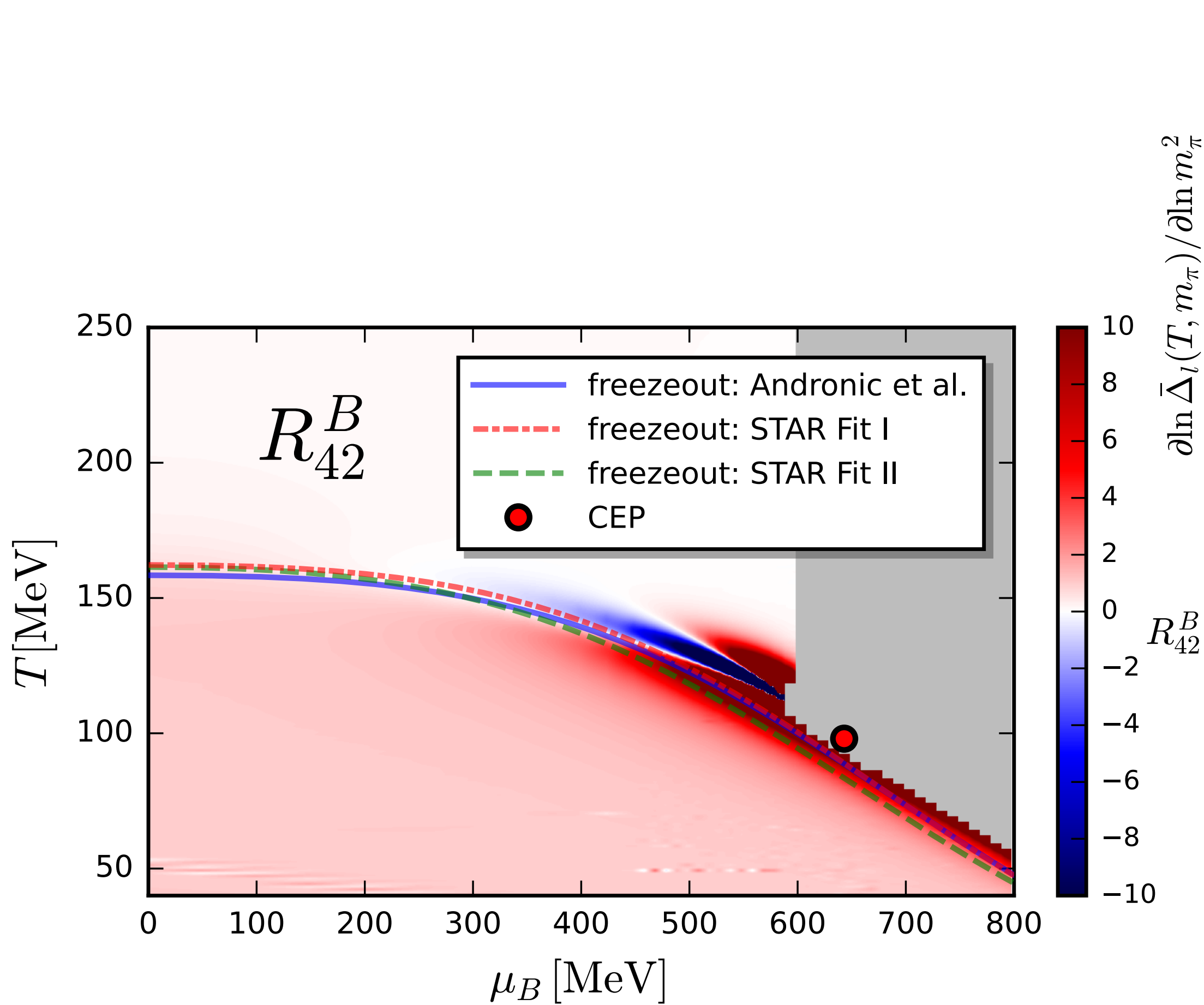
From GCE to CE



Different location of CEP & freeze-out curves

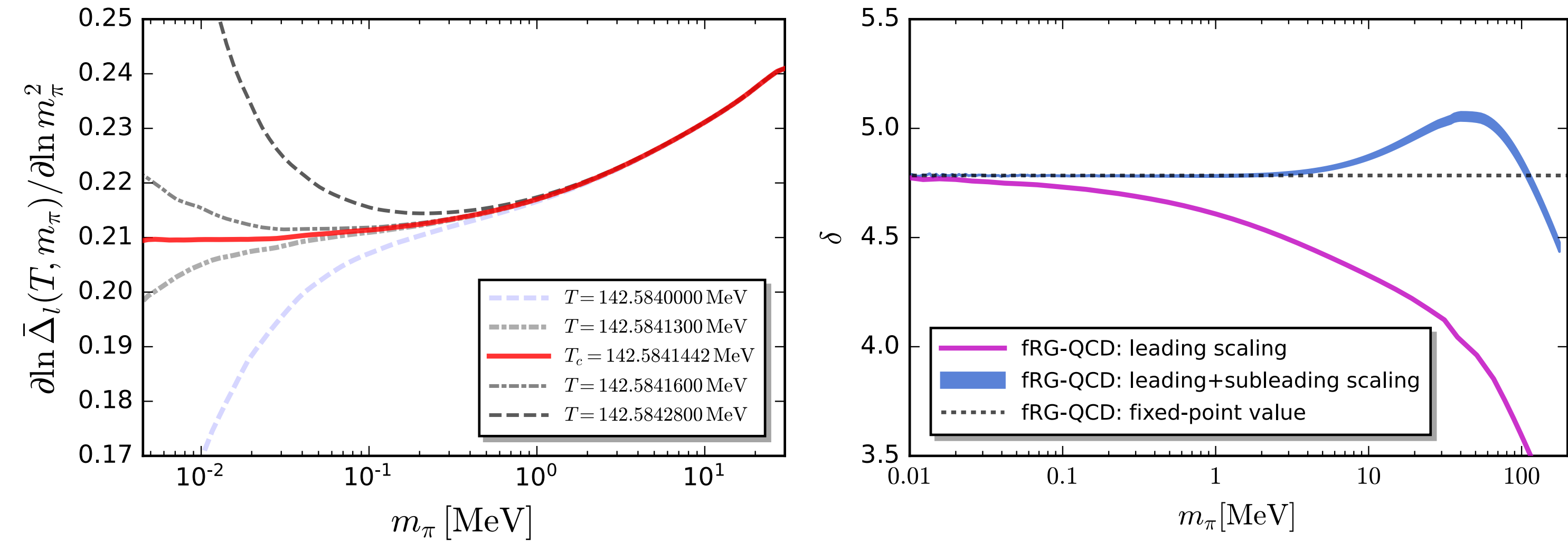


Non-monotonic energy dependence



Fu, Luo, Pawłowski, Rennecke, Wen, SY, *arXiv:2308.15508*

Jens Braun et al.; *arXiv: 2310.19853*



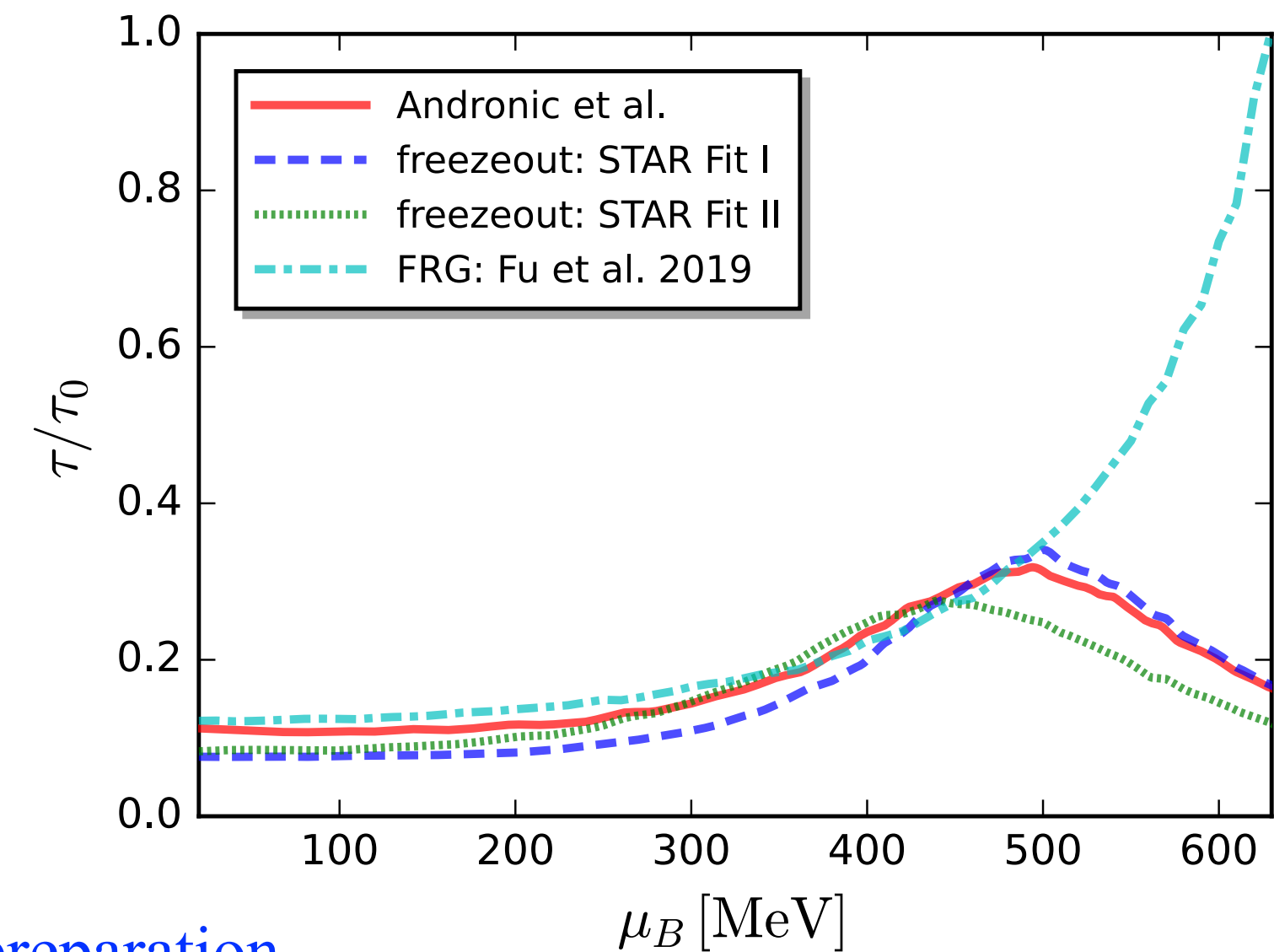
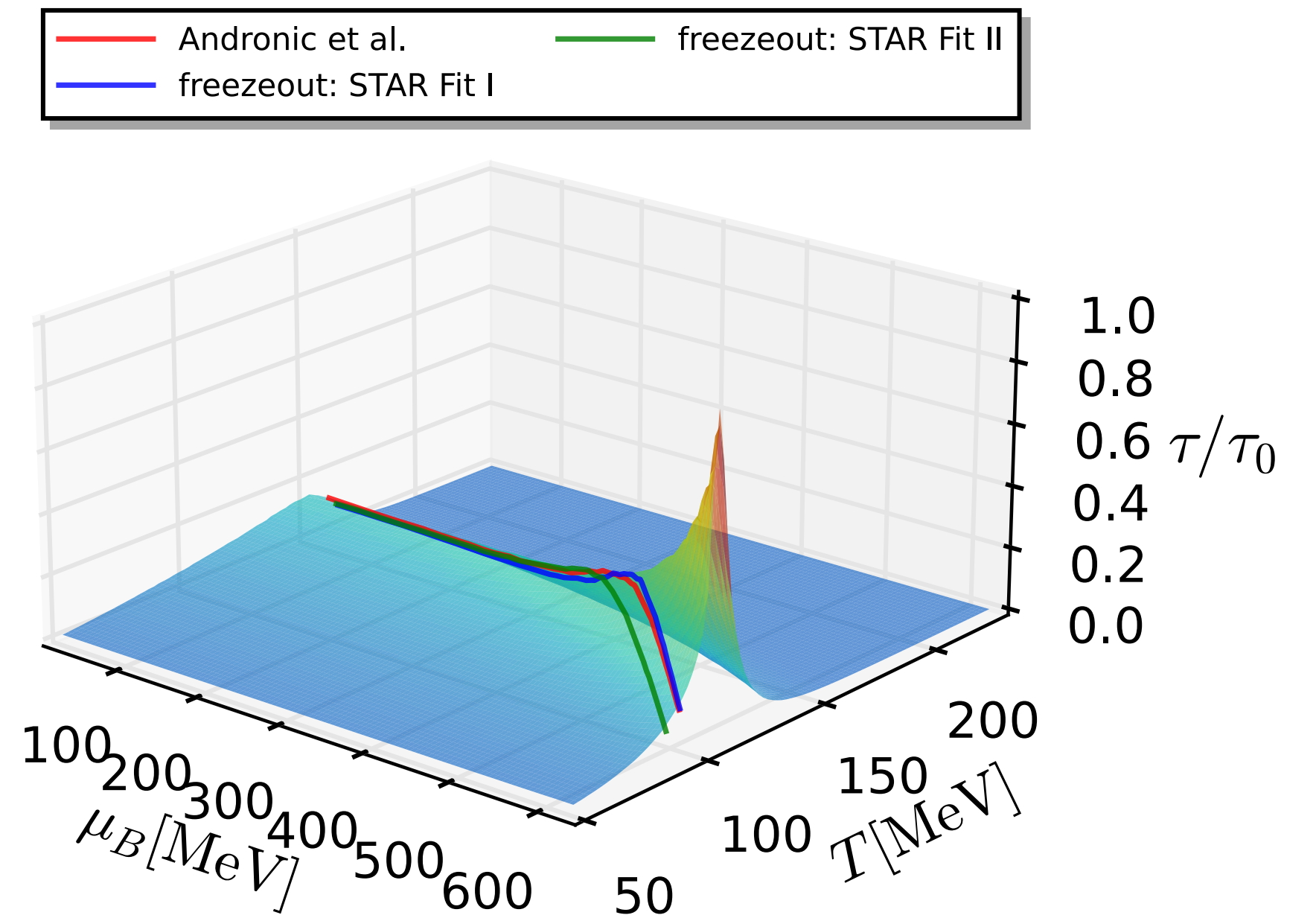
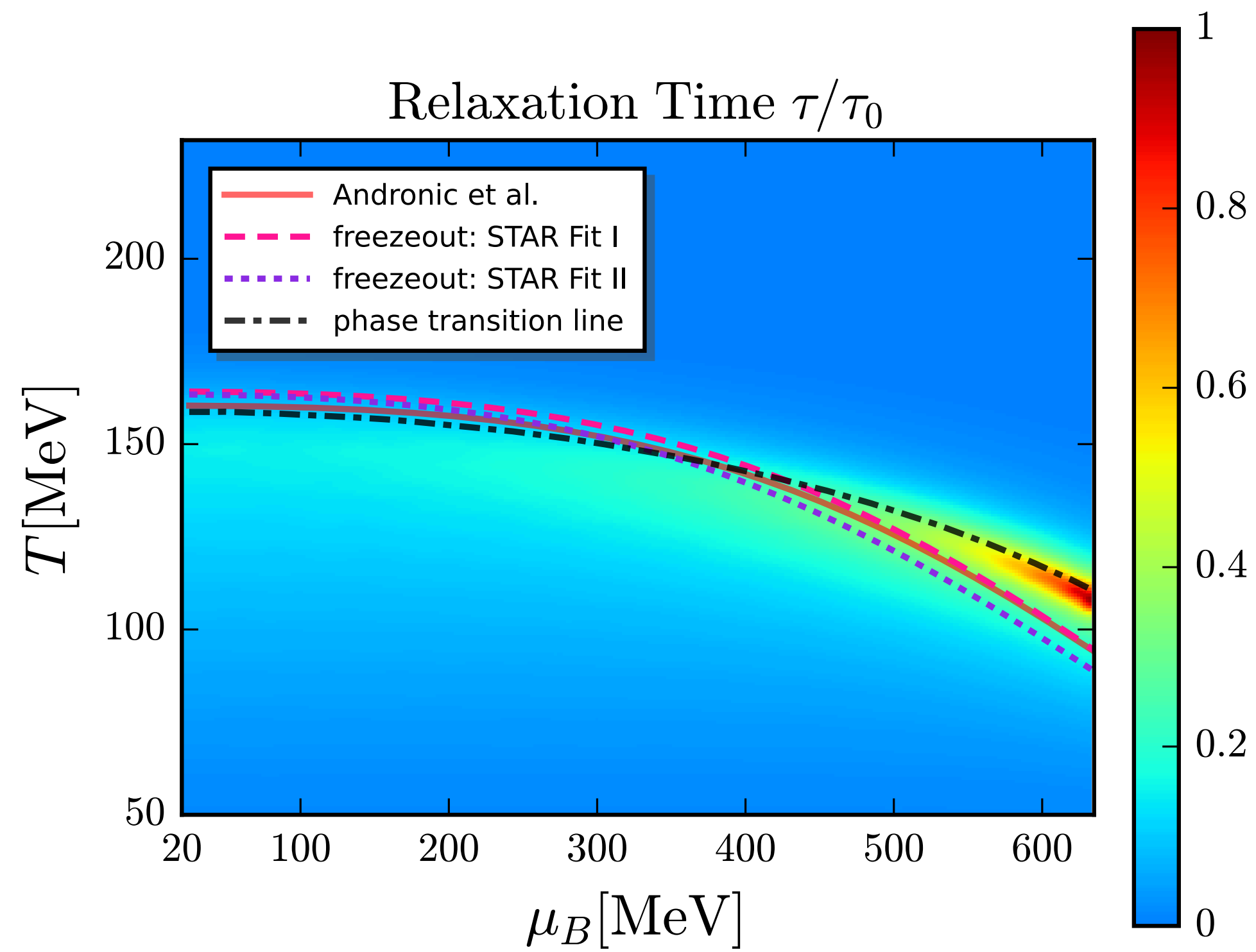
$$\bar{\Delta}_l^{\text{crit}}(m_\pi) = B_c m_\pi^{2/\delta} (1 + a_m m_\pi^{2\theta_H})$$

★ Non-monotonic energy dependence comes from gradually sharp crossover

★ Freeze-out curves are outside the critical region

Non-equilibrium effect

Obtained by Langevin equation with fRG-QCD data



Fu, Pawlowski, Rennecke, Tan, SY; in preparation

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Summary

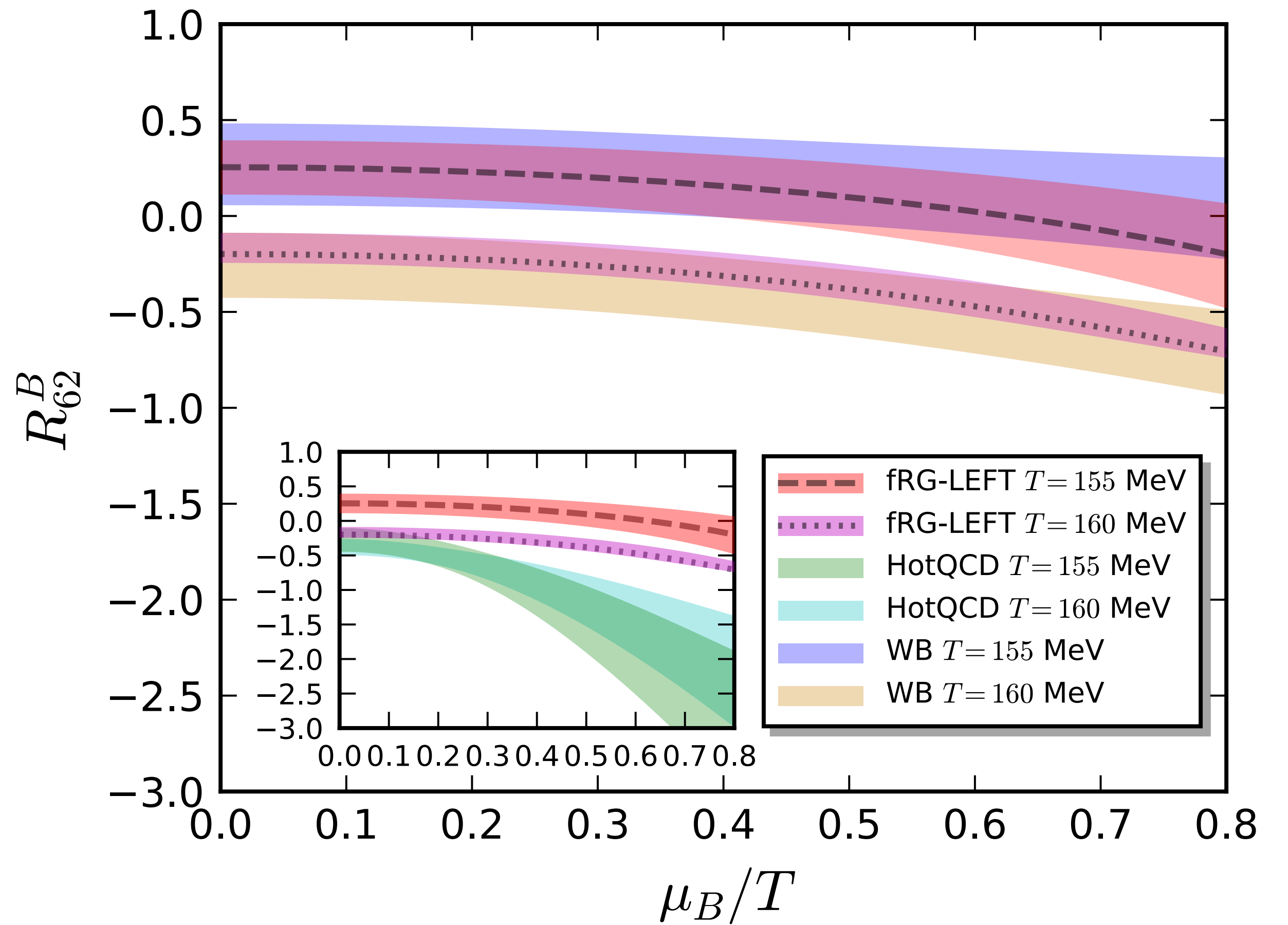
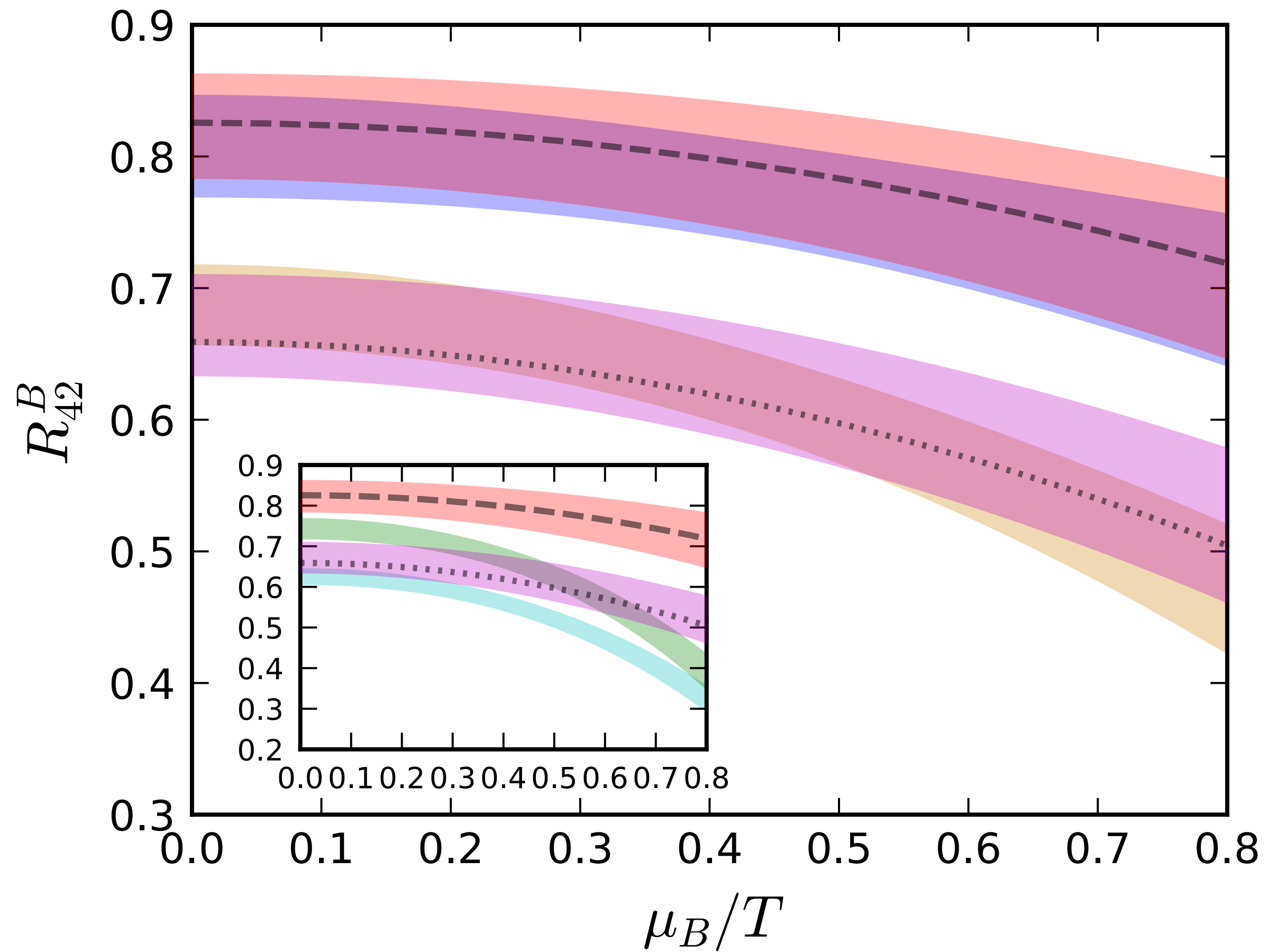
- Two steps toward quantitative calculation of baryon number fluctuations
 - QCD-assisted LEFT
 - Updated QCD-assisted LEFT
- Included global charge conservation effect
 - Subensemble acceptance method

Outlooks

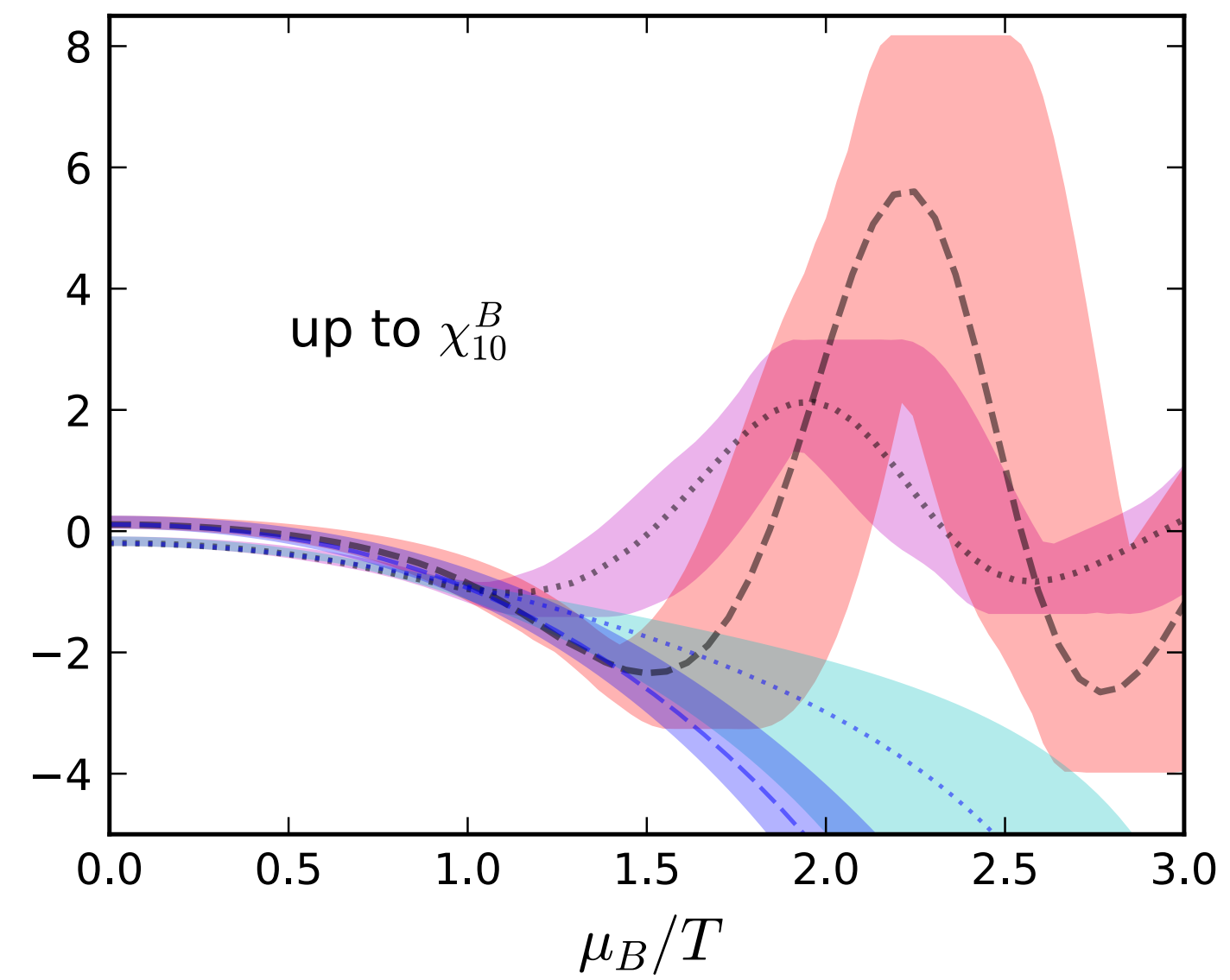
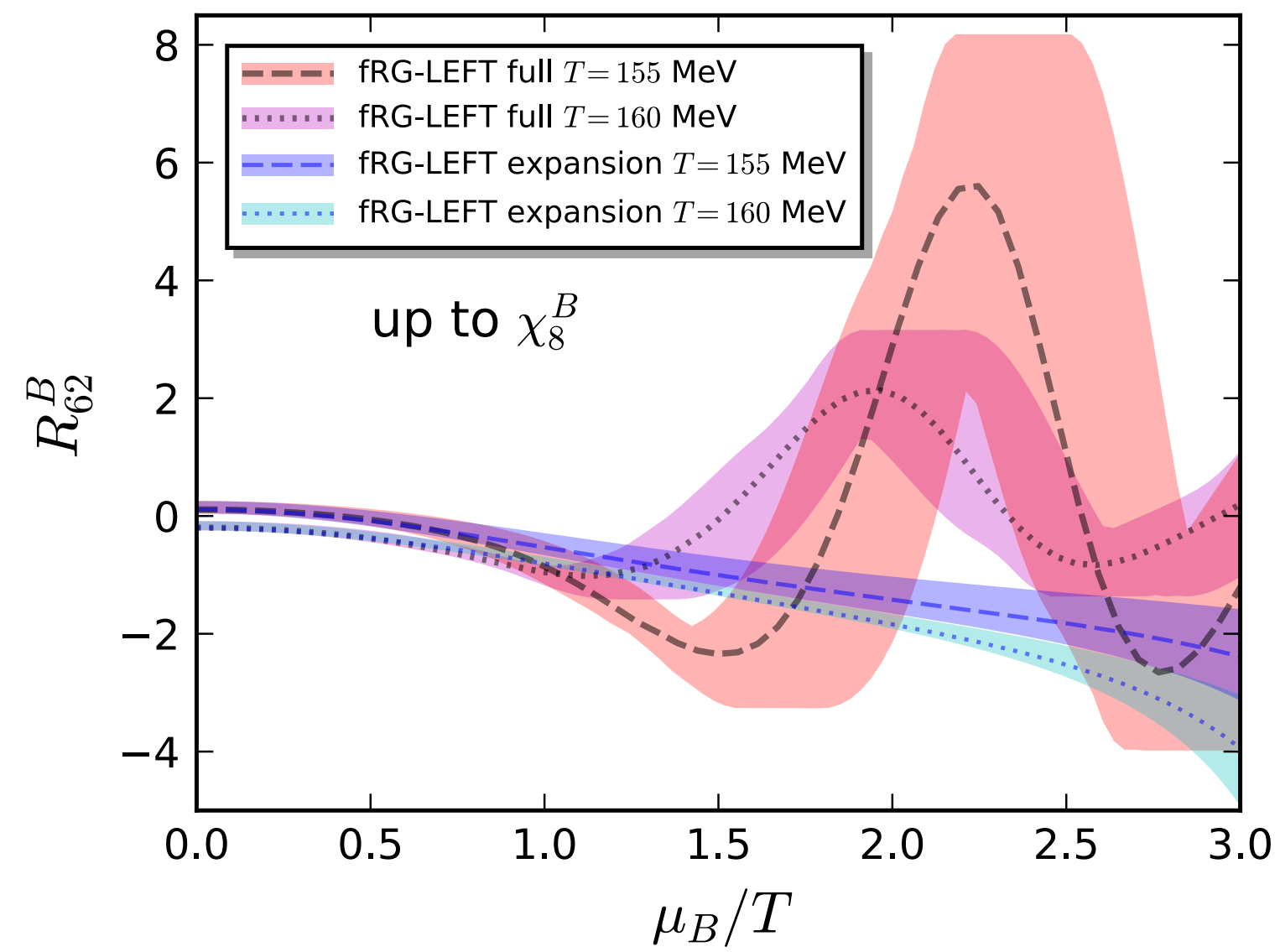
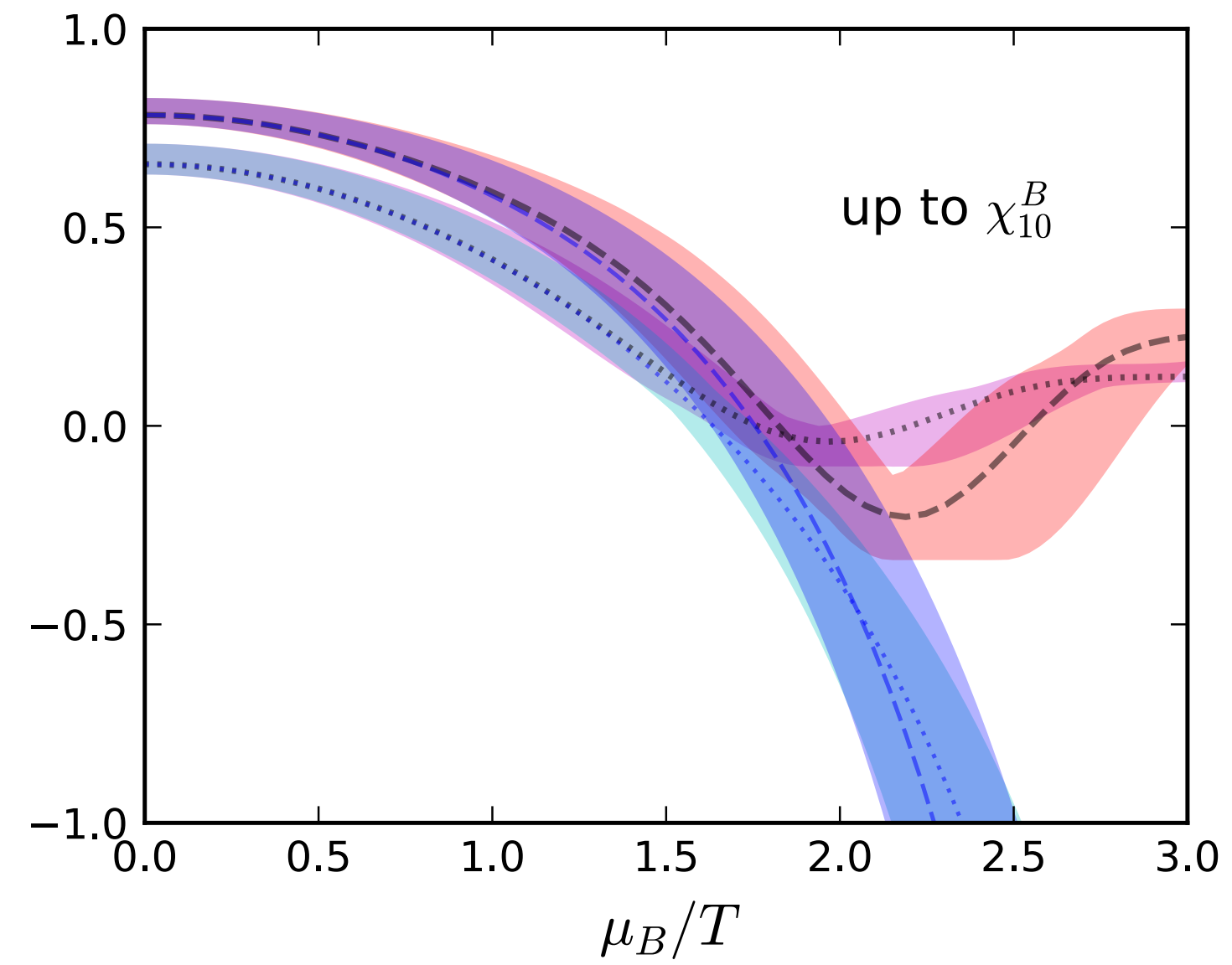
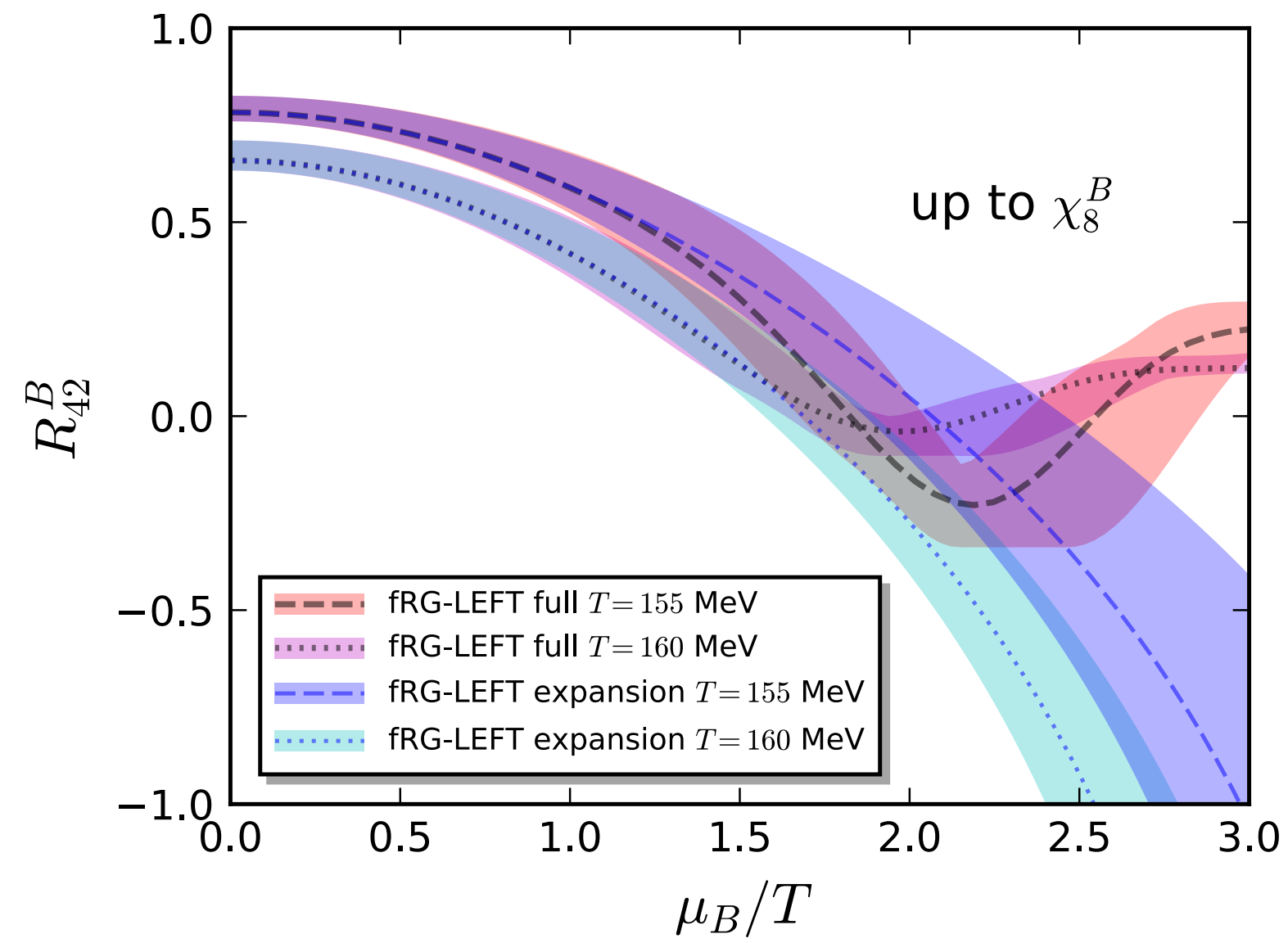
- The third step: Fluctuations from full QCD system
- Include non-equilibrium effect
- Looking forward to more accurate freeze-out curve

Thank you
very much !!!!

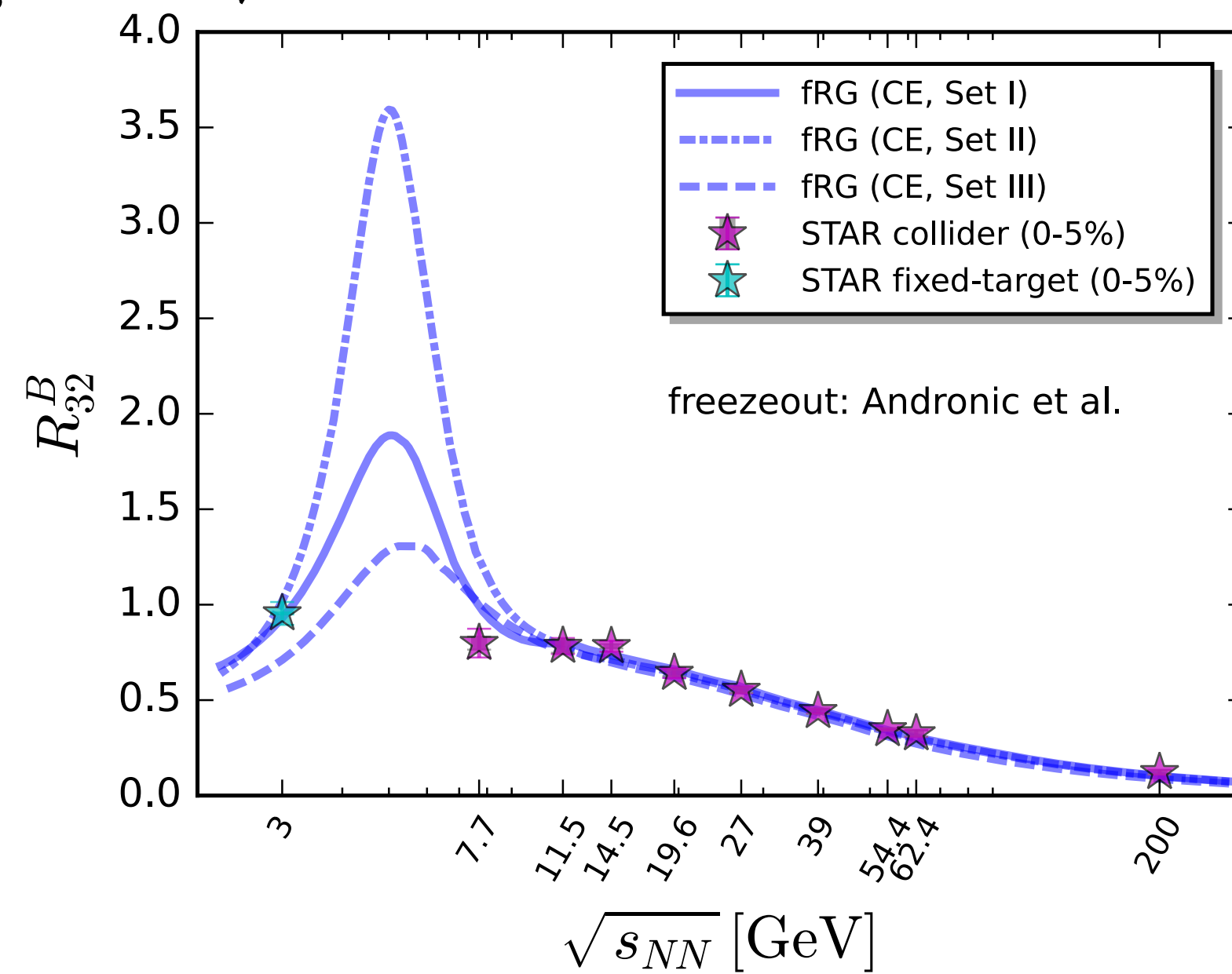
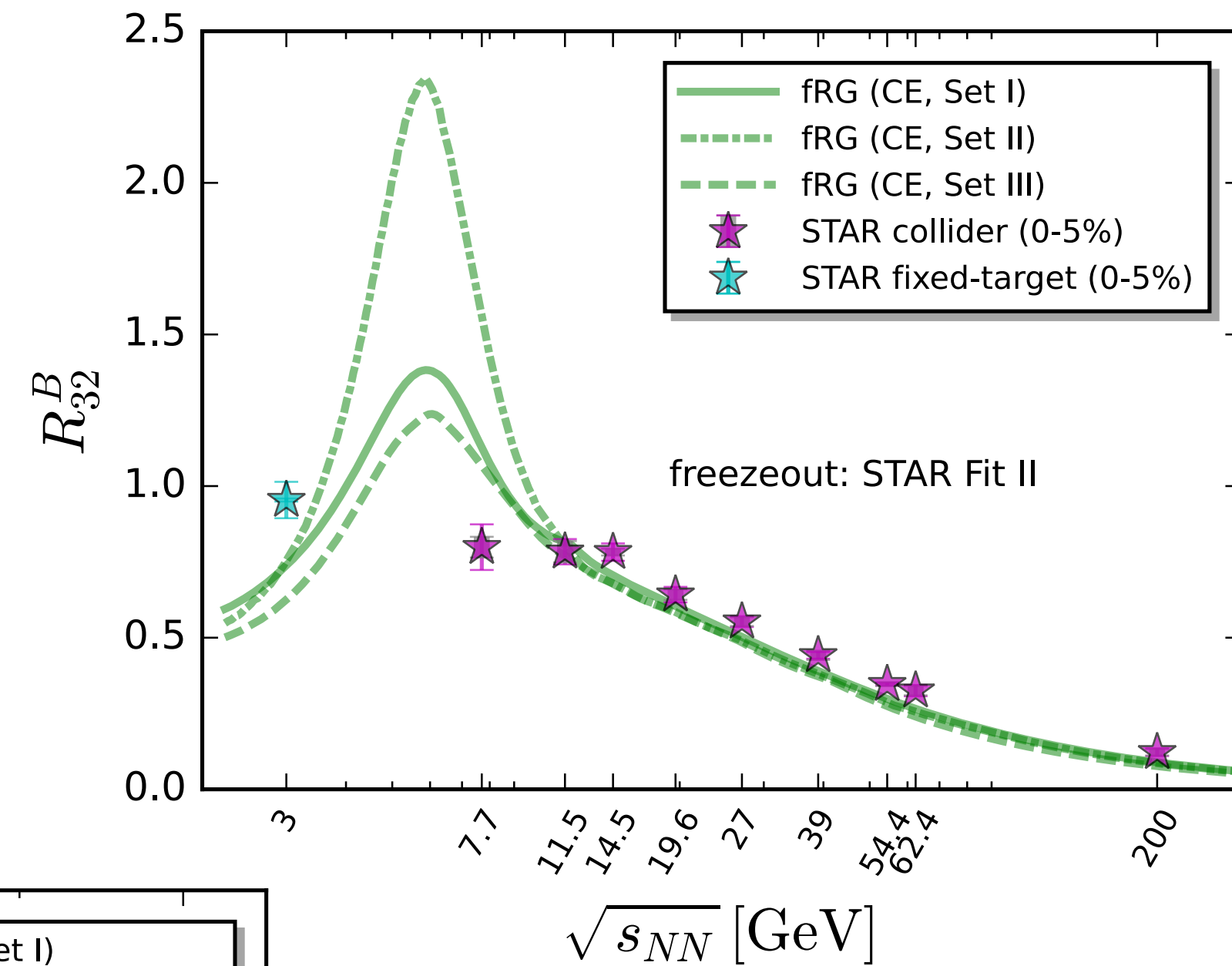
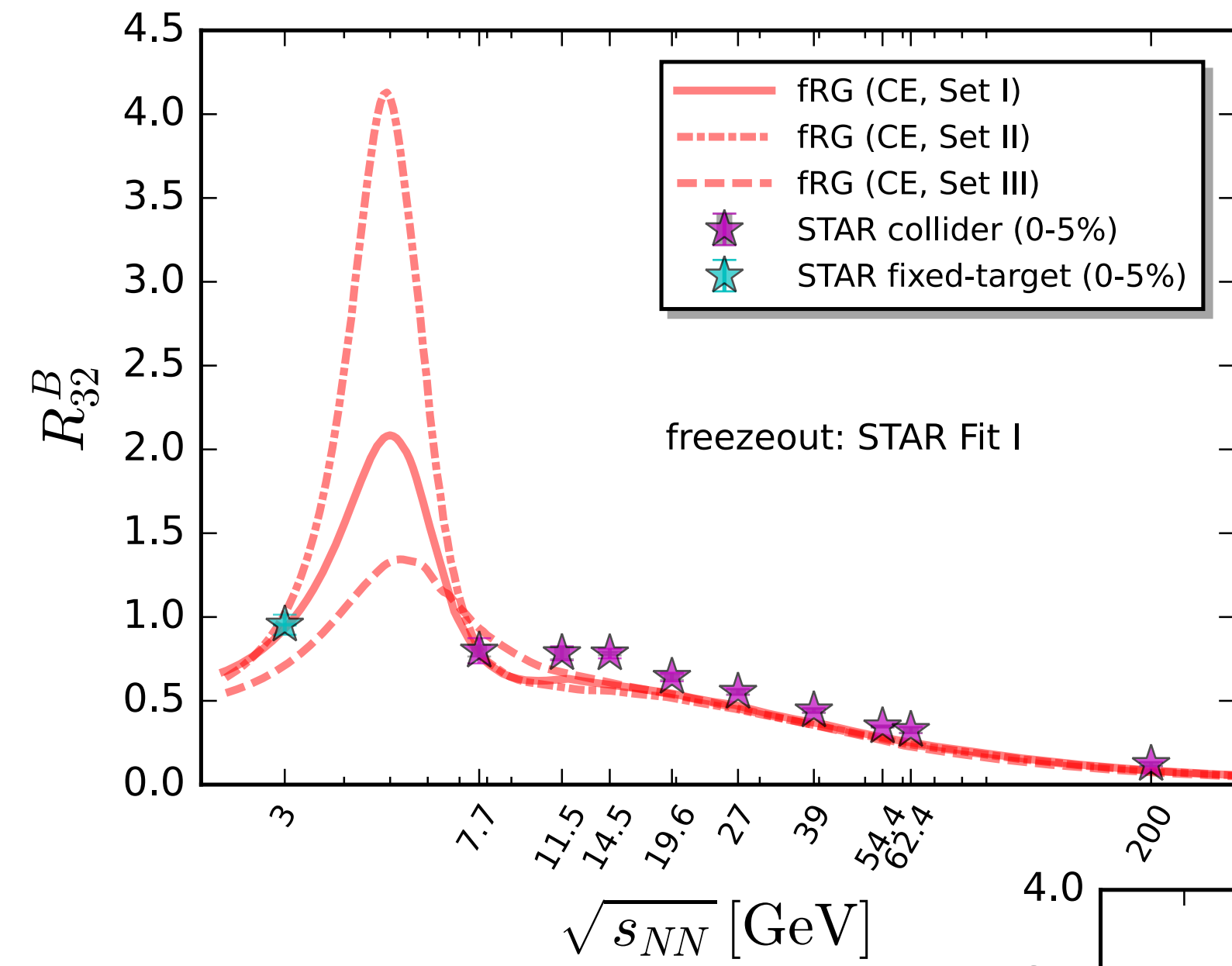
Backup



Backup



Backup



Backup

