Data-driven pole determination of (overlapping) resonances

(based on Phys. Lett. B 839, 137809 (2023))

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Lunch Club Seminar, Giessen, May 17, 2023







7th Lunch Club talk

Spectral Functions and Transport Coefficients from the Functional Renormalization Group

CRC-TR 21

Ralf-Arno Tripolt¹, Lorenz von Smekal², Jochen Wambach^{1,3}

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Seminar "Theoretische Kern- und Hadronenphysik" Justus-Liebig-Universität Gießen, October 14, 2015

Spectral functions with the FRG

Ralf-Arno Tripolt

In collaboration with

Chris Jung, Fabian Rennecke, Naoto Tanii,

Lorenz von Smekal, Jochen Wambach, Johannes Wevrich

Lunch Club Seminar, Justus-Liebig-Universität Giessen, November 21, 2018.

HGS-HIRe for FAIR

H-QM State Manager Barbard

clober 14, 2015 | Ruf Areo Youit | Seechst Functions and Transport Coefficients Iron the FRD

UNIVERSITÄT

The "Resonances Via Padé" (RVP) Method

Ralf-Arno Tripolt, ECT*, Trento, Italy

Based on arXiv: 1610.03252 Ralf-Arno Tripolt, Idan Haritan, Jochen Wambach, Nimrod Moisevev

Gießen, December 7th, 2016

ECT*

1/64

Fermionic spectral functions with the Functional Renormalization Group

Ralf-Arno Tripolt (Goethe University Frankfurt)

Lunch Club Seminar, Justus-Liebig-Universität Giessen, December 11, 2019





Theory overview on dileptons

Ralf-Arno Tripolt (Justus-Liebig-University Giessen)

Lunch Club Seminar

Giessen, 10 November, 2021







What we are going to do: Identify poles



What we are going to do: Try to identify poles



What we are going to do: Try very hard to identify poles



Outline

I) Padé approximants and the Schlessinger Point Method (SPM)

- comparison of different methods: single-point Padé, multi-point Padé, SPM
- analytic continuation to the complex plane

II) Application to experimental data

- setting up a suitable SPM algorithm
- \blacktriangleright benchmark on models for S-wave $J/\psi \to \gamma \pi^0 \pi^0$ data
- application to BESIII data

III) Summary and outlook

Padé approximants

and

the Schlessinger Point Method (SPM)

Single-point Padé

We want to describe a function f(x) in terms of a rational function, i.e. the Padé approximant of order [m/n]:

$$R_n^m(x) = \frac{P_m(x)}{1 + Q_n(x)} = \frac{\sum_{i=0}^m a_i x^i}{1 + \sum_{j=1}^n b_j x^j}$$

The coefficients a_i and b_j are obtained from **derivatives** of f(x) at a single point, e.g. at x = 0:

$$f(x) = P_m(x) - f(x)Q_n(x)$$

$$f'(x) = P'_m(x) - f'(x)Q_n(x) - f(x)Q'_n(x)$$

$$f''(x) = P''_m(x) - f''(x)Q_n(x) - f(x)Q''_n(x) - 2f'(x)Q'_n(x)$$

[G. A. J. Baker, Essentials of Padé Approximants, Academic Press, Cambridge, Massachusetts, 1975]
 [P. Dimopoulos et. al. (C. Schmidt), Phys. Rev. D 105, 034513 (2022)]

. . .

Example:

Find the Padé approximant of order [m/n] = [2/3] of $f(x) = \cos(x)$:

$$R_3^2(x) = \frac{a_0 + a_1 x + a_2 x^2}{1 + b_1 x + b_2 x^2 + b_3 x^3}$$

We need up to the m + n = 5th derivative of f(x) at x = 0 to fix all coefficients!



Example:

Find the Padé approximant of order [m/n] = [2/3] of $f(x) = \cos(x)$:

$$R_3^2(x) = \frac{a_0 + a_1 x + a_2 x^2}{1 + b_1 x + b_2 x^2 + b_3 x^3}$$

We need up to the m + n = 5th derivative of f(x) at x = 0 to fix all coefficients:

$$R_3^2(x) \approx \frac{1 - 0.4167 \, x^2}{1 + 0.0833 \, x^2 + 0 \, x^3}$$



Multi-point Padé

We want to describe a function f(x) in terms of a rational function, i.e. the Padé approximant of order [m/n]:

$$R_n^m(x) = \frac{P_m(x)}{1 + Q_n(x)} = \frac{\sum_{i=0}^m a_i x^i}{1 + \sum_{j=1}^n b_j x^j}$$

The coefficients a_i and b_j are obtained from function values and derivatives of f(x) at multiple points:

$$P_m(x_1) - f(x_1)Q_n(x_1) = f(x_1)$$

$$P'_m(x_1) - f'(x_1)Q_n(x_1) - f(x_1)Q'_n(x_1) = f'(x_1)$$
...
$$P_m(x_2) - f(x_2)Q_n(x_2) = f(x_2)$$

$$P'_m(x_2) - f'(x_2)Q_n(x_2) - f(x_2)Q'_n(x_2) = f'(x_2)$$
...
$$P_m(x_N) - f(x_N)Q_n(x_N) = f(x_N)$$

$$P'_m(x_N) - f'(x_N)Q_n(x_N) = f'(x_N)$$

[G. A. J. Baker, Essentials of Padé Approximants, Academic Press, Cambridge, Massachusetts, 1975]
 [P. Dimopoulos et. al. (C. Schmidt), Phys. Rev. D 105, 034513 (2022)]

Example:

Find the Padé approximant of order [m/n] = [2/3] of $f(x) = \cos(x)$:

$$R_3^2(x) = \frac{a_0 + a_1 x + a_2 x^2}{1 + b_1 x + b_2 x^2 + b_3 x^3}$$

We use function values and first derivatives at three different points to determine all coefficients.



Example:

Find the Padé approximant of order [m/n] = [2/3] of $f(x) = \cos(x)$:

$$R_3^2(x) = \frac{a_0 + a_1 x + a_2 x^2}{1 + b_1 x + b_2 x^2 + b_3 x^3}$$

We use function values and first derivatives at three different points:

$$R_3^2(x) \approx \frac{1 + 0.143 \, x - 0.497 \, x^2}{1 + 0.140 \, x + 0.0194 \, x^2 + 0.0374 \, x^3}$$



Schlessinger Point Method (SPM)

Given a finite set of N data points (x_i, y_i) we construct the rational interpolant p(x)/q(x) with polynomials p(x) and q(x) that is given by the continued fraction

$$p(x)/q(x) = C_N(x) = \frac{y_1}{1 + \frac{a_1(x - x_1)}{1 + \frac{a_2(x - x_2)}{\vdots a_{N-1}(x - x_{N-1})}}}$$

where the coefficients a_i are given recursively by $a_1 = \frac{y_1/y_2 - 1}{x_2 - x_1}$ and

$$a_{i} = \frac{1}{x_{i} - x_{i+1}} \left(1 + \frac{a_{i-1}(x_{i+1} - x_{i-1})}{1 + 1} \frac{a_{i-2}(x_{i+1} - x_{i-2})}{1 + 1} \cdots \frac{a_{1}(x_{i+1} - x_{1})}{1 - y_{1}/y_{i+1}} \right)$$

The polynomials (p(x), q(x)) are of order (N/2 - 1, N/2) for an **even** number of input points and of order ((N - 1)/2, (N - 1)/2) for an **odd** number of input points

[[]L. Schlessinger, Physical Review, Volume 167, Number 5 (1968)]

[[]R.W. Haymaker and L. Schlessinger, Mathematics in Science and Engineering, Volume 71, Chapter 11 (1970)]

[[]H.J. Vidberg and J.W. Serene, Journal of Low Temperature Physics, Vol. 29, Nos. 3/4 (1977)]

Example:

Find the SPM interpolator of $f(x) = \cos(x)$ using 6 input points:

$$C_6(x) = \frac{a_0 + a_1 x + a_2 x^2}{1 + b_1 x + b_2 x^2 + b_3 x^3}$$



Example:

Find the SPM interpolator of f(x) = cos(x)using 6 input points:

$$C_6(x) \approx \frac{1 - 0.0011 \, x - 0.41 \, x^2}{1 + 0.00027 \, x + 0.11 \, x^2 + 0.0032 \, x^3}$$



Gives the same result as multi-point Padé when only using function values (no derivatives) and choosing the corresponding order of polynomials (e.g. [2/3] for 6 input points)!

Identifying branch cuts with the SPM

Example:

We construct the SPM interpolator of $f(x) = \sqrt{x}$ using 15 input points:

- branch cut shows up in the SPM interpolator as a series of poles!
- the location of the branch point is given by the first pole



Complex poles from data and fits



Comparison of reconstructions

exact

MEM

BG

SPM

1000

100

0.100

0.010

0.001

200

ρ [GeV⁻²]

Analytic structure from DSE/lattice



[R.-A. T., I. Haritan, J. Wambach, N. Moiseyev Phys. Lett. B 774, 411-416 (2017)]



ω [MeV]

600

800

400

[D. Binosi, R.-A. T., Phys. Lett. B 801, 135171 (2020)]

For N = 2 input points:

- ► since N is even, the resulting rational interpolant will be of order (N/2 − 1, N/2)
- \blacktriangleright for N=2 we expect

$$C_N(x) = \frac{p(x)}{q(x)} = \frac{p_0}{q_0 + q_1 x}$$



For N = 2 input points:

- ► since N is even, the resulting rational interpolant will be of order (N/2 - 1, N/2)
- \blacktriangleright for N=2 we expect

$$C_N(x) = \frac{p(x)}{q(x)} = \frac{p_0}{q_0 + q_1 x}$$

 the resulting SPM interpolator is given by

$$C_N(x) = \frac{1.2}{1 - 0.2 x}$$



For N = 3 input points:

- ► since N is odd, the resulting rational interpolant will be of order (N/2, N/2)
- \blacktriangleright for N=3 we expect

$$C_N(x) = \frac{p(x)}{q(x)} = \frac{p_0 + p_1 x}{q_0 + q_1 x}$$



For N = 3 input points:

- ▶ since N is odd, the resulting rational interpolant will be of order (N/2, N/2)
- \blacktriangleright for N = 3 we expect

 $C_N(x) = rac{p(x)}{q(x)} = rac{p_0 + p_1 x}{q_0 + q_1 x}$

- ▶ for exact data, the linear function is recovered, i.e. ~ x/1
- if we add noise of the order of $\mathcal{O}(10^{-15})$ we get, e.g.,

$$C_N(x) = \frac{1.81 \cdot 10^{-14} + x}{1 - 1.78 \cdot 10^{-15} x} \approx x$$



For N = 4 input points:

- ► since N is even, the resulting rational interpolant will be of order (N/2 - 1, N/2)
- \blacktriangleright for N = 4 we expect

$$C_N(x) = \frac{p(x)}{q(x)} = \frac{p_0 + p_1 x}{q_0 + q_1 x + q_2 x^2}$$



For N = 4 input points:

- ► since N is even, the resulting rational interpolant will be of order (N/2 - 1, N/2)
- ▶ for N = 4 we expect

$$C_N(x) = \frac{p(x)}{q(x)} = \frac{p_0 + p_1 x}{q_0 + q_1 x + q_2 x^2}$$

- ▶ however, for exact data, the linear function is recovered, i.e. ~ x/1
- ▶ if we add noise of the order of $\mathcal{O}(10^{-15})$ we get, e.g.,

$$C_N(x) = \frac{-3.2 \cdot 10^{-14} + x}{1 + 1.2 \cdot 10^{-14} x - 1.3 \cdot 10^{-15} x^2}$$



We use a Breit-Wigner function of the form

$$f(x) = \frac{A}{(x^2 - M^2)^2 + \gamma^2 M^2}$$

with parameters $A=100\text{, }M=4\text{, }\gamma=2$

For N = 14 input points we expect

$$C_N(x) = rac{p(x)}{q(x)} \sim rac{\mathcal{O}(x^6)}{\mathcal{O}(x^7)}$$



Breit-Wigner function

We use a Breit-Wigner function of the form

$$f(x) = \frac{A}{(x^2 - M^2)^2 + \gamma^2 M^2}$$

with parameters $A=100\text{, }M=4\text{, }\gamma=2$

For ${\cal N}=14$ input points we expect

$$C_N(x) = rac{p(x)}{q(x)} \sim rac{\mathcal{O}(x^6)}{\mathcal{O}(x^7)}$$

with noise of the order of $\mathcal{O}(10^{-15})$ we get, e.g.,

$$C_N(x) \approx \frac{0.313 - 0.2 x + 0.04 x^2 - 0.003 x^3 - 10^{-14} x^4 + 10^{-16} x^5 + 10^{-17} x^6}{1 - 0.7 x + 0.04 x^2 + 0.06 x^3 - 0.01 x^4 - 0.001 x^5 + 0.0004 x^6 - 0.00002 x^7}$$

 \rightarrow artificial poles (e.g. at $x\approx2.3)$ in denominator cancel with zeros in numerator!

2.0

15

€ 1.0

0.5

0.0 0 f(x) Input

SPM

2

4

х

6

Breit-Wigner function

SPM as an extrapolation:

- evaluate $C_N(x)$ outside the input range
- describes Breit-Wigner function well for precise data!



Breit-Wigner function

Analytic continuation:

- ▶ evaluate SPM interpolation function $C_N(x)$ for complex arguments: $x \rightarrow x + iy$
- describes Breit-Wigner function well for precise data!



We now add noise of $\mathcal{O}(10^{-1})$ to the same Breit-Wigner function of the form

$$f(x) = \frac{A}{(x^2 - M^2)^2 + \gamma^2 M^2}$$

with parameters $A=100\text{, }M=4\text{, }\gamma=2$

For ${\cal N}=14$ input points we expect

$$C_N(x) = rac{p(x)}{q(x)} \sim rac{\mathcal{O}(x^6)}{\mathcal{O}(x^7)}$$



Breit-Wigner function with noise

2.0 We now add noise of $\mathcal{O}(10^{-1})$ to the same f(x) Breit-Wigner function of the form Input 1.5 SPM $f(x) = \frac{A}{(x^2 - M^2)^2 + \gamma^2 M^2}$ with parameters $A = 100, M = 4, \gamma = 2$ × 1.0 For N = 14 input points we expect 0.5 $C_N(x) = \frac{p(x)}{q(x)} \sim \frac{\mathcal{O}(x^5)}{\mathcal{O}(x^7)}$ 0.0^L 2 Λ 6 and get, e.g., х $C_N(x) \approx \frac{0.345 - 0.6 x + 0.3 x^2 - 0.1 x^3 + 0.01 x^4 - 0.001 x^5 + 0.00003 x^6}{1 - 1.9 x + 1.3 x^2 - 0.5 x^3 + 0.1 x^4 - 0.02 x^5 + 0.001 x^6 - 0.00005 x^7}$

 \rightarrow artificial poles in denominator no longer cancel perfectly, but they appear randomly!

Breit-Wigner function with noise

Analytic continuation to the complex plane:

- ▶ main "particle" poles are slightly shifted by the noise but still clearly visible
- ▶ artificial poles can be identified due to their random location



Application to experimental data

S-wave $J/\psi ightarrow \gamma \pi^0 \pi^0$ data from BESIII

- ▶ results show three peaks above 1 GeV, likely associated with $f_0(1500)$, $f_0(1710)$ and $f_0(2020)$
- ▶ together with the f₀(1370) they are the main candidates for the lightest glueball
- process contributing to $J/\psi \rightarrow \gamma \pi^0 \pi^0$:





intensities for the 0^{++} amplitudes as a function of $M(\pi_0\pi_0)$

[M. Ablikim et al. (BESIII Collaboration), Phys. Rev. D 92, 052003 (2016)]
 [JPAC Collaboration: A. Rodas, A. Pilloni, M. Albaladejo, C. Fernandez-Ramirez, V. Mathieu, A. P. Szczepaniak, Eur.Phys.J.C 82, 80 (2022)]
 [D. Binosi, A. Pilloni, R.-A. T., Phys. Lett. B 839, 137809 (2023)]

Models

S- and D-wave intensities and relative phase of $J/\psi \rightarrow \gamma \pi^0 \pi^0$ and $\rightarrow \gamma K^0_S K^0_S$ were parametrized by JPAC based on the coupled-channel N/D formalism (using unitarity and dispersion relations):

$$I_i^J(s) = \mathcal{N}p_i \left| a_i^J(s) \right|^2$$
$$a_i^J(s) = E_{\gamma} p_i^J \sum_k n_k^J(s) \left[D^J(s)^{-1} \right]_{ki}$$
$$p_i = \sqrt{s - 4m_i^2}/2 \qquad E_{\gamma} = (m_{J/\psi}^2 - s)/(2\sqrt{s})$$

with $i=\pi\pi$ or $K\bar{K}$ and the invariant mass squared s.

The resulting **5 models (A-E)** contain either 3 or 4 poles and sometimes an additional branch point at E = 1.52 GeV from the stable $\rho\rho$ channel.



 [[]JPAC Collaboration: A. Rodas, A. Pilloni, M. Albaladejo, C. Fernandez-Ramirez, V. Mathieu, A. P. Szczepaniak, Eur.Phys.J.C 82 (2022) 1, 80]
 [D. Binosi, A. Pilloni, R.-A. T., Phys. Lett. B 839, 137809 (2023)]
 [J.R. Pelaez, Physics Reports 658 (2016) 1]

Statistical SPM



[D. Binosi, A. Pilloni, R.-A. T., Phys. Lett. B 839, 137809 (2023)]

100,000 runs: separating signal from noise



- step 1: increase transparency
- step 2: use morphological pattern recognition (e.g. with Mathematica: MorphologicalBinarize and MorphologicalComponent)
- step 3: ignore 'horizontal' noise clusters close to real axis
- ▶ final signal pole clusters contain e.g. 50,000 poles which can be used to compute the standard deviation
- \blacktriangleright on average, the number of signal poles is $\sim 20\%$ of all poles

Validation of the method - Model A - exact data



- stars denote exact pole locations (red: model artifact)
- all poles are correctly identified for exact data (i.e. data without uncertainties)!
- > noise poles are almost exactly canceled by corresponding zeros in numerator

Validation of the method - Model A

stars denote exact pole locations

▶ different colors correspond to different numbers of input points for the SPM (M = 20,...,60)

▶ 3 rows correspond to different amount of full experimental errors: 10%, 33%, 100%

 ellipses indicate 1σ (i.e. 68%) confidence regions



Validation of the method - Model A



all physical poles are correctly identified!

Validation of the method - Model B



- model B turns out to be the most difficult for the SPM
- third pole is not identified correctly for large errors (too deep in the comlex plane)

Validation of the method - Model C



- branch cut does not affect the pole reconstruction
- all poles are identified correctly

Validation of the method - Model D



- third and fourth pole seem to be reconstructed as a single "average" pole
- as the error increases the signal of the fourth pole faints and the correct position of the pole is recovered

Validation of the method - Model E



- branch cut does not affect the pole reconstruction
- > poles are correctly identified for even values of M (i.e., the ones with a vanishing asymptotic)

S-wave $J/\psi ightarrow \gamma \pi^0 \pi^0$ data from BESIII

- ▶ results show three peaks above 1 GeV, likely associated with $f_0(1500)$, $f_0(1710)$ and $f_0(2020)$
- ▶ together with the f₀(1370) they are the main candidates for the lightest glueball
- process contributing to $J/\psi \to \gamma \pi^0 \pi^0$:





intensities for the 0^{++} amplitudes as a function of $M(\pi_0 \pi_0)$

[M. Ablikim et al. (BESIII Collaboration), Phys. Rev. D 92, 052003 (2016)]
 [JPAC Collaboration: A. Rodas, A. Pilloni, M. Albaladejo, C. Fernandez-Ramirez, V. Mathieu, A. P. Szczepaniak, Eur.Phys.J.C 82, 80 (2022)]
 [D. Binosi, A. Pilloni, R.-A. T., Phys. Lett. B 839, 137809 (2023)]

Application to BESIII data: results



[D. Binosi, A. Pilloni, R.-A. T., Phys. Lett. B 839, 137809 (2023)] [JPAC Collaboration: A. Rodas, A. Pilloni, M. Albaladejo, C. Fernandez-Ramirez, V. Mathieu, A. P. Szczepaniak, Eur.Phys.J.C 82, 80 (2022)] [Bonn-Gatchina: A. V. Sarantsev, I. Denisenko, U. Thoma, E. Klempt, Phys. Lett. B816, 136227 (2021)] [S. Ropertz, C. Hanhart, B. Kubis, Eur. Phys. J. C78, 1000 (2018)]

Application to **BESIII** data

Pole positions obtained from different data analyses (in MeV):

	SPM (this work)	JPAC	Bonn-Gatchina	Ropertz <i>et al.</i>
$f_0(1500)$	$(1449 \pm 24) - i(100 \pm 32)/2$	$(1450 \pm 10) - i(106 \pm 16)/2$	$(1483 \pm 15) - i(116 \pm 12)/2$	$(1465\pm18)-i(101\pm20)/2$
$f_0(1710)$	$(1763 \pm 23) - i(104 \pm 34)/2$	$(1769 \pm 8) - i(156 \pm 12)/2$	$(1765 \pm 15) - i(180 \pm 20)/2$	/
$f_0(2020)$	$(1983 \pm 31) - i(143 \pm 54)/2$	$(2038 \pm 48) - i(312 \pm 82)/2$	$ \begin{array}{l} (1925\pm25)-i(320\pm35)/2\\ (2075\pm20)-i(260\pm25)/2 \end{array} $	$(1901 \pm 41) - i(401 \pm 76)/2$

- ▶ extracted poles (likely associated with $f_0(1500)$, $f_0(1710)$ and $f_0(2020)$) in good agreement with literature
- other determinations include information from other channels, SPM only uses $J/\psi \rightarrow \gamma \pi^0 \pi^0$ data!

[D. Binosi, A. Pilloni, R.-A. T., Phys. Lett. B 839, 137809 (2023)] [JPAC Collaboration: A. Rodas, A. Pilloni, M. Albaladejo, C. Fernandez-Ramirez, V. Mathieu, A. P. Szczepaniak, Eur.Phys.J.C 82, 80 (2022)] [Bonn-Gatchina: A. V. Sarantsev, I. Denisenko, U. Thoma, E. Klempt, Phys. Lett. B816, 136227 (2021)] [S. Ropertz, C. Hanhart, B. Kubis, Eur. Phys. J. C78, 1000 (2018)] We presented a data-driven method for the determination of complex poles associated to resonances from experimental data:

- based on the Schlessinger Point Method (SPM) which interpolates a given data set by a continued-fraction expression
- > analytic continuation to the complex plane allows to study pole structure
- ▶ SPM represents a ('model-independent') analysis technique for the determination of complex poles

Outlook:

• extension to coupled-channel datasets (S- and D-wave for $J/\psi \rightarrow \gamma \pi^0 \pi^0$ and $J/\psi \rightarrow \gamma K_S^0 K_S^0$)